

QUADRUPOLE MOMENTS OF HIGH-SPIN ISOMERS: TEST OF THE TILTED-AXIS CRANKING MODEL*

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We report the results of recent measurements of the spectroscopic quadrupole moments of high-spin isomers. For the $K^\pi = \frac{35}{2}^-$ five-quasiparticle isomer in ^{179}W we measured $Q_s = 4.00^{(+0.83)}_{(-1.06)} \text{ eb}$. It corresponds to a smaller deformation compared to the ground states of the W isotopes and is in disagreement with the current theoretical predictions. We also measured the quadrupole moment of the $I^\pi = 11^-$ isomer in ^{196}Pb , $Q_s = (-)3.41(66) \text{ eb}$. It has the same proton ($s_{1/2}^{-2}h_{9/2}i_{13/2}$) configuration as the one suggested for the $I^\pi = 16^{(-)}$ magnetic bandhead which allows to deduce the quadrupole moment of the 16^- state as $Q_s = -0.316(97) \text{ eb}$. This small value proves the near sphericity of the bandhead.

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1. Introduction

Several measurements of spectroscopic quadrupole moments of high-spin isomers were performed at the CYCLONE facility at Louvaine-la-Neuve, using the LEvel Mixing Spectroscopy technique (LEMS) [1].

We measured the quadrupole moment of the five-quasiparticle isomer in ^{179}W , addressing the question: is the nuclear deformation of the ground state and of the high-seniority multiquasiparticle excitations the same?

Next, we studied the isomers in $N = 82$ nucleus ^{196}Pb , for which several bands with enhanced M1 transitions have been established [2]. They are built on states which result from the coupling of high-spin particle to high-spin hole excitations. In the Pb nuclei, high-spin particle states are formed by proton excitations into the $\pi h_{9/2}$ and $\pi i_{13/2}$ orbitals and the high-spin hole states result from neutron excitations in the $\nu i_{13/2}$ shell. The specific coupling of the spin vectors of these two excitations causes a magnetic dipole moment with a large component perpendicular to the total spin. Its rotation around the spin axis gives rise to enhanced M1 radiation. Our study addresses the question: what is the deformation of the magnetic bandhead?

These experiments are related to the tilted-axis cranking (TAC) model, since both excitations, the high- K states and the magnetic rotation are described within this approach [3, 4]. In the former case the quenching of pairing is studied. The usual approach is to deduce the moments of inertia of the bands built on the high- K states and compare them to that of the ground-state (fully paired) rotational band [5]. Yet, the moment of inertia depends on both, the deformation and the pairing, which requires that the deformation is determined experimentally. In the case of magnetic rotation, the measured weak $B(E2)$ transition probabilities are an indication for small deformation [6]. However, the measurement of the spectroscopic quadrupole moment will provide the stringent test.

2. Experimental details

The LEMS set-up consists of a split-coil 4.4 T superconducting magnet, a target holder, allowing precise temperature control in the interval 4–600 K, and 4 Ge detectors, positioned at 0° and 90° , which monitor the target through the holes of the magnet. The magnetic field is oriented along the beam axis. The anisotropy of the γ -radiation is measured as a function of the external magnetic field. At small magnetic fields the initial anisotropy is perturbed due to the interaction of the quadrupole moment of the isomer of interest, Q_s , with the Electric-Field Gradient (EFG) of the LEMS host, V_{zz} . It is restored at high fields. The transition between the two regimes depends on the ratio of the quadrupole frequency $\nu_Q = \frac{e}{\hbar} Q_s V_{zz}$ and the magnetic moment of the isomer [1].

2.1. The case of ^{179}W

The $^{170}\text{Er}(^{13}\text{C},4n)$ reaction at 63 MeV was used to populate high-spin states in ^{179}W . The target was $500\text{ }\mu\text{g}/\text{cm}^2$ self-supporting, enriched ^{170}Er , which allowed an in-beam implantation of the ^{179}W recoils into a Tl foil at $T = 473(1)\text{ K}$, serving as a LEMS host. The obtained LEMS curve is presented in Fig. 1. A quadrupole frequency $\nu_Q = 53(8)\text{ MHz}$ was deduced. A more complete report of this experiment was published elsewhere [7].

