

PRODUCTION OF B_c MESONS IN PHOTON-PHOTON
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We discuss two-photon and hadronic production of B_c mesons in non-relativistic bound state approximation and to lowest order in the coupling constants α and α_s . It is shown that in photon-photon collisions, heavy quark fragmentation is dominated by recombination of \bar{b} and c quarks up to the highest accessible transverse momenta. In contrast, in hadroproduction, which at high energies mainly involves gluon-gluon collisions, the fragmentation mechanism dominates at transverse momenta $p_T > m_{B_c}$, providing a simple and satisfactory approximation of the complete $O(\alpha_s^4)$ results in the high- p_T regime. Contradictions in previous publications on hadroproduction of B_c mesons are clarified. We also present predictions for cross sections and differential distributions at present and future accelerators.

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Since the top quark is too short lived for the formation of quarkonium-like resonances, B_c mesons are most probably the only flavoured heavy

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quark resonances in nature. Because of flavour conservation in strong and electromagnetic interactions, the B_c ground state must decay weakly. The nonrelativistic nature of these bound states provides unique possibilities to compute genuine nonperturbative quantities such as fragmentation functions and weak matrix elements, and to study interesting aspects of the strong and weak dynamics of hadrons.

A limit on B_c production has been reported recently by the CDF-collaboration at the Tevatron [1]. At the LHC, the production rates are predicted to be large enough for a detailed study of the production and decay properties [2]. Also at linear colliders in the TeV energy range, B_c mesons produced in collisions of Compton or bremsstrahlung photons may come into experimental reach [3].

In this talk, we report about two recent studies [2, 3] of B_c production. We discuss the fragmentation and recombination mechanisms and compare the relative importance of them in photon-photon and hadron-hadron collisions. As a main result, we clarify quantitatively the validity of the hard scattering description in terms of heavy quark fragmentation functions. Furthermore, on the basis of two completely independent calculations, we resolve contradictions in previous publications on hadronic production of B_c mesons [4–7]. Finally, we present the most relevant integrated cross sections and differential distributions as predicted in lowest-order perturbation theory and nonrelativistic bound state approximation.

In photon-photon collisions, B_c mesons are produced in association with b - and \bar{c} -quark jets:

$$\gamma\gamma \rightarrow B_c b \bar{c}. \quad (1)$$

This is also the case in hadronic collisions, where gluon-gluon scattering,

$$gg \rightarrow B_c b \bar{c}, \quad (2)$$

is the dominant subprocess at high energies. In general, one can distinguish two production mechanisms, namely heavy quark fragmentation and recombination. We have found these mechanisms to contribute quite differently in the two reactions, which makes a comparison very interesting.

The twenty $O(\alpha^2\alpha_s^2)$ Feynman diagrams of process (1) can be classified in three gauge invariant subsets characterized in Fig. 1: subset (I_b) with the b -quark line coupled to the primary photons, subset (I_c) with the c -quark line coupled to the primary photons, and subset (II) with the b -quark line coupled to one of the primary photons and the c -quark line to the other one. The three subsets can be interpreted physically as describing, respectively, $b\bar{b}$ production and subsequent \bar{b} fragmentation, $\bar{b} \rightarrow B_c \bar{c}$, $c\bar{c}$ production and subsequent c fragmentation, $c \rightarrow B_c b$, and simultaneous production of a $b\bar{b}$ and $c\bar{c}$ pair and recombination of the \bar{b} - and c -quark into a B_c meson.

The gluon-fusion process (2) involves a larger number of Feynman diagrams because of the presence of the gluon self-coupling. To $O(\alpha_s^4)$, one has thirty six diagrams in total. More importantly, they cannot be divided up in gauge invariant subsets corresponding to the production mechanisms typified in Fig. 1.

