THE RISE OF THE PROTON STRUCTURE FUNCTION F_2 TOWARDS LOW x*

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Results on the derivative of $\log(F_2)$ with respect to $\log(x)$ at fixed Q^2 are presented. The measured derivatives are within errors independent of x for $Q^2 \geqq 0.85\,\mathrm{GeV}^2$ and increase linearly with $\log(Q^2)$ for $10^{-4} \leqq x \leqq 0.01$ and $Q^2 \gtrapprox 3\,\mathrm{GeV}^2$. The results are based on preliminary and published H1 data which at Q^2 below $2\,\mathrm{GeV}^2$ are combined with NMC and ZEUS data.

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1. Introduction

The rise of the proton structure function F_2 towards small Bjorken x has been discussed since the existence of QCD. In the double asymptotic limit (large energies, *i.e.* small x, and large photon virtualities Q^2) the DGLAP evolution equations [1] can be solved [2] and F_2 is expected to rise approximately like a power of x towards low x. A power like behaviour is also expected in the BFKL approach [3]. However, it soon was discussed [4] that this rise may eventually be limited by gluon self interactions in the nucleon, or more generally due to unitarity constraints.

Experimentally this rise towards small x was first observed in 1993 in the HERA data [5]. Meanwhile the precision of the F_2 data is much improved and the rise can be studied in great detail.

2. Procedure

The low x behaviour of F_2 at fixed Q^2 is studied locally by the measurement of the derivative $\lambda \equiv -(\partial \ln F_2/\partial \ln x)_{Q^2}$ as function of x and Q^2. The results are based on preliminary H1 F_2 data presented to this conference [6]

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covering the range $0.5 < Q^2 < 3.5\,\mathrm{GeV}^2$ and published H1 data [7,8] which cover the range $1.5 < Q^2 < 150\,\mathrm{GeV}^2$. The low Q^2 F_2 data were obtained by shifting the *ep* interaction vertex by 70 cm in proton beam direction [6]. At $Q^2 < 2\,\mathrm{GeV}^2$ the H1 data are also shown combined with data of NMC [9] and ZEUS [10]. The derivative $\lambda(x, Q^2)$ is evaluated using data points at adjacent values of x at fixed Q^2 taking into account error correlations and x spacing corrections. The derivatives are compared with the next to leading order (NLO) QCD fit to the H1 cross section data [7] and a "fractal" fit [11] where self-similar properties of the proton structure are assumed.

3. Results

The x and Q^2 dependence of $\lambda = -(\partial \ln F_2/\partial \ln x)_{Q^2}$ is shown in Fig. 1.

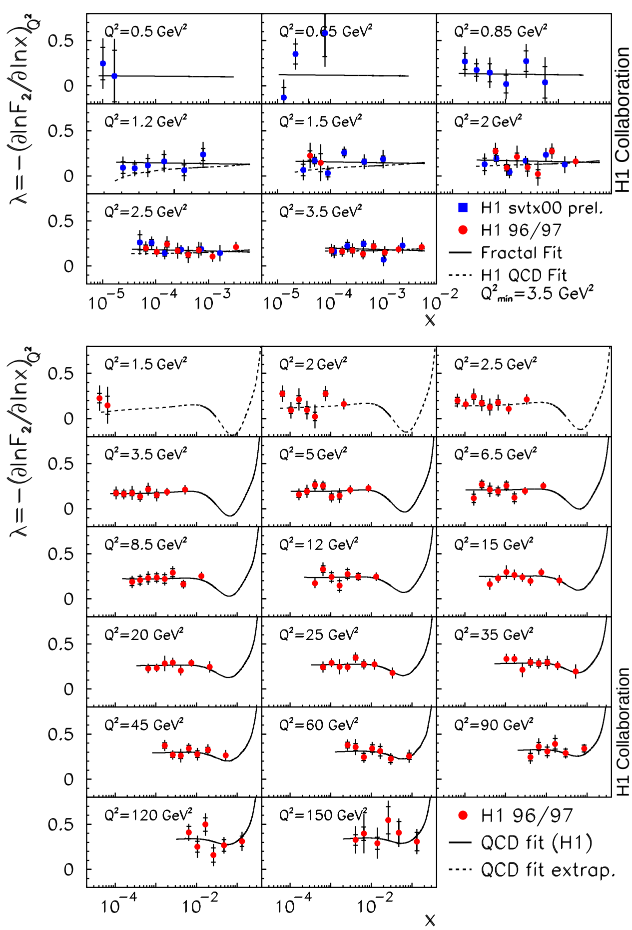

Fig. 1. Derivative λ compared with the QCD analysis of Ref. [7] and a "fractal" fit [11] for $0.5 < Q^2 < 3.5\,\mathrm{GeV}^2$ (upper plot) and for $1.5 < Q^2 < 150\,\mathrm{GeV}^2$ (lower plot).

The new shifted vertex and the published data agree well in the overlap region. The derivative λ is constant within experimental uncertainties for fixed Q^2 in the range $x < 0.01$, implying that the data are consistent with the power behaviour $F_2 = c(Q^2)\,x^{-\lambda(Q^2)}$. Fitting this form for each Q^2 bin to the data at $x < 0.01$, results in the λ and c values presented in Fig. 2.

