

## INCLUSIVE HADRON-JET PRODUCTION AT THE LHC\*

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We suggest the inclusive detection at the LHC of a light charged hadron and of a jet widely separated in rapidity as a new probe channel for the study of the BFKL resummation. Predictions for cross section and azimuthal correlations, shaped on the CMS and CASTOR acceptances, are presented.

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### 1. Introduction

Semi-hard reactions represent a challenging testfield for perturbative QCD in the high-energy limit. Here, the fixed-order description undershoots the effect of large energy logarithmic contributions, which balance the smallness of the QCD coupling constant and hence must be resummed to all orders. The Balitsky–Fadin–Kuraev–Lipatov (BFKL) [1] approach is the most powerful tool to resum to all orders these large logarithms both in the leading (LLA) and the next-to-leading (NLA) approximation. In the last years, a good number of processes have been proposed as candidate probes of the high-energy regime, namely: the inclusive hadroproduction of two jets with high  $p_T$  and well-separated in rapidity (better known as the Mueller–Navelet process [2]), for which a richness of theoretical predictions have appeared so far [3–16], the inclusive detection of two identified, light charged hadrons

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[17–19], and more recently, multi-jet hadroproduction [20–26] and heavy-quark pair photoproduction [27]. In this work, a novel semi-hard reaction, *i.e.* the concurrent detection of a light charged hadron,  $\pi^\pm, K^\pm, p(\bar{p})$ , and a jet, both with high  $p_T$  and separated by a large rapidity interval, is proposed and investigated in the NLA BFKL approach. Although being, *de facto*, a hybridization of two already studied processes, this channel presents some own characteristics which can make it worthy of consideration in future analyses at the LHC. First, the tag of two different kinds of final-state object leads to a natural asymmetric configuration in  $p_T$ , allowing us to suppress the Born contribution and enhancing the discrepancy between BFKL and DGLAP, as pointed out in [10, 13, 28]. Second, having just one final-state hadron identified, instead of two, should quench “minimum-bias” contaminations, thus facilitating the comparison with experimental analyses. Third, one can use this reaction to compare different sets of fragmentation functions (FFs) and jet algorithms, dealing with *linear* expressions in the corresponding functions.

## 2. Inclusive hadron-jet production at the LHC

The process under investigation is<sup>1</sup>

$$\text{proton}(p_1) + \text{proton}(p_2) \rightarrow \text{hadron}(k_H, y_H) + X + \text{jet}(k_J, y_J), \quad (1)$$

when a light charged hadron,  $\pi^\pm, K^\pm, p(\bar{p})$ , and a jet, featuring large transverse momenta,  $k_{H,J} \gg \Lambda_{\text{QCD}}$ , and well-separated in rapidity,  $Y \equiv y_H - y_J$ , are produced together with an undetected hadronic system, X. The differential cross section of the process can be presented as

$$\frac{d\sigma}{dy_H dy_J d|\vec{k}_H| d|\vec{k}_J| d\phi_H d\phi_J} = \frac{1}{(2\pi)^2} \left[ \mathcal{C}_0 + \sum_{n=1}^{\infty} 2 \cos(n\phi) \mathcal{C}_n \right], \quad (2)$$

where  $\phi \equiv \phi_H - \phi_J - \pi$ , with  $\phi_{H,J}$  the hadron/jet azimuthal angles, while  $\mathcal{C}_0$  is the  $\phi$ -averaged cross section and the other coefficients  $\mathcal{C}_n$  determine the azimuthal-angle distribution of the final state. In order to match realistic kinematic configurations used at the LHC, we integrate the coefficients over the phase space for two detected objects and keep fixed the rapidity interval,  $Y$ , between the hadron and the jet

$$C_n = \int_{y_H^{\min}}^{y_H^{\max}} dy_H \int_{y_J^{\min}}^{y_J^{\max}} dy_J \int_{k_H^{\min}}^{k_H^{\max}} dk_H \int_{k_J^{\min}}^{k_J^{\max}} dk_J \delta(y_H - y_J - Y) \mathcal{C}_n. \quad (3)$$

<sup>1</sup> This process has much in common with the inclusive  $J/\Psi$ -meson plus backward jet production, considered recently in [29].