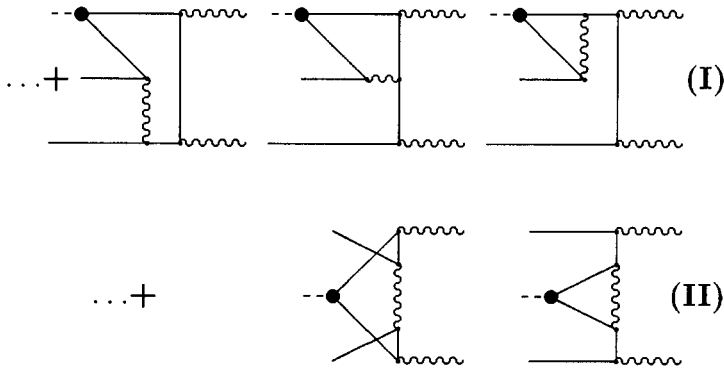


Fig. 1. Characteristic topologies of the lowest-order Feynman diagrams contributing to $\gamma\gamma \rightarrow B_c b\bar{c}$.

In the following, we concentrate on the production of the pseudoscalar (B_c) and vector (B_c^*) bound states. Furthermore, we consider the nonrelativistic limit, in which the relative momentum of the constituents and their binding energy are neglected relative to the b - and c -quark masses. In this limit, the masses of the bound states are equal to the sum of m_b and m_c , and the momenta of both constituent quarks are proportional to the bound state momentum. Moreover, the amplitudes for the production of S -waves factorize into hard scattering amplitudes for $\gamma\gamma$ (or gg) $\rightarrow b\bar{b}c\bar{c}$ times the S -wave function at the origin. The latter can be related to the $B_c^{(*)}$ decay constants. All this results in the simple substitution rule [8, 3, 2]

$$v(p_b)\bar{u}(p_c) = \frac{f_{B_c^{(*)}}}{\sqrt{48}}(\not{p} - M)\Pi_{SS_Z}, \quad (3)$$

indicated in Fig. 1 by the black blob, where $v(p_b)$ and $\bar{u}(p_c)$ denote the \bar{b} - and c -quark spinors, respectively, and $\Pi_{SS_Z} = \gamma_5(\not{p})$ is the spin projector for B_c (B_c^*). Note that the colour structure is not accounted for in Eq. (3).

The squared matrix elements for the processes (1) and (2) have been calculated independently using two different methods, the traditional trace technique and the method of helicity amplitudes. The results are found to be in perfect agreement. In addition, we have checked external gauge invariance, that is the vanishing of the matrix elements when the polarization

vector of any of the initial photons or gluons is substituted by its momentum. For process (1), we have also tested internal gauge invariance, that is the independence of the matrix element of the gauge parameter in the gluon propagator. Finally, we have double-checked the phase space integration using two different Monte Carlo routines. For more details of the calculation we refer to Refs [2] and [3].

In addition to the above $O(\alpha^2\alpha_s^2)$ and $O(\alpha_s^4)$ calculations, we have also studied the factorized description of the processes (1) and (2) in terms of $b\bar{b}(c\bar{c})$ production followed by $\bar{b}(c)$ fragmentation [12]:

$$d\sigma_{B_c} = d\sigma_{b\bar{b}} \otimes D_{\bar{b} \rightarrow B_c \bar{c}}(z), \quad (d\sigma_{c\bar{c}} \otimes D_{c \rightarrow B_c b}(z)). \quad (4)$$

The relevant fragmentation functions $D_{\bar{b}}(z)$ and $D_c(z)$ have been derived from perturbation theory [9]. They are known to provide a perfect approximation of the energy distribution $d\sigma/dz$ in $e^+e^- \rightarrow B_c \bar{c}b$ [9, 10] in order $\alpha^2\alpha_s^2$. Clearly, in photon-photon and hadron-hadron production such a factorized description cannot be expected to work close to threshold and at small p_T , where the quark masses play a role. The question is how well this approximation works at high- p_T . Previous studies of this issue only give qualitative and partly contradicting answers [4, 6]. This motivated us to compare the factorized approximation (4) with our complete lowest-order calculations, and to determine the region of validity of (4) quantitatively.