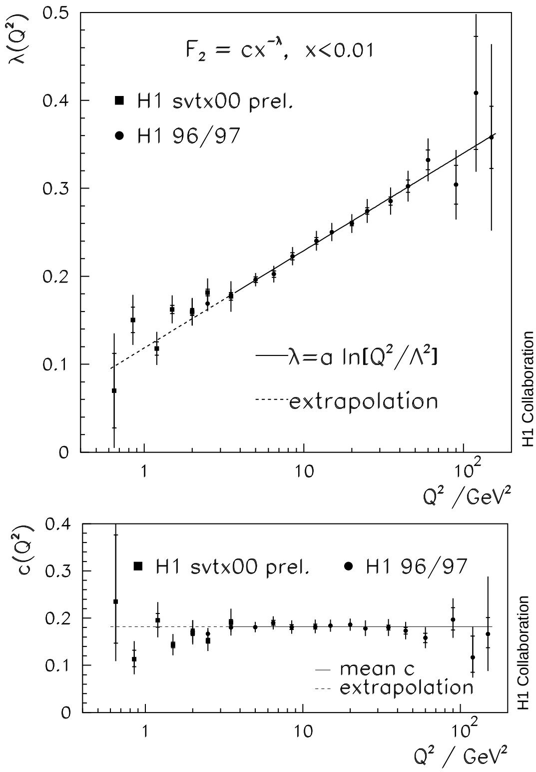

Fig. 2. $\lambda(Q^2)$ and $c(Q^2)$ from fits of the form $F_2 = c(Q^2)\,x^{-\lambda(Q^2)}$ to the H1 structure function data [7] and [11].

The results show that the F_2 data at low x for $Q^2 \gtrapprox 3.5\,\mathrm{GeV}^2$ can be well described by the very simple parameterisation

$$F_2 = c\,x^{-\lambda(Q^2)}\,, \qquad \text{where} \qquad \lambda(Q^2) = a\,\ln\left[\frac{Q^2}{\Lambda^2}\right] \tag{1}$$

with $a = 0.0481 \pm 0.0013 \pm 0.0037$ and $\Lambda = 292 \pm 20 \pm 51$ MeV and $c \approx 0.18$.

At low Q^2 the deviation of λ from the logarithmic Q^2 dependence and the decrease of $c(Q^2)$ is more significant if the H1 data are combined with NMC [9] and ZEUS [10] data (see Fig. 3).

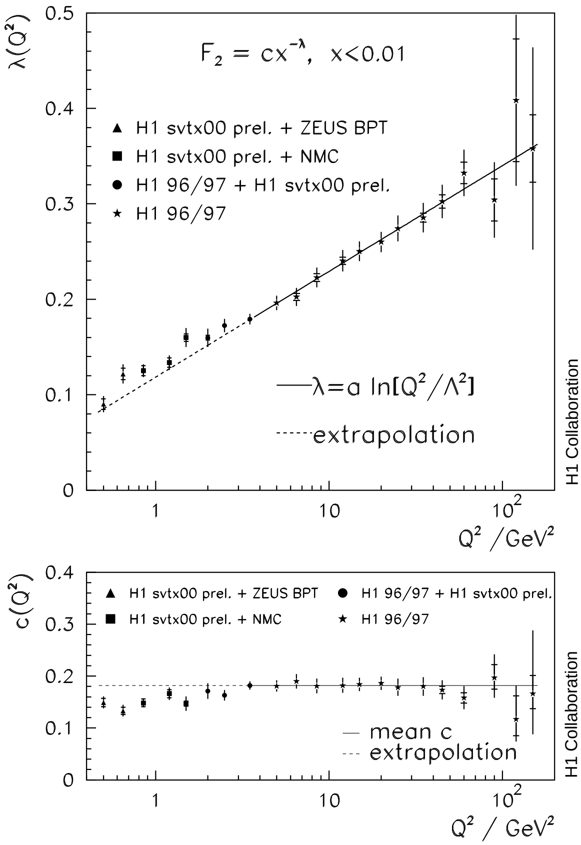

Fig. 3. $\lambda(Q^2)$ and $c(Q^2)$ from fits of the form $F_2 = c(Q^2)\,x^{-\lambda(Q^2)}$ combining the H1 structure function data of [7] and [11] and the H1 data with data of NMC [9] and ZEUS [10].

The deviations from a simple constant respectively logarithmic behaviour occur at about such Q^2 values below which perturbative QCD fits (*e.g.* [7]) are not supposed to be valid. At small Q^2 the structure function F_2 can be related to the total virtual photon absorption cross section by

$$\sigma_{\rm tot}^{\gamma^* p} = 4\pi\alpha^2 \frac{F_2}{Q^2} \sim \frac{x^{-\lambda}}{Q^2}\,, \tag{2}$$

where the total $\gamma^* p$ energy squared is given by $s = Q^2/x$. For $Q^2 \to 0$ we can expect $c(Q^2) \to 0$ and $\lambda(Q^2) \to\approx 0.08$. The latter value corresponds to the energy dependence of soft hadronic interactions $\sigma_{\rm tot} \sim s^{\alpha_{I\!P}(0)-1}$ with $\alpha_{I\!P}(0) - 1 \approx 0.08$ [12] which is approximately reached at $Q^2 = 0.5\,{\rm GeV}^2$.

4. Conclusion

No significant deviation from the power behaviour $F_2 \sim x^{-\lambda}$ at fixed Q^2 is visible at present energies and $Q^2 \gtrsim 0.85\,\mathrm{GeV}^2$. More specifically

- For $x < 0.01$ the derivative $\lambda \equiv -(\partial \ln F_2/\partial \ln x)_{Q^2}$ is independent of x within errors.
- λ is proportional to $\ln(Q^2)$ for $Q^2 \gtrsim 3\,\mathrm{GeV}^2$, *i.e.* in the pQCD region.
- Here the data can be very simply parametrised by $F_2 = cx^{-\lambda(Q^2)}$.
- At $Q^2 \lesssim 3\,\mathrm{GeV}^2$ deviations from the logarithmic Q^2 dependence of λ are observed.
- At low Q^2 ($Q^2 \lesssim 1\,\mathrm{GeV}^2$) the energy rise is similar as in soft hadronic interactions.

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