

# GEOMETRIC OBJECTS RELATED TO THE POTENTIAL OF ELECTRIC CHARGES

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*Dedicated to the memory of Professor Jan Rzewuski*

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We derive explicit formulas for curvature and torsion of a line of the field of  $n$  electric charges. These formulas show that in general the torsion of a field line is not zero if  $n \geq 3$ . We also propose a geometric interpretation of the derived formulas. In the second part of the paper we present an outline of a new description of equipotential surfaces of two and three electric charges. In this description the golden section appears in a natural way when two electric charges are equal. This approach also relates an equipotential surface of three charges to the classic cubic surface containing twenty seven straight lines.

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## 1. Introductory remarks

This paper consists of two parts. In the first part (Section 2) we follow [1] in deriving explicit formulas for curvature and torsion of a line of the field of  $n$  electric charges. In particular, these formulas imply that for two charges the torsion is always zero, and in general, is not zero when  $n \geq 3$ . Section 3 gives a geometric interpretation of the formulas derived in Section 2.

The second part (Sections 4 and 5) contains an outline of a new description of equipotential surfaces of two and three electric charges. This description is based on the notion of an auxiliary line  $L_2$  and an auxiliary surface  $L_3$ . The auxiliary objects  $L_2$  and  $L_3$  give a geometric insight into a natural parametrization of equipotential surfaces of two and three electric charges. In Section 4 we derive formulas which show that the famous golden section, [2], appears when two charges are equal. In Section 5, making use of the finite Fourier transform  $F(3)$ , [3], we relate the auxiliary surface  $L_3$  to the classic cubic surface containing twenty seven straight lines, [4–6].

## 2. Curvature and torsion of a field line of $n$ electric charges

We use the following notation:

- $q_i$  ( $i = 1, 2, \dots, n$ ) —  $i$ -th electric charge,  
 $\vec{r}(x, y, z)$  — position vector of the point  $(x, y, z)$ ,  
 $\vec{r}_i = (x_i, y_i, z_i)$  — position vector of  $q_i$ ,  
 $R_i = |\vec{R}_i| = |\vec{r} - \vec{r}_i|$  — distance between  $(x, y, z)$  and  $q_i$ .

Then the potential  $\Phi$  and the electric field  $\vec{E}$  of  $q_1, q_2, \dots, q_n$  are given by the classic formulas

$$\Phi = \sum_{i=1}^n q_i R_i^{-1}, \quad (2.1)$$

$$\vec{E} = -\text{grad } \Phi. \quad (2.2)$$

Let

$$\vec{r}(t) = (x(t), y(t), z(t)) \quad (2.3)$$

be a line parametrized by  $t$ . It is well known that the line  $\vec{r}(t)$  is by definition a line of the field  $\vec{E}$  if  $\vec{E}(\vec{r}(t))$  is tangent to the line  $\vec{r}(t)$  for any value of the parameter  $t$ :

$$\vec{r}'(t) = \vec{E}(\vec{r}(t)), \quad (2.4)$$

where

$$\vec{r}'(t) = \frac{d\vec{r}}{dt}. \quad (2.5)$$

The curvature (first curvature)  $k_1$  and torsion (second curvature)  $k_2$  of a line  $\vec{r}(t)$  are given by the well known formulas (see, for example, [7, 8]):

$$k_1 = \frac{|\vec{r}' \times \vec{r}''|}{|\vec{r}'|^3}, \quad (2.6)$$

$$k_2 = -\frac{[\vec{r}', \vec{r}'', \vec{r}''']}{|\vec{r}' \times \vec{r}''|^2}, \quad (2.7)$$

where  $[\vec{r}', \vec{r}'', \vec{r}''']$  denotes the mixed product of  $\vec{r}'$ ,  $\vec{r}''$  and  $\vec{r}'''$ :

$$[\vec{r}', \vec{r}'', \vec{r}'''] = (\vec{r}' \times \vec{r}'') \cdot \vec{r}'''. \quad (2.8)$$

