INCLUSIVE CROSS SECTIONS IN THE DUAL PARTON MODEL FOR PIONS AND NEUTRON TARGETS

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We compare the dual parton model predictions for $\pi^-\pi^- \rightarrow \pi^{\pm}X$ and $\pi^-n \rightarrow \pi^{\pm}X$ with the available experimental data. A good agreement is achieved.

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The dual parton model (DPM) [1] has been extensively compared with the experimental data [2]. The model gives a good description of a wide variety of data: rising of the inclusive cross sections [3], charge and strangeness distributions [4], long range correlations [5], dependence of the average transverse momentum on the multiplicity [6] and multiplicity distributions [7]. Also the extension of the model to hadron-nucleus [8] and nucleus--nucleus [9] interactions is in agreement with the experimental data. In DPM a valence quark from each of the two initial hadrons has, in average a small momentum fraction x. Then, the hadrons produced near y = 0 must come mainly from the fragmentation of the "held back" valence quarks. This effect has been studied for different beams using protons as targets in different inclusive cross sections [10]. In this note we study the cases where the targets are pions and neutrons. We use the data of references [11] and [12]. The inclusive pion target data were obtained from the momentum distributions of pions emitted from the recoiling system X⁻⁻ in the reaction $\pi^-n \rightarrow pX^{--}$ assuming the one pion exchange model. The considered inclusive cross sections are $\pi^-\pi^- \rightarrow \pi^+X$, $\pi^-\pi^- \rightarrow \pi^-X$, $\pi^-n \rightarrow \pi^-X$ and $\pi^-n \to \pi^+X$. For the two first reactions, the contribution comes from the diagrams of Fig. 1 which are given by the formula

$$\frac{dN^{\pi^-\pi^-\to\pi X}}{dy}(s, y) = \int_0^1 dx_1 \int_0^1 dx_2 \varrho_{d\tilde{u}}^{\pi^-\pi^-}(x_1, x_2)$$

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$$\times \left[\frac{dN_{\rm S}^{({\rm d},\bar{\rm u})\to\pi}}{dy} (y-\Delta_{\rm S},P_{\rm S}) + \frac{dN_{\rm L}^{(\bar{\rm u},{\rm d})\to\pi}}{dy} (y-\Delta_{\rm L},P_{\rm L}) \right] + \int_{0}^{1} dx_{1} \int_{0}^{1} dx_{2} \varrho_{\rm u\bar{u}}^{\pi^{-}\pi^{-}} (x_{1},x_{2}) \left[\frac{dN_{\rm A}^{({\rm d},\bar{\rm u})\to\pi}}{dy} (y-\Delta_{\rm A},P_{\rm A}) + \frac{dN_{\rm B}^{(\bar{\rm u},{\rm d})\to\pi}}{dy} (y-\Delta_{\rm B},P_{\rm B}) \right]$$
(1)

where we use the standard notation.



Fig. 1. Diagrams contributing to $\pi^-\pi^- \rightarrow \pi X$

For the structure functions we use the standard form

$$\varrho_{\overline{uu}}^{\pi^-\pi^-}(x_1, x_2) = \varrho_{\overline{du}}^{\pi^-\pi^-}(x_1, x_2) = 0.75 \frac{x_1}{\sqrt{1-x_1}} 0.75 \frac{x_2}{\sqrt{1-x_2}}$$
(2)

and the chain distributions are obtained in the usual way from the fragmentation functions

$$D_{u \to \pi^+}(x) = D_{\bar{d} \to \pi^+} = D_{\bar{u} \to \pi^-} = D_{d \to \pi^-} = \frac{cf(x)}{1 + w(x)},$$
(3)

$$D_{u \to \pi^{-}}(x) = D_{\bar{d} \to \pi^{-}} = D_{\bar{u} \to \pi^{+}} = D_{d \to \pi^{+}} = \frac{cf(x)w(x)}{1+w(x)},$$
(4)

where

$$f(x) = \frac{(1-x)^2}{1-0.5x}, \quad w(x) = \frac{1-x}{1+x}$$
(5)

and c = 1.1.

In Fig. 2 we plot the inclusive $\pi^-\pi^- \rightarrow \pi^+X$ cross section at $P_{lab} = 360 \text{ GeV}/c$ together with the experimental data. In Fig. 3 we present the ratio between the inclusive cross sections for π^- and π^+ . From the figures we see a remarkable agreement between theory and experiment.

For the $\pi^-n \to \pi^+X$ and $\pi^-n \to \pi^-X$ cross sections we proceed in an analogous way. The obtained $\pi^-n \to \pi^-X$ inclusive cross sections at 360 GeV/c are plotted in Fig. 4 together



Fig. 2. $\pi^+\pi^- \rightarrow \pi^+ X$ inclusive cross section at $P_{\text{Lab}} = 360 \text{ GeV}/c$



Fig. 3. Ratio between $\pi^-\pi^- \rightarrow \pi^- X$ and $\pi^-\pi^- \rightarrow \pi^+ X$ inclusive cross sections at $P_{Lab} = 360 \text{ GeV}/c$



with the experimental data. In Fig. 5 we plot the ratio between the inclusive cross sections for π^- and π^+ at 360 GeV/c. Again a remarkable agreement between theory and experiment is achieved. This agreement is also obtained at the other two energies 21 GeV/c and 205 GeV/c, where there are experimental data. A sample of that is shown in Fig. 4.

In conclusion, the DPM describes rightly the inclusive cross section $\pi^-\pi^- \rightarrow \pi^{\pm}X$ and $\pi^-n \rightarrow \pi^{\pm} X$. Notice that recombination type models [13–14] have problems [12] in the description of the same data.

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