# HIGH RESOLUTION PES INVESTIGATIONS ON THE PROTOTYPE HEAVY FERMION COMPOUND $CeCu_6^*$

D. EHM<sup>a</sup>, F. REINERT<sup>a</sup>, J. KROHA<sup>b</sup>, O. STOCKERT<sup>c</sup> AND S. HÜFNER<sup>a</sup>

 <sup>a</sup>Universität des Saarlandes, F.R. 7.2 Experimentalphysik ImStadtwald Bau 22 66123 Saarbrücken, Germany
<sup>b</sup>Institut für Theorie der Kondensierten Materie, Universität Karlsruhe Postfach 6980, 76128 Karlsruhe, Germany
<sup>c</sup>Max-Planck Institute for Chemical Physics of Solids Nöthnitzer Str. 40 01187 Dresden, Germany

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We present high resolution photoemission investigations on the heavy fermion compound CeCu<sub>6</sub>. By application of a normalization procedure we can resolve the Kondo resonance and its crystal field structures in both the He II<sub> $\alpha$ </sub> and He I<sub> $\alpha$ </sub> data. A comparison with NCA calculations allows a quantitative determination of the Kondo temperature and the crystal field energies.

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## 1. Introduction

Since more than two decades heavy fermion (HF) compounds are under intensive investigation by photoemission spectroscopy (PES). This method allows to study directly core-level spectra, the valence states and the occupied part of the 4f spectral function close to the Fermi level ( $E_{\rm F}$ ). In contrast, inverse PES has access to the spectral function above  $E_{\rm F}$ , but due to its finite energy resolution — is not able to resolve in detail the 4ffeatures, *i.e.* the Kondo resonance (KR) and its crystal field (CF) and spinorbit (SO) partners close to the Fermi level. Recently it has been shown [1] that high resolution PES can resolve the KR in  $\gamma$ -like CeCu<sub>2</sub>Si<sub>2</sub>. This direct observation of the KR was possible by application of a well known normalization procedure based on the division by the Fermi-Dirac distribution (FDD) [3], that allows to recover the thermally occupied spectrum

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up to ~  $5k_{\rm B}T$  above  $E_{\rm F}$ . At higher temperatures this energy range even includes the CF structures above  $E_{\rm F}$ . To determine  $T_{\rm K}$  and the CF splittings of the J = 5/2 multiplet ( $\Delta_{\rm CF}$ ) one can directly compare the PES 4f spectrum with the spectral function extracted from calculations in the non crossing approximation (NCA) based on the Single Impurity Anderson Model (SIAM) [2]. By an iterative fitting of the data one is able to determine all relevant model parameters, from which  $T_{\rm K}$  and  $\Delta_{\rm CF}$  can be extracted and systematically compared to numbers determined e.g. by inelastic neutron scattering (INS). Surprisingly a reasonable agreement can be found between PES and INS results, although PES investigates the surface region of the sample, whereas INS reflects bulk properties.

Here we turn our attention to the prototype heavy fermion compound  $CeCu_6$ , which is characterized by a huge value of the Sommerfeld coefficient ( $\gamma_0 = 1.6 \text{ JK}^{-2}\text{mole}^{-1}$ ) and a very low  $T_{\rm K}$  of about 5 K. There exist only few PES studies in the literature of this compound [4–6] because of its low  $T_{\rm K}$ value and the corresponding weak 4f spectral features in the PES data. We present temperature dependent high resolution ( $\Delta E = 5.4 \text{ meV}$ ) He II<sub> $\alpha$ </sub> and He I<sub> $\alpha$ </sub> PES spectra, which clearly show the 4f spectral features. Contrary to previous investigations [7], we find that the 4f PES cross-sections at He I<sub> $\alpha$ </sub> are still strong enough to observe the KR and its SO structure.

### 2. Experimental setup

The PES experiments have been performed using a SCIENTA SES 200 analyzer in combination with a monochromatized GAMMADATA VUV lamp using photon energies of  $h\nu = 21.2$  eV (He I<sub> $\alpha$ </sub>) and  $h\nu = 40.8$  eV (He II<sub> $\alpha$ </sub>) [8]. The energy resolution of the system was set to 5.4 meV to have sufficient photoemission intensity during the He II<sub> $\alpha$ </sub> measurements. The single-crystalline CeCu<sub>6</sub> samples were grown by Czochralski technique in a high-purity argon atmosphere using a tungsten crucible.

#### 3. Results

Typical high resolution near- $E_{\rm F}$  spectra on CeCu<sub>6</sub> are displayed in Fig. 1, measured at different temperatures with He I<sub> $\alpha$ </sub> (right panel) and He II<sub> $\alpha$ </sub> (left panel) radiation. It should be mentioned that although the photoionization cross-sections of the Ce 4f states increase about a factor 3.2 in going from He I<sub> $\alpha$ </sub> to He II<sub> $\alpha$ </sub> radiation [9] the PES spectra over an extended energy range (not shown here) clearly show the 4f derived spectral features, namely the SO structure at around  $E_B = 250$  meV and the tail of the KR directly at  $E_{\rm F}$ . The latter can be observed with He II<sub> $\alpha$ </sub> and even with He I<sub> $\alpha$ </sub> radiation as shown in the insets of Fig. 1, but with different intensities. The spectral



