

COMPOSITE PARTICLES IN MEDIUM — EFFECTS OF SUBSTRUCTURE*

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The role of phase-space occupation effects for the formation of two- and three-particle bound states in a dense medium is investigated for systems with short-range interactions. While for two-fermion bound states due to the Pauli blocking in a dense medium the binding energy is reduced and vanishes at a critical density (Mott effect), for three-fermion bound states, it is shown to be nonzero and positive. Therefore, beyond the Mott density of the two-fermion bound state, three-fermion bound states can exist in a medium and, therefore, be denoted as the in-medium Borromean states.

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In strongly coupled systems where the interaction saturates at short distances, the effects of phase-space occupation (Pauli blocking and Bose enhancement) are dominating medium effects. They can be nicely discussed in the algebraic Lipkin model [1] by generalising it for fermion–boson pairs (composite fermions) as compared to boson–boson pairs [2]. The problem of stability of three-particle bound states in dense matter is interesting for applications like the dissociation of baryons in dense quark/nuclear matter [3–6] and to fermionic atoms in traps [7].

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In the present note, we discuss the case of composite fermions in medium by considering the phase-space occupation factor that results from the Matsubara summation in the intermediate propagation in the N -particle T-matrix approach. The N -particle T-matrix in ladder approximation fulfils the Bethe–Goldstone equation

$$T_N = V_N + V_N G_N^0 T_N = V_N (1 - V_N G_N^0)^{-1}, \quad (1)$$

where V_N is the interaction potential and G_N^0 is the free N -particle Greens function in which the phase-space occupation effects become apparent. Let us examine the case of 2-particle and 3-particle states. The free two-fermion propagator in medium depends on the bosonic Matsubara frequency $\Omega_{12} = \omega_1 + \omega_2$ and is obtained by performing the Matsubara summation over ω_1

$$G_2^0(\Omega, e_1, e_2) = \sum_{\omega_1} \frac{1}{\omega_1 - e_1} \frac{1}{\Omega_{12} - \omega_1 - e_2} = \frac{Q_{12}}{\Omega_{12} - e_{12}}. \quad (2)$$

The energy denominator has a pole at $\Omega_{12} = e_{12} = e_1 + e_2$ and in the numerator occurs the phase-space occupation (Pauli blocking) factor $Q_{12} = 1 - f_1 - f_2$ with the Fermi functions $f_i = [\exp(e_i/T) + 1]^{-1}$. The free three-particle propagator is obtained by considering a pair of fermion and (composite) boson with the fermionic Matsubara frequency $\Omega_{123} = \Omega_{12} + \omega_3$

$$\begin{aligned} G_3^0(\Omega, e_1, e_2, e_3) &= \sum_{\omega_3} \frac{1}{\omega_3 - e_3} \frac{Q_{12}}{\Omega_{123} - \omega_3 - e_{12}} = \frac{(1 - f_3 + g_{12})Q_{12}}{\Omega - e_{12} - e_3} \\ &= \frac{Q_{123}}{\Omega_{123} - e_1 - e_2 - e_3}, \end{aligned} \quad (3)$$

where $Q_{123} = 1 - f_1 - f_2 - f_3 + f_1 f_2 + f_1 f_3 + f_2 f_3$ is the three-particle phase-space occupation factor and we have used the identity $g_{12}(1 - f_1 - f_2) = f_1 f_2$.

The in-medium Borromean property of three-fermion states can be seen in the simple example when $f_1 = f_2 = 0.5$, leading to a blocking of the two-particle state, $Q_{12} = 0$ (Mott effect). The three-particle state, however, has for the same case only a reduction of the effective coupling since $Q_{123} = 0.25 \neq 0$ [6], but can still be bound. In a next step, for a liquid of composite fermions such as nucleons, one obtains a weakly attractive boson exchange interaction that leads to fermionic superfluidity, as in nuclear matter [7].

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