

CONTINUUM COUPLING EFFECTS IN SPECTRA
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Continuum-coupling correction to binding energies in neutron rich oxygen and fluorine isotopes and to excitation energies of 2_1^+ states in ^{36}Ca and ^{36}S mirror nuclei are studied using the real-energy continuum shell model.

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1. Introduction

Atomic nuclei form a network of coupled systems communicating with each other through decays and captures [1]. If continuum states are neglected, this communication is broken and each system becomes an isolated closed quantum system (CQS). It is obvious, that the CQS description of atomic nuclei (*e.g.* the nuclear shell model (SM)) becomes self-contradictory for weakly-bound or unbound states.

A classic example of a continuum coupling is the Thomas–Ehrmann shift [2] which manifests itself in the asymmetry in the energy spectra between mirror nuclei having different particle emission thresholds. A consistent description of the interplay between scattering states, resonances, and bound state requires an open quantum system (OQS) formulation. Comprehensive many-body theory of weakly bound/unbound states has been advanced recently in the time-asymmetric quantum mechanics using the complete ensemble of single-particle states consisting of resonant (Gamow) states and the complex-energy, non-resonant continuum of scattering states from which the complete many-body basis of OQs can be obtained [3]. Another formulation of the continuum shell model is obtained by embedding

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standard SM in the continuum of decay channels. This approach provides a unified description of nuclear structure and nuclear reaction aspects [4–7].

In this paper, we study the effect of the continuum coupling on spectra of ^{36}Ca and ^{36}S mirror nuclei. We show that the continuum coupling explains naturally not only the appearance of the asymmetry in spectra but also provides a major part of the Thomas–Ehrmann shift. We shall also discuss salient effects of a continuum coupling in the binding energy systematics, in particular the anti-odd–even staggering (anti-OES) and the effective range of energies in which various decay channels are correlated with each other with the discrete many-body states changing significantly their energy and wave function.

2. Shell model embedded in the continuum

In the shell model embedded in the continuum (SMEC), nucleus is described as an OQS [8]. The total function space consists of the set of \mathcal{L}^2 -functions, as in the standard SM, and the set of scattering states (decay channels). These two sets are obtained by solving the Schrödinger equation for discrete states of the closed subsystem (closed quantum system (CQS)): $H_{\text{SM}}\Phi_i = E_i^{(\text{SM})}\Phi_i$, and for scattering states of the external environment: $\sum_{c'}(E - H_{cc'})\xi_E^{c'(+)} = 0$, where H_{SM} is the SM Hamiltonian, and $H_{cc'} = H_0 + V_{cc'}$ is the coupled-channel (CC) Hamiltonian. Channels: $c \equiv [J_i^{A-1}; (l, j)]^{J_k^A}$, are determined by the motion of an unbound particle with orbital angular momentum l and total angular momentum j relative to the residual nucleus with $A - 1$ bound particles in a SM state Φ_j^{A-1} . $\xi_E^{c(+)}$ are channel projected scattering states with outgoing asymptotics. States of the daughter nucleus are assumed to be stable with respect to the particle emission. By means of two functions sets: $\mathcal{Q} \equiv \{\Phi_i^A\}$, $\mathcal{P} \equiv \{\zeta_E^{c(+)}\}$, one can define the corresponding projection operators and the projected Hamiltonians: $\hat{Q}H\hat{Q} \equiv H_{QQ}$ and $\hat{P}H\hat{P} \equiv H_{PP}$ [4, 8]. H_{SM} is identified with the CQS Hamiltonian H_{QQ} and H_{cc} with H_{PP} . The coupling term H_{PQ} is given by the two-body residual interaction [8].