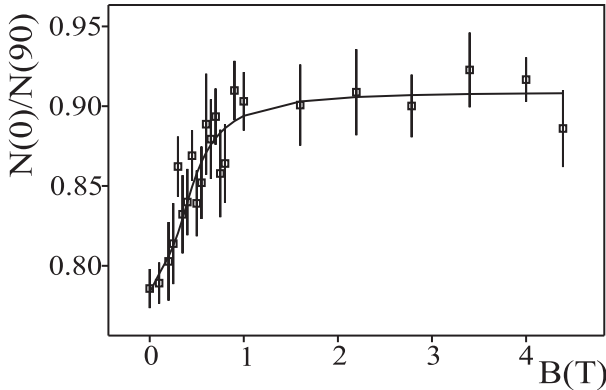


Fig. 1. LEMS curve for the $I = K = \frac{35}{2}^-$ isomer in ^{179}W . The recoiling nuclei were implanted into a Tl polycrystalline foil at a temperature of $473(1)\text{ K}$.

In order to determine the EFG of $\text{W}\underline{\text{Tl}}$ we performed band-structure calculations based on density functional theory using the full-potential Linearized Augmented Plane Wave (LAPW) method as implemented in the WIEN97 package [8]. A value of $V_{zz}(\text{W}\underline{\text{Tl}}) = 2.54 \times 10^{21}\text{ V/m}^2$ at 0 K was obtained. The EFGs in Tl decrease strongly with temperature. The temperature dependence of the EFGs in non-cubic metals follows the $T^{3/2}$ law: $V_{zz}(T) = V_{zz}(0) \times [1 - bT^{3/2}]$ [9]. In a dedicated experiment the temperature dependence factor $b = 7.6^{(+0.2)}_{(-0.4)} \times 10^{-5}\text{ K}^{-3/2}$ was derived, resulting in a value $V_{zz}(\text{W}\underline{\text{Tl}}) = 0.55^{(+0.12)}_{(-0.08)} \times 10^{21}\text{ V/m}^2$ at 473 K [10]. Thus, the spectroscopic quadrupole moment of the $K^\pi = \frac{35}{2}^-$ isomer in ^{179}W is found to be $Q_s = 4.00^{(+0.83)}_{(-1.06)}\text{ eb}$. This value is smaller than the ground state quadrupole moments of the $Z=74$ nuclei, which were derived from the reduced transition probabilities [11], and is about 2σ off the current theoretical estimates.

A critical reader might question the reliability of this result, as it relies on a calculated EFG for $\text{W}\underline{\text{Tl}}$, as well as on the assumption for a $T^{3/2}$ temperature dependence of V_{zz} . However, in a parallel experiment Ionescu-Bujor *et al.* measured the quadrupole moment of the $K^\pi = 14^+$ isomer in ^{176}W [12]. The experiment was performed under similar conditions as

our measurement (in-beam implantation of W in Tl at 464 K) and yields a quadrupole frequency $\nu_Q = 92(10)$ MHz for the isomer. This frequency corresponds to a ratio of the quadrupole moments of the two isomers:

$$\frac{Q_s(^{176}\text{W})}{Q_s(^{179}\text{W})} = 1.6 \pm 0.3, \quad (1)$$

and to value for the quadrupole moment of the 14^+ isomer, which is close to the ground-state values. This is in disagreement with the current theoretical predictions and leads to the logical question: do we understand the deformations of the high- K excitations?