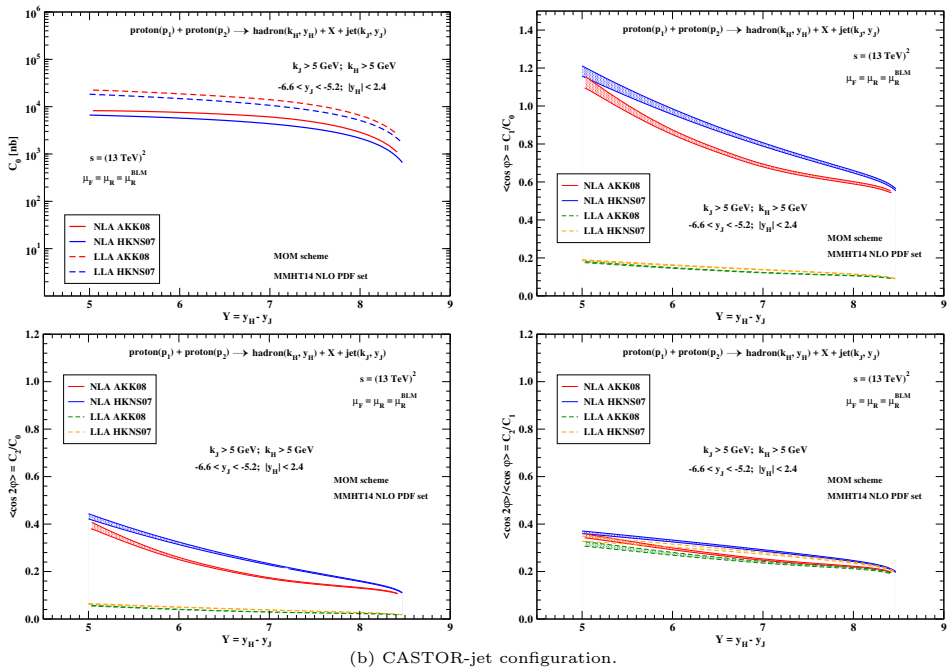
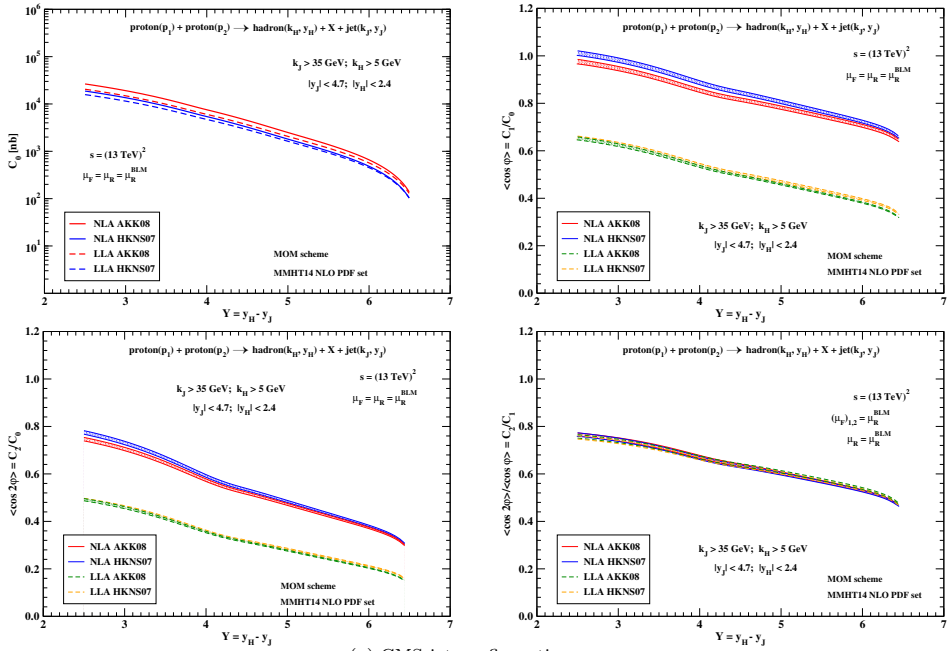


Fig. 1.  $Y$ -dependence of  $C_0$  and of several ratios  $C_m/C_n$  for  $\sqrt{s} = 13$  TeV.