Schrödinger equation in the total function space splits into the two equations for projected operators H_{QQ} , H_{PP} . Assuming $\mathcal{Q} + \mathcal{P} = I_d$, one can determine a third wave function $\omega_i^{(+)}$ which is a continuation of SM wave function Φ_i^A in the scattering continuum. $\omega_i^{(+)}$ is obtained by solving the CC equations with the source term [8]. Using the three function sets: $\{\Phi_i^A\}$, $\{\zeta_E^{c(+)}\}$ and $\{\omega_i^{(+)}\}$, one obtains a solution in the total function space:

$$\Psi_E^c = \zeta_E^c + \sum_{i,k} (\Phi_i^A + \omega_i^{(+)}(E)) \langle \Phi_i^A | (E - \mathcal{H}_{QQ}(E))^{-1} | \Phi_k^A \rangle \langle \Phi_k^A | H_{QP} | \zeta_E^c \rangle, \quad (1)$$

where E is the total energy and $\mathcal{H}_{QQ}(E)$ is the energy-dependent effective Hamiltonian in \mathcal{Q} subspace:

$$\mathcal{H}_{QQ}(E) = H_{QQ} + H_{QP}G_P^{(+)}(E)H_{PQ}, \quad (2)$$

where $G_P^{(+)}(E)$ in (2) is the Green function in \mathcal{P} . \mathcal{H}_{QQ} takes into account a modification of the CQS Hamiltonian (H_{QQ}) by couplings to the environment of decay channels. \mathcal{H}_{QQ} is a complex-symmetric matrix above the particle-emission threshold $E^{(\text{thr})}$ and Hermitian below it. Diagonalisation of \mathcal{H}_{QQ} by an orthogonal and, in general, non-unitary transformation yields complex eigenvalues $\tilde{E}_i - \frac{1}{2}i\tilde{\Gamma}_i$, which depend on the energy E of the particle in the continuum. Energies and widths of the resonance states follow from: $E_i = \tilde{E}_i(E = E_i)$, $\Gamma_i = \tilde{\Gamma}_i(E = E_i)$, where $\tilde{E}_i(E)$ and $\tilde{\Gamma}_i(E)$ are the eigenvalues of $\mathcal{H}_{QQ}(E)$. Details of the SMEC calculations can be found in Refs. [4, 8].

3. Features of the continuum-coupling energy correction to eigenvalues of the closed quantum system

In this chapter, we shall discuss salient features of the continuum-coupling energy correction: $E_i^{(\text{corr})}(E) = \langle \Phi_i | H_{QP}G_P^{(+)}(E)H_{PQ} | \Phi_i \rangle$, for the ground state (g.s.) of neutron-rich oxygen and fluorine isotopes. Details of the SMEC calculations for continuum-coupling energy correction to binding energies can be found in Ref. [6]. In these calculations, all possible couplings of the g.s. of A nucleus to the states of $A - 1$ nucleus are taken into account incoherently.

3.1. Binding systematics in neutron-rich nuclei

For $T = 1$ couplings, the average behaviour of $E_{g.s.}^{(\text{corr})}$ is determined by the strong dependence on $E_n^{(\text{thr})}$ (c.f. Fig. 1a of Ref. [6] for oxygen isotopes)¹. Close to the neutron drip line, this leads to an effective enhancement of nn continuum-coupling strength which cannot be compensated by the $E_n^{(\text{thr})}$ -independent correction of monopole terms. On the top of this behavior, one can see the odd-even staggering (OES) of $E_{g.s.}^{(\text{corr})}(N)$. Blocking of the virtual scattering to continuum states by an unpaired nucleon diminishes the nn continuum-coupling energy correction in odd- N nuclei. For a fixed $E_n^{(\text{thr})}$, the coupling of $[N = 2k, Z]$ system to $[N = 2k - 1, Z] \otimes n$ decay channels is enhanced and the coupling of $[N = 2k + 1, Z]$ system to $[N = 2k, Z] \otimes n$ decay channels becomes weaker with respect to an averaged behavior.

¹ In the chosen *sdfp* model space, *np* continuum couplings are absent in oxygen isotopes.

This drip-line effect is seen only in a narrow range of excitation energies around $E_n^{(\text{thr})}(N) \simeq 0$ and vanishes for $E_n^{(\text{thr})}(N) \geq 4$ MeV. For more realistic values of separation energies, as given by nuclear SM [6], $E_n^{(\text{thr})}$ exhibits the pairing induced OES which becomes partially compensated close to a drip line by the anti-OES effect induced by couplings to decay channels, both opened and closed.

