MAGNETIC PROPERTIES OF CeRhIn₅ UNDER PRESSURE PROBED BY ¹¹⁵In-NQR*

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We report the pressure (P)- induced evolution of the antiferromagnetism in CeRhIn₅ which undergoes a superconducting transition at $T_c \sim 2.1$ K at pressurs exceeding $P_c \sim 1.6$ GPa. From measurements of ¹¹⁵In nuclear-spin-lattice-relaxation time (T_1) under P, we found that Néel temperature T_N is reduced above P = 1.23 GPa, which is accompanied by an emergent *pseudogap* behavior.

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1. Introduction

Since the discovery of cerium (Ce) based heavy-fermion (HF) superconductivity in CeCu₂Si₂ [1], many experimental works on CeCu₂Ge₂, [2] CeIn₃ [3–5] and CePd₂Si₂ [6] have suggested that antiferromagnetism and superconductivity are closely related each other. The discovery of pressure (P)-induced HF superconductivity in Ce-based antiferromagnetic (AF) compounds has stimulated further experimental works under P [7–10].

A HF AF CeRhIn₅ undergoes a P induced superconducting (SC) transition at a lower critical $P \sim 1.6$ GPa, yet reaching a $T_{\rm c} \sim 2$ K [7] higher than previous examples. Systematic studies on CeRhIn₅ seem to allow us

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to unravel an interplay between antiferromagnetism and superconductivity under P. CeRhIn₅ forms in the tetragonal HoCoGa₅ structure that is composed of alternating layers of CeIn₃ and RhIn₂ stacked sequentially along the [001] direction [7]. It orders antiferromagnetically with an incommensurate wave vector (1/2, 1/2, 0.297) [11]. The Néel temperature $T_{\rm N}=3.8~{\rm K}$ at P = 0 slightly increases up to $P \sim 1.0$ GPa, followed by an onset of superconductivity at $T_{\rm c} = 2.1$ K at pressures exceeding a critical pressure $P_{\rm c} \sim 1.63$ GPa [7]. In the previous paper [10], the ¹¹⁵In-NQR study of $CeRhIn_5$ has clarified the *P*-induced anomalous magnetism and unconventional superconductivity. In the AF region, $T_{\rm N}$ which is evidenced from clear shift and splits of NQR spectrum exhibits a moderate increase up to P = 1.0 GPa, but decreases above P = 1.23 GPa. By contrast, the internal field H_{int} at the In(1) site, that is associated with the magnetic ordering, is almost linearly reduced against increasing P. Note that the In(1) site is equivalent to the In site in the CeIn₃. In the SC region at P = 2.1 GPa, it was reported that the $1/T_1$ reveals a T^3 dependence without the coherence peak just below $T_{\rm c}$, consistent with the existence of line-node gap [10]. It is, however, not yet clear in $CeRhIn_5$ how antiferromagnetism evolves into superconductivity with increasing P. Here, we focus on magnetic properties of CeRhIn₅ near $P_{\rm c} \sim 1.6$ GPa probed by ¹¹⁵In-NQR T_1 measurements in P = 0 - 1.6 GPa.

2. Experimental result

The ¹¹⁵In (I = 9/2)-NQR spectrum at the paramagnetic state consists of four transitions that are equally spaced with a quadrupole frequency $n \times \nu_{\rm Q}$ with n = 1, 2, 3 and 4. $1/T_1$ was measured at the transitions of 2 $\nu_{\rm Q}$ $(\pm 3/2 \leftrightarrow \pm 5/2)$ and 3 $\nu_{\rm Q}$ $(\pm 5/2 \leftrightarrow \pm 7/2)$ at the In(1) site, using the conventional saturation-recovery method in the T = 1, 4 - 100 K at P = 0.46,1.23, and 1.6 GPa. Hydrostatic pressure was applied by utilizing a BeCu piston-cylinder cell, filled with Daphne oil (7373) as a pressure-transmitting medium. To calibrate a value of pressure at low temperatures, the shift in $T_{\rm c}$ of Sn metal under P was monitored by resistivity measurement.

The *T* dependence of $1/T_1T$ below 10 K is shown in Fig. 1 at P = 0.46, 1.23, and 1.6 GPa. At the respective values of *P*, it is clearly seen that AF order occurs at $T_{\rm N}=4.0$, 3.6 and 2.8 K, as evidenced by the clear peak in $1/T_1T$ due to critical magnetic fluctuations toward the AF ordering. This indicates that $T_{\rm N}$ reduces above P = 1.23 GPa as *P* approaches $P_{\rm c} \sim 1.63$ GPa. At P = 1.23 and 1.6 GPa, $1/T_1T$ shows a broad peak around $T_{\rm PG}$ well above $T_{\rm N}$. This resembles the pseudogap behavior found in high- $T_{\rm c}$ copper oxide superconductors [12]. Likewise, when *P* approaches the critical pressures $P_{\rm c}$, the low-energy spectral weight of magnetic fluctuations is sup-

pressed before an ordering occurs. We note that the pseudogap behavior has been found in either two- or lower-dimensional strongly correlated electron systems [12]. Very recently, in CeRhIn₅, anisotropic three dimensional AF fluctuation was reported from neutron scattering at P = 0 with an energy scale of less than 1.7 meV at temperatures as high as 3 T_c [13]. On the other hand, as P further increases up to P = 2.1 GPa where the SC transition appears, $1/T_1T$ continues to increase down to $T_c = 2.2$ K without any signature for the pseudogap behavior as seen in the previous report [10]. The Tvariation of $1/T_1T$ is consistent with the three dimensional AF Fermi-liquid model described by the self-consistent renormalized (SCR) theory for nearly AF metals [10, 14]. Thus the P-induced evolution in the magnetic fluctuations, from a magnetic regime of reduced dimensionality to a more isotropic one, may take place in a narrow P window of 1.6–2.1 GPa, when the AF order evolves into the SC one. The pressure-temperature phase diagram for CeRhIn₅ is summarized in Fig. 2.



Fig. 1. T dependence of 1 / T_1 at P = 0.46, 1.23 and 1.6 GPa. The solid and dotted arrows indicate T_{PG} and T_N .

In conclusion, we have reported the *P*-induced evolution of the magnetic properties in CeRhIn₅ on the basis of the ¹¹⁵In-NQR T_1 measurements. In the itinerant antiferromagnet CeRhIn₅ at P = 0, T_N slightly increases up to P = 1.0 GPa, but it turns to decrease at P = 1.23 GPa when approaching $P_c \sim 1.63$ GPa at which superconductivity sets in. This reduction in T_N coincides with the emergence of the pseudogap behavior that is evidenced by the suppression in $1/T_1T$ above T_N . This *P*-induced evolution in the magnetic fluctuations, from a magnetic regime of reduced dimensionality to a more isotropic one, may take place in a narrow *P* window of 1.6–2.1 GPa when the AF order evolves into the SC one.



Fig. 2. P-T phase diagram for CeRhIn₅. The open marks are determined from ac-susceptibility measurement [10].

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