From the above formulas it follows that the curvature  $k_1$  and torsion  $k_2$  can be written in the form

$$k_1 = \frac{|\vec{E} \times \vec{F}|}{|\vec{E}|^3}, \quad (2.9)$$

$$k_2 = -\frac{[\vec{E}, \vec{F}, \vec{G}]}{|\vec{E} \times \vec{F}|^2}, \quad (2.10)$$

where

$$\vec{F} = \sum_{i=1}^n b_i \vec{R}_i, \quad (2.11)$$

$$\vec{G} = \sum_{i=1}^n c_i \vec{R}_i, \quad (2.12)$$

$$a_i = q_i R_i^{-3}, \quad (2.13)$$

$$b_i = a'_i = -3R_i^{-2} a_i A_i, \quad (2.14)$$

$$c_i = a''_i = -3R_i^{-2} a_i [(-5R_i^{-2} A_i + A) A_i + B_i + E^2], \quad (2.15)$$

$$A_i = \vec{R}_i \vec{E}, \quad (2.16)$$

$$B_i = \vec{R}_i \vec{F}, \quad (2.17)$$

$$A = \sum_{i=1}^n a_i. \quad (2.18)$$

It is easy to verify that the vector  $\vec{E} \times \vec{F}$  and the mixed product  $[\vec{E}, \vec{F}, \vec{G}]$  can be written in the form

$$\vec{E} \times \vec{F} = \sum_{\substack{i,1=1 \\ (i < j)}} \mu_{ij} \vec{S}_{ij}, \quad (2.19)$$

$$[\vec{E}, \vec{F}, \vec{G}] = \sum_{\substack{i,j,k=1 \\ (i < j < k)}}^n \sigma_{ijk} V_{ijk}, \quad (2.20)$$

where

$$\mu_{ij} = \begin{vmatrix} a_i & a_j \\ b_i & b_j \end{vmatrix}, \quad (2.21)$$

$$\sigma_{ijk} = \begin{vmatrix} a_i & a_j & a_k \\ b_i & b_j & b_k \\ c_i & c_j & c_k \end{vmatrix}, \quad (2.22)$$

$$\vec{S}_{ij} = \vec{R}_i \times \vec{R}_j, \quad (2.23)$$

$$V_{ijk} = [\vec{R}_i, \vec{R}_j, \vec{R}_k]. \quad (2.24)$$

### 3. A geometric interpretation of the formulas (2.9)–(2.10) as (2.19)–(2.20)

We introduce the following terminology:

- $a_i \vec{R}_i$  — the first scaling of  $\vec{R}_i$ ,  
 $b_i \vec{R}_i$  — the second scaling of  $\vec{R}_i$ ,  
 $c_i \vec{R}_i$  — the third scaling of  $\vec{R}_i$ ,

$$\left. \begin{aligned} \vec{E} &= \sum_{i=1}^n a_i \vec{R}_i \\ \vec{F} &= \sum_{i=1}^n b_i \vec{R}_i \\ \vec{G} &= \sum_{i=1}^n c_i \vec{R}_i \end{aligned} \right\} \begin{array}{l} \text{Resulting vectors of the first, second} \\ \text{and third scaling of } \vec{R}_1, \dots, \vec{R}_n. \end{array}$$

- $P(\vec{A}, \vec{B})$  — parallelogram spanned on the vectors  $\vec{A}$  and  $\vec{B}$ ,  
 $P(\vec{A}, \vec{B}, \vec{C})$  — parallelepiped spanned on the vectors  $\vec{A}$ ,  $\vec{B}$  and  $\vec{C}$ ,  
 $\vec{A} \times \vec{B}$  — surface vector of  $P(\vec{A}, \vec{B})$ ,  
 $[\vec{A}, \vec{B}, \vec{C}]$  — volume of  $P(\vec{A}, \vec{B}, \vec{C})$ ,  
 $|\vec{S}_{ij}| = |\vec{R}_i \times \vec{R}_j|$  — elementary surface,  
 $V_{ijk} = [\vec{R}_i, \vec{R}_j, \vec{R}_k]$  — elementary volume,  
 $\mu_{ij} \vec{S}_{ij}$  — surface scaling of  $\vec{S}_{ij}$ ,  
 $\sigma_{ijk} V_{ijk}$  — volume scaling of  $V_{ijk}$ ,  
 $NES(n) = \binom{n}{2}$  — number of elementary surfaces,  
 $NEV(n) = \binom{n}{3}$  — number of elementary volumes,