In fluorine isotopes (*c.f.* Fig. 2a of Ref. [6]), $E_{\text{g.s.}}^{(\text{corr})}$ is dominated by np continuum couplings ($T = 0, 1$). The strength of np continuum coupling can be deduced by comparing SMEC results for binding energies with experimental data [6]. An optimal value of $V_0^{(np)}/V_0^{(nn)}$ close to the neutron drip line is $V_0^{(np)} \simeq (1/2)V_0^{(nn)}$, whereas for nuclei close to the valley of beta-stability a standard choice is: $V_0^{(np)} \simeq 2V_0^{(nn)}$. Since the np couplings provide a major part of this correction, therefore the gradual reduction of $V_0^{(np)}/V_0^{(nn)}$ toward the neutron drip line leads to a non-linear dependence of the $E_{\text{g.s.}}^{(\text{corr})}$ and, hence, to the dependence of two-body monopole terms on the neutron number. A similar dependence is expected if effective three-body interactions are included in the two-body framework of the SM [9].

The np continuum coupling in odd-odd ($[N = 2k+1, Z = 2m+1]$) fluorine isotopes is increased as compared to the neighboring even-odd ($[N = 2k+2, Z = 2m+1]$ and $[N = 2k, Z = 2m+1]$) nuclei. Contrary to the continuum-coupling correction for like particles (nn or pp couplings), the characteristic anti-OES of $E_{\text{g.s.}}^{(\text{corr})}$ persists even for $E_n^{(\text{thr})} > 4$ MeV. This effect is further enhanced by the OES of one-neutron (1n) emission thresholds which yields lower $E_n^{(\text{thr})}$ and, hence, larger continuum-coupling correction for odd- N systems. The np continuum-coupling energy correction attenuates the OES and can even wash it out close to drip lines if the ratio $V_0^{(np)}/V_0^{(nn)}$ would not be strongly reduced from its accepted value $\simeq 2$ in well-bound nuclei.

In BCS formalism, the OES is associated with the blocking of a quasi-particle state close to the Fermi energy by an unpaired neutron (proton, respectively) what weakens nn (pp , respectively) pairing correlation in odd- N (odd- Z , respectively) isotopes. The proximity of continuum states in weakly-bound nuclei makes the blocking mechanism less effective, reducing the OES of one-nucleon separation energies. This reduction appears even though the strength of nn (pp) pairing correlations is essentially unchanged.

3.2. Anatomy of the continuum-coupling correction

A typical behaviour of the total continuum-coupling correction $E_{\text{g.s.}}^{(\text{corr})}$ to the g.s. SM energy (the CQS eigenvalue) for different oxygen isotopes is presented in Fig. 1 as a function of the neutron energy. This correction

is an incoherent sum of contributions from couplings to all SM states in the daughter nucleus. $E = 0$ corresponds to a position of the first $1n$ emission threshold. Rapid change of E_{corr} related to opening of next $1n$ emission threshold can be seen for ^{27}O at $E \simeq 3$ MeV. The continuum-coupling energy correction rises with number of valence neutrons and in general is bigger in even- N isotopes. Few notable exceptions can be seen right after the closure of the sd shell (*c.f.* ^{29}O in Fig. 1) and at the beginning of the sd shell. Behavior of the continuum-coupling correction depends on l of the neutron wave involved in a decay channel $[J_i^{\pi; A-1}; (l, j)]^{J_{\text{g.s.}}^{\pi; A}}$. This correction is largest exactly at the threshold of a channel $[J_i^{\pi; A-1}; (l, j)]^{J_{\text{g.s.}}^{\pi; A}}$ only for $l = 0$ neutrons. For higher l -values or for protons, the centrifugal barrier and/or the Coulomb barrier shift the maximum of energy correction above the threshold. This can be seen for ^{28}O , where $l = 2$ neutron wave dominates in the g.s. to g.s. coupling.