The numerical results plotted in Figs. 2–6 have been obtained with the following values of the parameters:

$$m_b = 4.8 \text{ GeV}, \quad m_c = 1.5 \text{ GeV}, \quad \alpha = \frac{1}{129}, \quad f_{B_c} = f_{B_c^*} = 0.4 \text{ GeV} \quad [11]. \quad (5)$$

Furthermore, we have used the running coupling constant $\alpha_s(Q^2)$ in leading logarithmic approximation for five flavours and normalized to $\alpha_s(m_Z^2) = 0.113$. Additional specifications are given when needed.

Fig. 2 illustrates transverse momentum distributions of the B_c in $\gamma\gamma$ -production (1). We see that the description (4) in terms of $b\bar{b}$ production and \bar{b} -fragmentation and the corresponding subset (I_b) of the diagrams of Fig. 1 give distributions (dotted and dashed curves, respectively, labeled by b) which are very similar in shape and normalization except in the low- p_T region, where the factorized approximation is expected to break down. As far as the shape is concerned this is also true for $c\bar{c}$ production and c fragmentation (curves labeled by c). However, in this case the approximation (4) fails to reproduce the correct magnitude of the cross section predicted by the diagrams subset (I_c) of Fig. 1. Note that primary c -quark production is enhanced to b -quark production by a factor 16 due to the ratio $(Q_c/Q_b)^4$ of the electric charges. On the other hand, the radiation of a $c\bar{c}$ -pair from a b -quark leads to a harder p_T spectrum for the B_c bound states than the

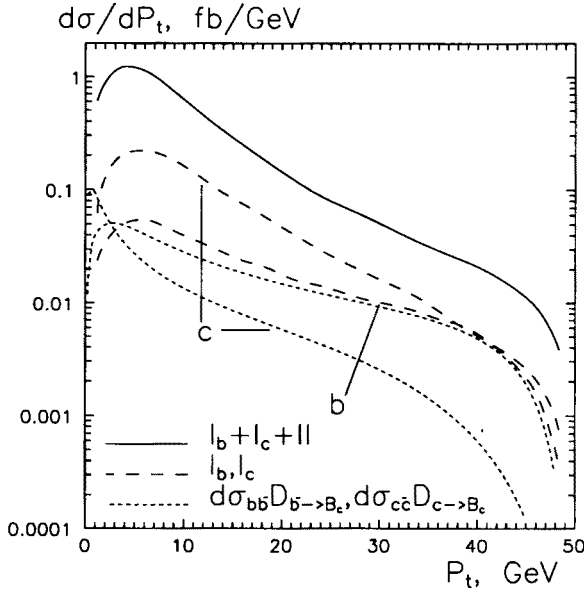


Fig. 2. Transverse momentum distributions of B_c mesons produced in $\gamma\gamma \rightarrow B_c b\bar{c}$ at $\sqrt{s} = 100$ GeV for the different production mechanisms and approximations described in the text.

radiation of a $b\bar{b}$ -pair from c -quarks. Finally, the most important observation is that the recombination mechanism, represented by the diagrams (II) of Fig. 1, dominates B_c production not only at low p_T , as one could have expected, but also in the high- p_T region up to the kinematical limit. In other words, the familiar description of high- p_T hadron production in terms of the production and fragmentation of quarks is inadequate for single B_c production in $\gamma\gamma$ -scattering.

The main features of the transverse momentum distributions of the B_c produced in the subprocess $gg \rightarrow B_c b\bar{c}$ are illustrated in Fig. 3. Here, we see that the fragmentation description (4) indeed approaches the p_T -distributions resulting from the complete $O(\alpha_s^4)$ calculation, but only in the tails of the distributions. In order to demonstrate the effect of the quark masses on the fragmentation kinematics, we have assumed three different relations between daughter and parent momenta:

$$p_T = z\sqrt{\frac{\hat{s}}{4} - \mu^2} \sin \theta_b, \quad p_L = p_T \cot \theta_b, \quad \mu = M \text{ (I)}, \quad m_b \text{ (II)}, \quad 0 \text{ (III)}, \quad (6)$$

where $\sqrt{\hat{s}}$ is the gluon-gluon c.m. energy. Case I obeys the physical phase space boundaries. The choices II and III have been considered in Refs [12] and [13], respectively. Fig. 3 shows that the mass ambiguities of the frag-