2.2. The case of ^{196}Pb

High-spin isomers in ^{196}Pb (as well as in ^{194}Pb) have been populated in the $^{\text{nat}}\text{Re}(^{14}\text{N}, 5n)$ reaction. The $50\ \mu\text{m}$ $^{\text{nat}}\text{Re}$ foil served as a target and a LEMS host. Quadrupole frequencies of $\nu_Q = 38(3)$ MHz and $\nu_Q = 199(32)$ MHz were measured for the $I^\pi = 12^+$ and the $I^\pi = 11^-$ isomers, respectively. This study is discussed in more detail elsewhere [10]. The LEMS curve for the $I^\pi = 11^-$ isomer is displayed in Fig. 2. The quadrupole moment of the 12^+ isomer is known, $Q_s = 0.65(5)$ eb [13]. Also the magnetic moments of both isomers are known [13, 14]. This provides the possibility for an internal calibration and allows us to deduce the quadrupole moment of the 11^- isomer from the ratio of the measured frequencies:

$$Q_s(11^-) = \frac{\nu_Q(11^-)}{\nu_Q(12^+)} Q_s(12^+) = (-)3.41 \pm 0.66 \text{ eb}. \quad (2)$$

The negative sign comes from systematics. The quoted uncertainty includes the uncertainties of the magnetic moments as well. Assuming an axial symmetry for this state ($K = I = 11$), the quadrupole deformation is $\beta = (-)0.16(3)$, which is in a perfect agreement with the theory [15].

The $I^\pi = 11^-$ and the $I^\pi = 12^+$ isomers have the $\pi(s_{1/2}^{-2}h_{9/2}i_{13/2})$ and the $\nu(i_{13/2}^{-2})$ configurations, respectively, and the $\pi(s_{1/2}^{-2}h_{9/2}i_{13/2})_{11^-} \otimes \nu(i_{13/2}^{-2})_{12^+}$ $2p$ - $2h$ configuration is suggested for the $I^\pi = 16^{(-)}$ magnetic bandhead [2]. It is possible to derive the quadrupole moment of this state, $Q_s(16^-) = -0.32(10)$ eb, using the additivity of the quadrupole operator, $Q_2^0 = Q_2^0(\pi) + Q_2^0(\nu)$:

$$Q_s(16^-) = \langle 16, 16 | Q_2^0(\pi) + Q_2^0(\nu) | 16, 16 \rangle. \quad (3)$$

This value proves the small deformation of the $16^{(-)}$ state, which supports the concept of magnetic rotation.

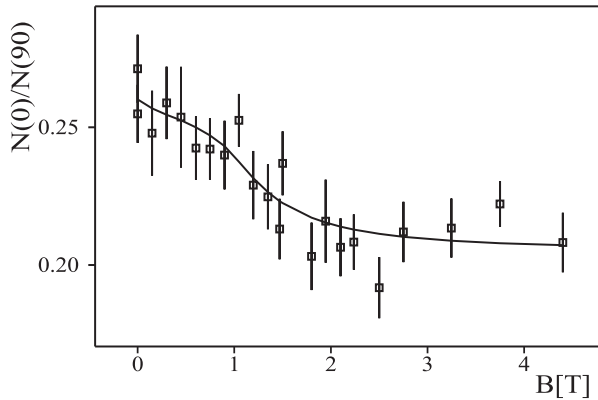


Fig. 2. LEMS curve for the $I^\pi = 11^-$ isomer in ^{196}Pb . The recoiling nuclei were implanted into a Re polycrystalline foil at room temperature.

3. Conclusions

We have measured the spectroscopic quadrupole moments of high-spin isomers in ^{179}W and ^{196}Pb using the LEMS technique. These experiments were performed to test the predictions of the TAC model [3,4]. We find that the deformation of the $K = \frac{35}{2}$ isomer in ^{179}W is smaller than the ground-state deformation, which questions the current understanding of these excitations. For the $I^\pi = 16^{(-)}$ magnetic bandhead in ^{196}Pb we derive a small quadrupole moment, which is in agreement with the concept of magnetic rotation.

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