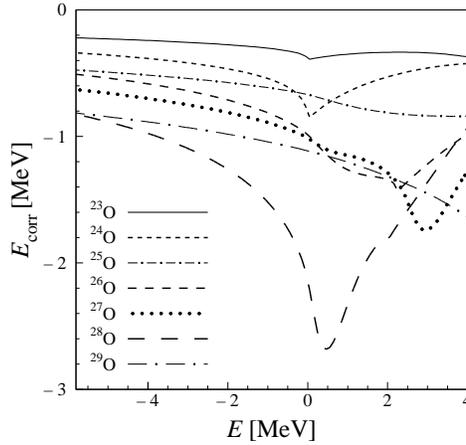


Fig. 1. The continuum-coupling energy correction to the SM g.s. energy for different oxygen isotopes is plotted as a function of the neutron energy. Contributions from couplings to all available states in the daughter nucleus are added incoherently.

Couplings to excited states in $A - 1$ nucleus are less important in even- N isotopes. On the contrary, in odd- N isotopes they dominate. The distribution of contributions to $E_{\text{g.s.}}^{(\text{corr})}$ as a function of the energy of corresponding states in $A - 1$ nucleus, reflects the nature of pairing correlations in odd- N and even- N isotopes. Strong dissimilarity of g.s. to g.s. couplings between odd- N and even- N oxygen isotopes is attenuated if couplings to decay channels involving excited states in $A - 1$ nucleus are taken into account. The fraction of $E_{\text{g.s.}}^{(\text{corr})}$ coming from coupling to the g.s. of daughter nucleus is shown in Figs. 2 and 3 for oxygen and fluorine isotopes, respectively.

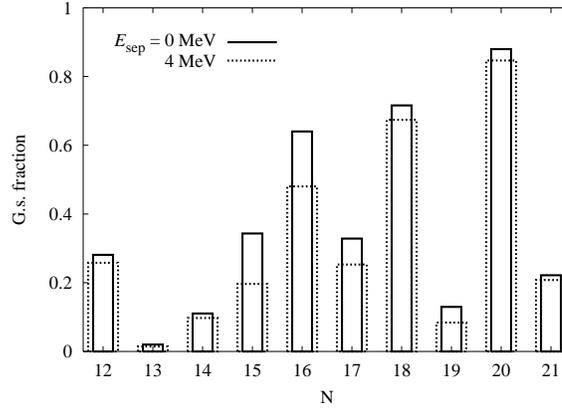


Fig. 2. The fraction of the total continuum-coupling energy correction $E_{\text{g.s.}}^{(\text{corr})}$ in oxygen isotopes, coming from a coupling to the g.s. of a daughter nucleus. SMEC calculations are performed for fixed $1n$ -threshold energies: $E_n^{(\text{thr})} = 0, 4 \text{ MeV}$.

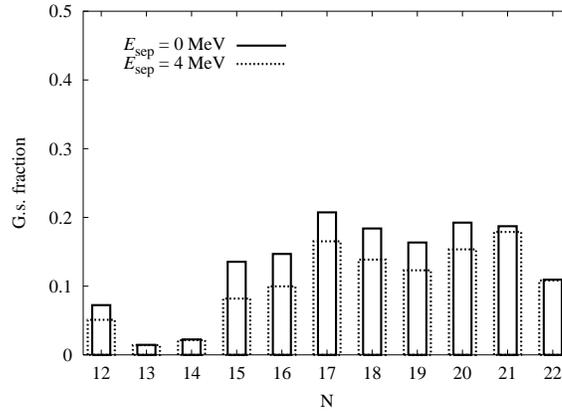


Fig. 3. The same as in Fig. 2 but for fluorine isotopes.