Making use of this terminology we obtain the following geometric interpretation of the formulas (2.9)–(2.10) and (2.19)–(2.20):

- (i) The curvature  $k_1$  is equal to the ratio of the surface  $|\vec{E} \times \vec{F}|$  to the third power of the length of  $\vec{E}$ .  
(ii) The torsion  $k_2$  is equal to minus the ratio of the volume  $[\vec{E}, \vec{F}, \vec{G}]$  to the square of the surface  $|\vec{E} \times \vec{F}|$ .  
(iii) The volume  $[\vec{E}, \vec{F}, \vec{G}]$  consists of  $\binom{n}{3}$  scaled elementary volumes  $\sigma_{ijk} V_{ijk}$ .  
(iv) The surface vector  $\vec{E} \times \vec{F}$  consists of  $\binom{n}{2}$  scaled elementary surface vectors  $\mu_{ij} \vec{S}_{ij}$ .

$$(v) \quad NEV(n) - NES(n) = \begin{cases} -2 & \text{if } n = 3, 4 \\ 0 & \text{if } n = 5 \\ > 0 & \text{if } n > 5 \end{cases}$$

(vi) If  $n = 1$  (one electric charge  $q_1$ ), then  $k_1 = k_2 = 0$  because the vectors  $\vec{E}$ ,  $\vec{F}$ ,  $\vec{G}$  are then colinear:

$$\begin{aligned}\vec{E} &= a_1 \vec{R}_1, \\ \vec{F} &= b_1 \vec{R}_1, \\ \vec{G} &= c_1 \vec{R}_1.\end{aligned}$$

(vii) If  $n = 2$  (two electric charges  $q_1$  and  $q_2$ ), then in general  $k_1 \neq 0$ , but  $k_2 = 0$  because the vectors  $\vec{E}$ ,  $\vec{F}$ ,  $\vec{G}$  are then linear combinations of  $\vec{R}_1$  and  $\vec{R}_2$ , and so are linearly dependent.

#### 4. The golden section and an equipotential surface of two electric charges

In this Section we consider the surface of a fixed potential  $\lambda$  (the  $\lambda$ -potential surface) of two electric charges  $q_1$  and  $q_2$ . Then the formula (2.1) takes the form

$$\lambda R_1 R_2 = q_1 R_2 + q_2 R_1. \quad (4.1)$$

We choose the position vectors  $\vec{r}_1$  and  $\vec{r}_2$  in the following way

$$\begin{cases} \vec{r}_1 = (0, 0, a), \\ \vec{r}_2 = (0, 0, 0). \end{cases}$$

Then the  $\lambda$ -potential surface is invariant under any rotation round the  $z$ -axis, and so we can restrict ourselves, for example, to the plane  $x = 0$ . For this choice

$$\begin{aligned}R_1 &= s = (w^2 - 2az + a^2)^{1/2}, \\ R_2 &= w = (y^2 + z^2)^{1/2},\end{aligned} \quad (4.3)$$

and the equation (4.1) can be rewritten in the form

$$uv = D_1 v + D_2 u, \quad (4.4)$$

where

$$u = \frac{s}{a}, \quad v = \frac{w}{a}, \quad (4.5)$$

$$D_i = \frac{q_i}{a\lambda} \quad (i = 1, 2). \quad (4.6)$$

Introducing new variables

$$\begin{aligned}\xi &= \frac{1}{2}(u + v), \\ \eta &= \frac{1}{2}(u - v),\end{aligned} \quad (4.7)$$