Fig. 1. Theoretical and experimental spectra of CeCu<sub>6</sub> after application of the mentioned normalization procedure. The PES data are taken with He II<sub> $\alpha$ </sub> (left) and He I<sub> $\alpha$ </sub> radiation (right) at different temperatures. Model parameters of the NCA:  $\epsilon_f = -1.05 \text{ eV}, D = 2.8 \text{ eV}, \Delta_{\rm CF} = 7.2/13.9 \text{ meV}, \Delta_{\rm SO} = 250 \text{ meV}, V = 116 \text{ meV}.$  The insets show the data on the same energy scale prior to the normalization. All spectra are normalized to the same intensity at  $\approx 100 \text{ meV}.$ 

intensity of the raw He I<sub> $\alpha$ </sub> and He II<sub> $\alpha$ </sub> spectra at energies of a few meV above  $E_{\rm F}$  and higher is suppressed by the FDD (see insets of Fig. 1). If we now apply the normalization method (division by the FDD) to the raw data, we can reconstruct the spectral information up to an energy of ~  $5k_{\rm B}T$  (see Fig. 1). Obviously there appears a narrow peak with a full width at half maximum (FWHM) of ~ 6 meV and a maximum at about 3 meV above  $E_{\rm F}$  both in the low temperature He I<sub> $\alpha$ </sub> and He II<sub> $\alpha$ </sub> data. This is sufficiently below  $5k_{\rm B}T \approx 7$  meV, which defines the limit of the usable energy range.

In order to analyse this narrow peak we have performed NCA calculations with a set of parameters that is given in the figure caption of Fig. 1. The striking coincidence between theory and experiment is obvious at the first glance. The NCA spectra are able to describe the lineshape and the energy position of the narrow peak just above  $E_{\rm F}$  — the Kondo resonance — both in the spectra taken with He I<sub> $\alpha$ </sub> and He II<sub> $\alpha$ </sub> radiation. In addition, the temperature dependence of the spectra is described consistently. Towards higher T, the linewidth of the Kondo resonance increases, while the maximum intensity becomes smaller and the KR and the CF structures are successively smeared out. Above T = 101 K the thermal broadening of the FDD is large enough to have access to the CF excitations at ~ 15 meV. Compared to the He  $I_{\alpha}$  spectra, a considerable increase of the relative intensity of the KR can be observed in the data taken with He II<sub> $\alpha$ </sub> radiation. If we compare both normalized low temperature spectra we find an intensity gain by a factor of 2 in going from He I<sub> $\alpha$ </sub> to He II<sub> $\alpha$ </sub> radiation. This is identical with the observed intensity loss in the SO feature by changing from He II<sub> $\alpha$ </sub> to He I<sub> $\alpha$ </sub>. The difference between this value and the theoretical one, can be explained by the influence of the Cu 3d, Cu 4s and Cu 4p, as well to the Ce 5d states, that partly give a considerable contribution to the spectra in this energy range.

In order to estimate  $T_{\rm K}$  from our data, we analysed the linewidth of the KR in the NCA spectra calculated at  $T = 0.1 \times T_{\rm K} \approx 1$  K (low temperature limit) with the parameter set extracted from the fit described above [1]. By this procedure we could determine  $T_{\rm K} = 4.6$  K for CeCu<sub>6</sub>. We are now able to compare this Kondo temperature and the CF energies with the values determined by INS measurements. From INS we find  $T_{\rm K} = 5.0 \pm 0.5$  K [10] and  $\Delta_{\rm CF} = 7.0/13.8$  meV [11]. Surprisingly our results are in reasonable agreement with the INS bulk values, considering especially that in the surface sensitive PES the spectral function, and hence  $T_{\rm K}$ , might be modified compared to the bulk values. The reason for this can be found in the fact that CeCu<sub>6</sub> belongs to the class of  $\gamma$ -like Ce compounds, where no real difference exists between surface and bulk electronic properties [12]. Thus we have shown that PES studies can contribute important information to the discussion of CF excitations in Ce compounds.

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## REFERENCES

- [1] F. Reinert et al., Phys. Rev. Lett. 87, 106401 (2001).
- [2] N.E. Bickers et al., Phys. Rev. **B36**, 2036 (1987).
- [3] T. Greber et al., Phys. Rev. Lett. 79, 4465 (1997).
- [4] F. Patthey et al., Phys. Rev. B34, 2967 (1986).
- [5] L. Schlapbach et al., Phys. Rev. C19, L63 (1986).
- [6] G. Chiaia et al., Phys. Rev. **B55**, 9207 (1997).
- [7] M. Garnier et al., Phys. Rev. Lett. 78, 4127 (1997).
- [8] F. Reinert et al., Phys. Rev. Lett. 85(18), 3930 (2000).
- [9] J. Yeh, I. Lindau, At. Data Nucl. Tables 32, 1 (1985).
- [10] J. Rossat-Mignod et al., J. Magn. Magn. Mater. 76&77, 376 (1988).
- [11] E.A. Goremychkin et al., Phys. Rev. **B47**, 14580 (1993).
- [12] C. Laubschat et al., Phys. Rev. Lett. 65, 1639 (1990).