In oxygen isotopes ($T = 1$ couplings), one can see a strong OES following the staggering of neutron pairing correlations in these isotopes. In general, strong nn pairing correlations in even- N isotopes increase the weight of a g.s. to g.s. contribution. On top of this effect, one sees different values of the g.s. fraction at the beginning of the shell (filling of $d_{5/2}$ subshell) and at the end of the shell (filling of $s_{1/2}$ and $d_{3/2}$ subshells). This picture changes qualitatively in fluorine isotopes (*c.f.* Fig. 3). As pointed out in Sec. 3, np continuum coupling dominates in this isotopic chain, removing most of the OES due to the nn correlations. As in oxygen isotopes, the fraction of $E_{\text{g.s.}}^{(\text{corr})}$ coming from coupling to the g.s. of $A - 1$ nucleus is smaller in $d_{5/2}$

subshell than in $s_{1/2}$ and $d_{3/2}$ subshells. Maximal value of this fraction is less than $\sim 20\%$ in fluorine isotopes, whereas for ^{24}O , ^{26}O and ^{28}O the g.s. fraction is $> 60\%$.

4. Mirror-symmetry breaking effect of a continuum coupling: Example of 2_1^+ states in ^{36}Ca and ^{36}S

Symmetry breaking effects in spectra of mirror nuclei are directly or indirectly related to the Coulomb interaction. Direct effects (the Thomas–Ehrmann effect) have been extensively discussed in the literature. Much less known indirect effects of the Coulomb interaction result from different positions of n/p thresholds which modify continuum couplings effects and, hence, change spectra in mirror systems [8]. Recently, the energy of 2_1^+ state has been measured in a weakly-bound ^{36}Ca [10]. An excitation energy of this proton-unbound state differs by ~ 240 keV from an energy of the well-bound mirror state in ^{36}S . In this section, we will discuss in SMEC the mirror-symmetry breaking continuum-coupling effects for these states. Details of SMEC calculations and the choice of an effective interaction are the same as described in Sect. 3 and in Ref. [6].

SMEC excitation energies of 2_1^+ states in ^{36}Ca and ^{36}S are presented in the third column of Table I. The difference of excitation energies in those mirror states is close to the experimental value (*c.f.* the first column of Table I), but absolute excitation energies are ~ 0.5 MeV bigger than experimental ones and ~ 1 MeV bigger than a SM value (*c.f.* the second column of Table I). One should notice that proton continuum couplings for ^{36}Ca and neutron continuum couplings for ^{36}S result in a small difference of 2_1^+ excitation energies in ^{36}Ca and ^{36}S . Almost a whole difference of 2_1^+ excitation energies is due to the neutron-continuum couplings for ^{36}Ca and

TABLE I

Excitation energy of 2_1^+ states in ^{36}Ca and ^{36}S nuclei. Column labelled ‘SM’ contains the SM results. The following three columns present SMEC results with different couplings. The third column shows results with both neutron and proton continua included. The next columns give the 2_1^+ energy with only proton or neutron continuum couplings included. All energies are in MeV.

Exp. (^{36}Ca)	SM (^{36}Ca)	SMEC (^{36}Ca)	$^{36}\text{Ca}[^{35}\text{K} + p]$	$^{36}\text{Ca}[^{35}\text{Ca} + n]$
3.05 ± 0.05	2.641	3.403	3.019	3.514
Exp. (^{36}S)	SM (^{36}S)	SMEC (^{36}S)	$^{36}\text{S}[^{35}\text{S} + n]$	$^{36}\text{S}[^{35}\text{P} + p]$
3.2909	2.641	3.690	3.024	3.822

proton-continuum couplings for ^{36}S . This is due to the same structure of protons in ^{36}Ca and neutrons in ^{36}S for both 0_1^+ g.s. and 2_1^+ excited states (the shell closure).