and notation

$$\begin{aligned} E_1 &= \frac{1}{2}(D_1 + D_2), \\ E_2 &= \frac{1}{2}(D_1 - D_2), \end{aligned} \quad (4.8)$$

we obtain from (4.4) the following equation of hyperbola

$$(\xi - E_1)^2 - (\eta - E_2)^2 = D_1 D_2. \quad (4.9)$$

We call this hyperbola the auxiliary line  $L_2$  of the  $\lambda$ -equipotential surface of  $q_1$  and  $q_2$ . Let

$$P^2 = \begin{cases} D_1 D_2 & \text{if } q_1 q_2 > 0, \\ -D_1 D_2 & \text{if } q_1 q_2 < 0. \end{cases} \quad (4.10)$$

Then the equation (4.9) can be rewritten in the following two forms

$$(\xi - E_1)^2 - (\eta - E_2)^2 = P^2 \quad \text{if } q_1 q_2 > 0, \quad (4.11)$$

$$(\eta - E_2)^2 - (\xi - E_1)^2 = P^2 \quad \text{if } q_1 q_2 < 0. \quad (4.12)$$

Here we restrict ourselves to the first case ( $q_1 q_2 > 0$ ). The other case ( $q_1 q_2 < 0$ ) is discussed in [9]. Let us note that if  $q_1 q_2 > 0$ , then  $D_1$  and  $D_2$  are always positive. Indeed, by the formula (4.1)

$$q_1 > 0, \quad q_2 > 0 \Rightarrow \lambda > 0, \quad (4.13)$$

$$q_1 < 0, \quad q_2 < 0 \Rightarrow \lambda < 0, \quad (4.14)$$

and so, by (4.6),  $D_i > 0$  ( $i = 1, 2$ ).

The formula (4.11) implies that the auxiliary line  $L_2$  can be parametrized in terms of hyperbolic functions  $\cosh \gamma$  and  $\sinh \gamma$ :

$$\begin{cases} \xi - E_1 = P \cosh \gamma, \\ \eta - E_2 = P \sinh \gamma, \\ P = (D_1 D_2)^{1/2}. \end{cases} \quad (4.15)$$

Taking into account the formulas (4.7), (4.8) and (4.15) we can write the old variables  $u$  and  $v$  in the form

$$\begin{cases} u = D_1 + P e^\gamma, \\ v = D_2 + P e^{-\gamma}. \end{cases} \quad (4.16)$$

The formula (4.16) and the geometric meaning of  $u$  and  $v$  imply that

$$\begin{cases} \gamma_1 \leq \gamma \leq \gamma_2, \\ D_1 + D_2 + 2P \geq 1. \end{cases} \quad (4.17)$$

The lowest value  $\gamma_1$  and the highest value  $\gamma_2$  of  $\gamma$  can be determined from the following two conditions

$$u_1 + 1 = v_1, \quad (4.18)$$

$$u_2 = v_2 + 1, \quad (4.19)$$

where

$$\begin{cases} u_i = u|_{\gamma=\gamma_i}, \\ v_i = v|_{\gamma=\gamma_i}, \\ i = 1, 2. \end{cases} \quad (4.20)$$

Then  $\gamma_i$ ,  $u_i$  and  $v_i$  ( $i = 1, 2$ ) are given by the formulas

$$\begin{aligned} e^{\gamma_1} &= \frac{1}{2P} \{ D_2 - D_1 - 1 + [(D_2 - D_1 - 1)^2 + 4P^2]^{1/2} \}, \\ e^{\gamma_2} &= \frac{1}{2P} \{ D_2 - D_1 + 1 + [(D_2 - D_1 + 1)^2 + 4P^2]^{1/2} \}, \end{aligned} \quad (4.21)$$

$$\begin{aligned} u_1 &= \frac{1}{2} \{ D_1 + D_2 - 1 + [(D_2 - D_1 - 1)^2 + 4P^2]^{1/2} \}, \\ v_1 &= \frac{1}{2} \{ D_1 + D_2 + 1 + [(D_2 - D_1 - 1)^2 + 4P^2]^{1/2} \}, \\ u_2 &= \frac{1}{2} \{ D_1 + D_2 + 1 + [(D_2 - D_1 + 1)^2 + 4P^2]^{1/2} \}, \\ v_2 &= \frac{1}{2} \{ D_1 + D_2 - 1 + [(D_2 - D_1 + 1)^2 + 4P^2]^{1/2} \}. \end{aligned} \quad (4.22)$$

Let us denote that  $au_1$  is the minimal value  $s_{\min}$  of  $s$ , and  $av_2$  is the minimal value  $w_{\min}$  of  $w$ .