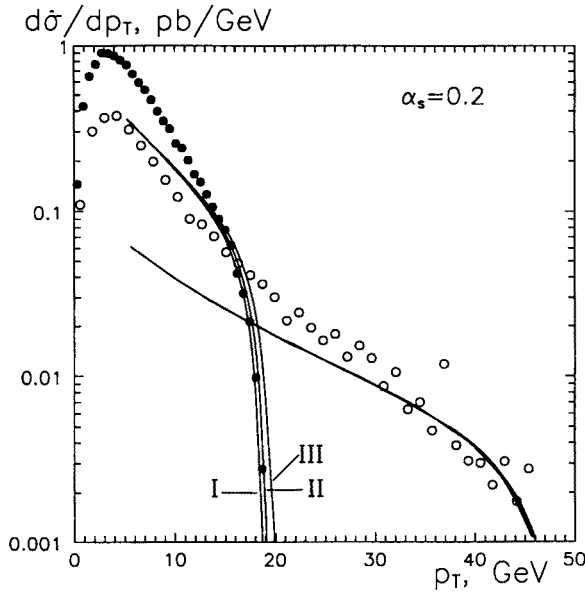


Fig. 3. Transverse momentum distributions of B_c mesons produced in $gg \rightarrow B_c b \bar{c}$ at $\sqrt{\hat{s}} = 40$ and 100 GeV: complete $O(\alpha_s^4)$ calculation (circles) and approximation Eq. (4) (solid curves). The labels I–III refer to the kinematics specified in Eq. (6).

mentation approach increase as \hat{s} decreases and that they become non-negligible at $\sqrt{\hat{s}} \leq 40$ GeV. Comparing the approximations with the $O(\alpha_s^4)$ results, one observes a slight preference for choice I or II. Most interesting, however, is the difference to $\gamma\gamma \rightarrow B_c b \bar{c}$ where heavy quark fragmentation is completely subdominant even at large p_T . Apparently, the presence of the gluon self-coupling and colour factors has a drastic influence on the relative importance of the fragmentation and recombination.

Predictions for the p_T distributions of B_c mesons in $p\bar{p}$ and pp collisions at Tevatron and LHC energies are obtained by convoluting the gg -subprocess distributions with the MRS(A') gluon structure functions [14]. The results with and without cuts in rapidity are plotted in Fig. 4 and compared with the fragmentation approximation. We have assumed the scale $Q^2 = p_T^2 + M^2$ in both $\alpha_s(Q^2)$ and the structure functions. The evolution effects in the fragmentation function $D_b(z)$ are ignored for consistency. These effects are studied in Ref. 12]. We see that after convolution the $O(\alpha_s^4)$ calculation and the fragmentation description (4) are in reasonable agreement at $p_T \geq 10$ GeV. This can be understood from the properties of the unfolded p_T -distributions illustrated in Fig. 3 and from the rise of the gluon density at small x which favours contributions from the smallest possible subenergies \hat{s} , and hence from the tails of the spectra. With decreasing p_T , the