We consider two distinct final-state ranges:

- (a) CMS-jet [30]: both the hadron and the jet tagged by the CMS detector in their typical kinematic configurations, *i.e.*:  $5 \text{ GeV} < k_H < 21.5 \text{ GeV}$ ,  $35 \text{ GeV} < k_J < 60 \text{ GeV}$ ,  $|y_H| \leq 2.4$ ,  $|y_J| \leq 4.7$ ;
- (b) CASTOR-jet [31]: a hadron always detected inside CMS in the range given above, together with a very backward jet tagged by CASTOR, lying in the range of  $5 \text{ GeV} < k_J \lesssim 17.68 \text{ GeV}$ ,  $-6.6 < y_J < -5.2$ .

We made all calculations with JETHAD, a Fortran code we recently developed, suited for the study of semi-hard reactions. We used the MMHT 2014 NLO PDF set [32] and two different NLO hadron FFs: AKK 2008 [33] and HKNS 2007 [34]. We took  $\mu_F = \mu_R$  and used the Brodsky–Lepage–Mackenzie (BLM) scheme [35] as derived in its “exact” version in [12]. All calculations were done in the MOM renormalization scheme. Predictions for  $C_0$  and for several ratios  $R_{nm} \equiv C_n/C_m$  at  $\sqrt{s} = 13 \text{ TeV}$  are shown in Fig. 1 (a) (CMS-jet) and Fig. 1 (b) (CASTOR-jet), while in Fig. 2 we compare  $C_0$  in different NLA BFKL processes: Mueller–Navelet jet, hadron-jet and di-hadron production at  $\sqrt{s} = 7, 13 \text{ TeV}$  in the CMS-jet case. We refer to Sections 3.3 and 3.4 of [36] for a detailed discussion of results, numerical tools and estimation of the uncertainties.

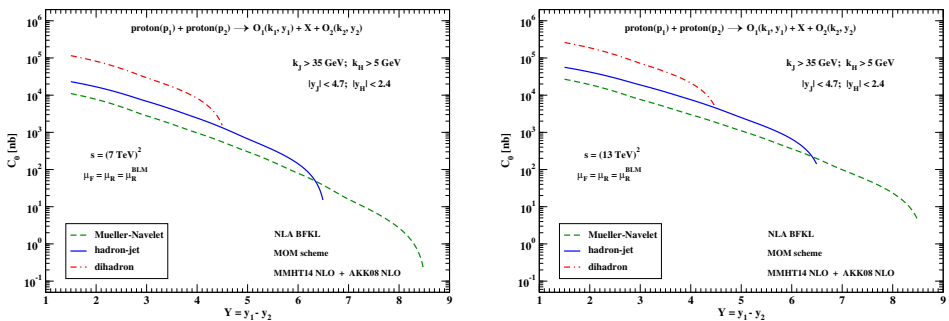


Fig. 2. Comparison of the  $\phi$ -averaged cross section  $C_0$  in different NLA BFKL processes for  $\sqrt{s} = 7, 13 \text{ TeV}$  in the CMS-jet configuration.

### 3. Summary

A new candidate probe of the BFKL mechanism, *i.e.* the inclusive hadron-jet production at the LHC, has been studied in the NLA accuracy. Distributions over the final-state rapidity interval are in accordance with the ones found for previously investigated semi-hard reactions, when the jet is detected by CMS, while new and interesting aspects emerged when the jet is tagged by CASTOR, which demand further, dedicated analyses.

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