Let  $D_1 = D_2 = 1$ . Then the formula (4.22) implies the following proportions of the golden section, [2]:

$$\frac{s_{\min}}{a} = \frac{w_{\min}}{a} = \frac{1}{2}(1 + \sqrt{5}). \quad (4.23)$$

## 5. Twenty seven straight lines and the auxiliary surface $L_3$ of three electric charges

Let  $C(3)$  denote the following cubic surface (in affine coordinates  $\eta_1, \eta_2, \eta_3$ ):

$$\eta_1^3 + \eta_2^3 + \eta_3^3 + f = 0. \quad (5.1)$$

The surface  $C(3)$  has remarkable properties. In particular  $C(3)$  contains 27 straight lines, [4–6]. The configuration of the 27 straight lines of  $C(3)$  is invariant with respect to the group  $G$  which is isomorphic to the Weyl group

of the exceptional Lie algebra  $E_6$ , i.e.  $G$  has 51840 elements and contains a simple group of order 25920, [10].

According to the formula (2.1), the  $\lambda$ -potential surface (equipotential surface of a fixed potential  $\lambda$ ) of three electric charges  $q_1$ ,  $q_2$  and  $q_3$  is given by the equation

$$\lambda R_1 R_2 R_3 = q_1 R_2 R_3 + q_2 R_1 R_3 + q_3 R_1 R_2. \quad (5.2)$$

The auxiliary surface  $L_3$  is by definition the surface obtained from (5.2) by its complexification ( $R_1$ ,  $R_2$  and  $R_3$  are then complex variables). In [9] we show that

- (i)  $L_3$  gives a geometric insight into a natural parametrization of the real  $\lambda$ -surface of  $q_1$ ,  $q_2$  and  $q_3$ .
- (ii)  $L_3$  can be related to  $C(3)$  via the following change of variables

$$\begin{pmatrix} R_1 - d_1 \\ R_2 - d_2 \\ R_3 - d_3 \end{pmatrix} = \frac{\sqrt{3}}{(d_1 d_2 d_3)^{\frac{1}{3}}} \begin{pmatrix} d_1 & 0 & 0 \\ 0 & d_2 \varepsilon_3 & 0 \\ 0 & 0 & d_3 \varepsilon_3^2 \end{pmatrix} F(3) \begin{pmatrix} \eta_1 \\ \eta_2 \\ \eta_3 \end{pmatrix}, \quad (5.3)$$

where

$$F(3) = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 1 & 1 \\ 1 & \varepsilon_3 & \varepsilon_3^2 \\ 1 & \varepsilon_3^2 & \varepsilon_3 \end{pmatrix}, \quad (5.4)$$

$$d_j = \frac{q_j}{\lambda} \quad (j = 1, 2, 3), \quad (5.5)$$

$$\varepsilon_3 = e^{2\pi i/3}. \quad (5.6)$$

- (iii)  $L_3$ , written in the new variables  $\eta_1$ ,  $\eta_2$ ,  $\eta_3$  (defined by (5.3)) intersects  $C(3)$  along a hyperbolic helix.
- (iv) The 27 straight lines of  $C(3)$  intersect  $L_3$  forming an “acupuncture” of  $L_3$ .
- (v) There is no invertible inhomogeneous linear transformation which transforms  $L_3$  onto  $C(3)$ .

Finally, let us note that the matrix  $F(3)$ , given by (5.4), is the so called finite Fourier transform which has remarkable properties collected in [3].



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