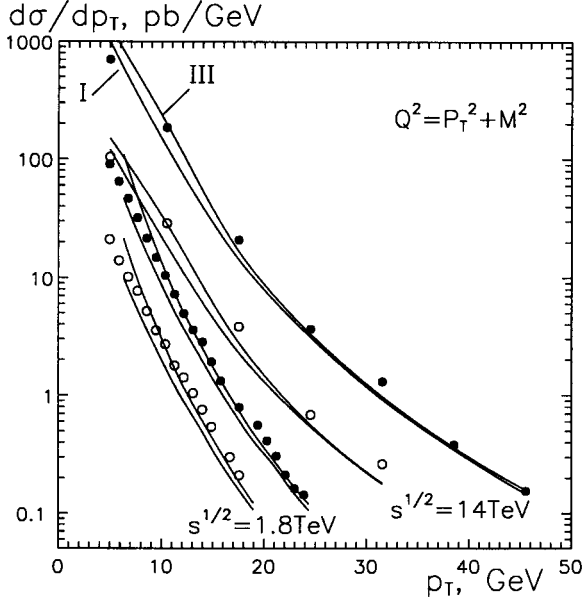


Fig. 4. Transverse momentum distribution of the B_c in $p\bar{p}(pp)$ collisions at the Tevatron (LHC) energies: complete $O(\alpha_s^4)$ calculation (circles) and fragmentation approximation Eq. (4) (solid curves). The labels I and III refer to the kinematics Eq. (6). Results are shown without a rapidity cut (full circles) and for $|y| \leq 0.5$ (empty circles).

fragmentation picture gradually breaks down and at $p_T < 5 \text{ GeV}$ only the complete $O(\alpha_s^4)$ calculation makes sense. Furthermore, the sensitivity to the kinematical prescription (6) decreases slowly with increasing p_T .

The total hadronic cross sections for B_c and B_c^* production are shown in Fig. 5. The typical rise of σ with energy is due to the rise of the gluon density as x approaches x_{\min} and the peaking of $\hat{\sigma}$ near threshold. In order to demonstrate the scale dependence of σ , we have indicated the B_c cross sections at $\sqrt{s} = 0.1, 1$ and 10 TeV for different choices of Q^2 . As one can see, at lower energies, the notorious scale ambiguity of leading logarithmic approximations leads to an uncertainty of more than one order of magnitude. Only at very high energies, the uncertainty shrinks to a factor two. It is interesting to note that the differences in the predictions resulting from different parametrizations of the gluon density, e.g. CTEQ2 [15] instead of MRS(A') [14], would be invisible in Fig. 5. Other uncertainties, connected with the decay constants $f_{B_c^{(*)}}$ and the effective quark masses m_b and m_c , amount at least to another factor of two.

The second main motivation for our own studies, besides examining the validity of the factorized fragmentation approach, has been the wish to re-

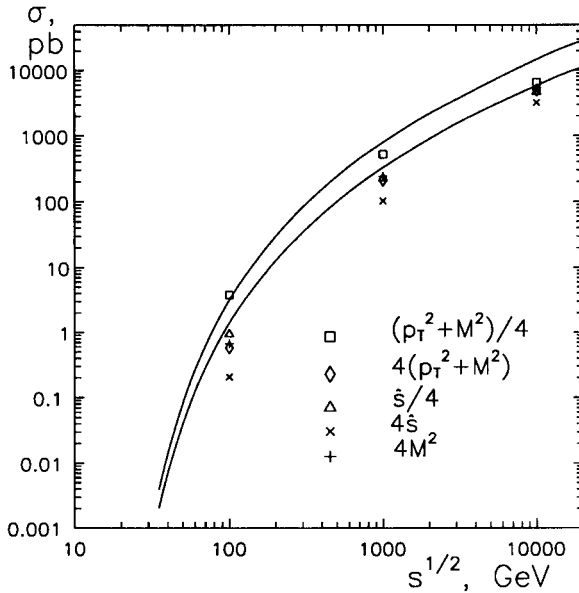


Fig. 5. Total cross sections for $pp(p\bar{p}) \rightarrow B_c^{(*)}b\bar{c} + X$ versus the c.m. energy. The lower (upper) curve show the predictions for B_c (B_c^*) and the scale $Q^2 = p_T^2 + M^2$. The symbols indicate expectations for other choices of Q^2 .

solve the confusion created by previous calculations [4–7] of hadroproduction of B_c mesons which have contradicted each other. Because of the ambiguities and uncertainties pointed out above, comparison was not always easy. We have been very careful in adjusting parameters and the gluon density to the choice in the respective calculation we considered. We also focused as much as possible on the comparison of gg cross sections and distributions being most transparent. In cases where gg results were not given, we compared the convoluted cross sections. The outcome of this comparison was surprising. As described in detail in Ref. [2], we were not able to reproduce any of the previous numerical predictions. In the meantime, the authors of Ref. [4] revised their work [16] and now agree with our results. Also the authors of [6] corrected a normalization error in [13] (this paper appeared almost simultaneously with [7]), so that the gg cross section coincide with what we found except at very high energies. However, the p_T distributions given in [13] still disagree substantially with ours.

In order to evaluate the observability of B_c mesons at the Tevatron and LHC, it is useful to integrate the p_T -distributions of Fig. 4 over $p_T \geq p_T^{\min}$. Assuming an integrated luminosity of 100 pb^{-1} , one can expect about 10^4 B_c mesons with $p_T^{\min} = 10 \text{ GeV}$ at the Tevatron, without taking into account contributions from the production and decay of B_c^* mesons and

heavier states. This rate should be sufficient for discovery in the channel $B_c \rightarrow J/\psi X$, for which a branching ratio of the order of 10% is predicted [17]. In fact, first results of such a search have already been reported in [1]. Finally, at the LHC for 100fb^{-1} , one can expect 10^7 direct B_c mesons at $p_T^{\text{min}} = 20\text{ GeV}$. This rate should then allow a more detailed study of production and decay properties.

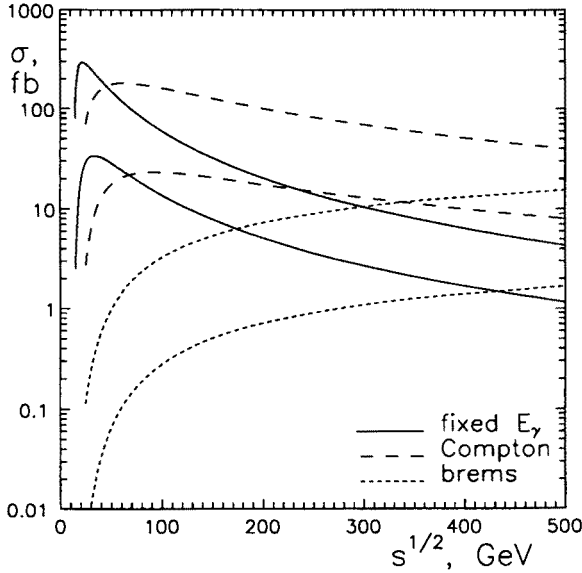


Fig. 6. Total cross sections for $\gamma\gamma \rightarrow B_c^{(*)}b\bar{c}$ versus fixed $\gamma\gamma$ c.m. energy, and after convolution with Compton and bremsstrahlung spectra, versus the e^+e^- c.m. energy. The lower (upper) curves correspond to B_c (B_c^*) production. The scale of α_s is chosen to be $p_T^2 + M^2$.

Predictions on the production rates of B_c mesons in photon-photon collisions at future e^+e^- machines are obtained by folding the total cross sections for $\gamma\gamma \rightarrow B_c^{(*)}b\bar{c}$ with the photon spectrum generated by Compton back-scattering of high intensity laser light on e^\pm beams [18], or with the Weizsäcker–Williams bremsstrahlung spectrum [19]. The expectations are illustrated in Fig. 6, where we have plotted the convoluted cross sections for B_c and B_c^* production versus the e^+e^- centre-of-mass energy, together with the unfolded cross sections as a function of fixed $\gamma\gamma$ centre-of-mass energy. Because of the long soft tail of the Compton spectrum, and the shape of the $\gamma\gamma$ cross section which peaks just above threshold, the convolution increases the cross sections substantially for energies above 100 GeV. At a 500 GeV linear collider and for an integrated luminosity of 10 fb^{-1} , one can produce about 100 B_c and 400 B_c^* . The yield of B_c mesons from

bremsstrahlung photons is invisibly small at LEP energies, but increases logarithmically with energy. In the TeV energy range bremsstrahlung photons become competitive with back-scattered laser photons in producing B_c mesons. Although, the prospect of B_c physics in $\gamma\gamma$ collisions are not very bright, observation of B_c mesons does not appear completely unfeasible.

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