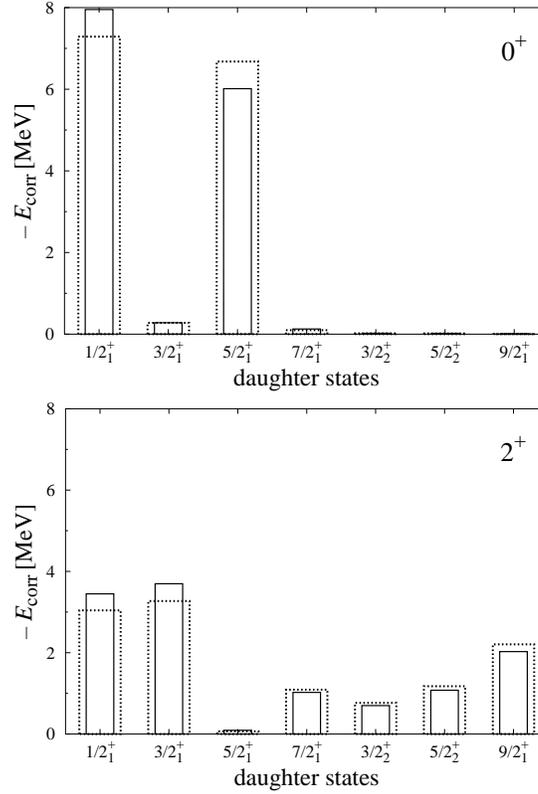


Fig. 4. The states $J_i^{\pi;A-1}$ in channel wave-functions $[J_i^{\pi;A-1}; (l, j)]^{J^{\pi;A}}$ ($J^{\pi;A} = 0_1^+, 2_1^+$) which contribute most to the continuum-coupling correction from neutron continuum in ^{36}Ca (^{35}Ca states — red narrow boxes) and proton continuum in ^{36}S (^{35}P states — green wide boxes). The upper (lower) panel is for 0_1^+ (2_1^+) state.

Let us take a closer look at the individual contributions from couplings via continuum to different states in $A - 1$ nuclei. In Fig. 4, contributions to the continuum-coupling energy corrections for 0_1^+ and 2_1^+ states in ^{36}Ca and ^{36}S are shown for selected daughter states. For 0_1^+ state, the dominant energy contributions from couplings to the g.s. $1/2_1^+$ and the second excited state $5/2_1^+$ of $A - 1$ nucleus are almost equal². The dominant component of SMEC g.s. wave function is:

² The balance of $E_{g.s.}^{(\text{corr})}$ in ^{36}Ca and ^{36}S is a direct consequence of the separation energies in these nuclei: $S_p(^{36}\text{Ca}) = 2.56$ MeV, $S_n(^{36}\text{S}) = 9.89$ MeV.

$$\begin{aligned}
& [(d_{5/2})^6(s_{1/2})^2] + \left[[(d_{5/2})^6(s_{1/2})^1]_{1/2^+} \otimes (s)_{(c)}^1 \right] \\
& + \left[[(d_{5/2})^5(s_{1/2})^2]_{5/2^+} \right] \otimes (d)_{(c)}^1 + \dots, \tag{3}
\end{aligned}$$

where index c denotes continuum state. Important contributions to $E_{2_1^+}^{(\text{corr})}$ are spread over six states in daughter nuclei. Contributions from couplings to the two lowest states ($1/2_1^+$ and $3/2_1^+$) favour ^{36}Ca state, whereas all contributions from couplings to higher excited states ($7/2_1^+$, $3/2_2^+$, $5/2_2^+$ and $9/2_1^+$) favour ^{36}S state. As a result, the 2_1^+ state in ^{36}Ca is shifted down with respect to ^{36}S . The dominant component of the SMEC 2_1^+ wave function is:

$$\begin{aligned}
& [(d_{5/2})^6(s_{1/2})^1(d_{3/2})^1] + \left[[(d_{5/2})^6(s_{1/2})^1]_{1/2^+} \otimes (d)_{(c)}^1 \right] \\
& + \left[[(d_{5/2})^6(d_{3/2})^1]_{3/2^+} \otimes (s)_{(c)}^1 \right] + \left[[(d_{5/2})^5(s_{1/2})^1(d_{3/2})^1]_{9/2^+} \otimes (d)_{(c)}^1 \right] \\
& + \left[[(d_{5/2})^5(s_{1/2})^1(d_{3/2})^1]_{7/2^+} \otimes (d)_{(c)}^1 \right] + \left[[(d_{5/2})^5(s_{1/2})^2]_{5/2^+} \otimes (s)_{(c)}^1 \right] \\
& + \left[[(d_{5/2})^5(s_{1/2})^1(d_{3/2})^1]_{3/2^+} \otimes (s)_{(c)}^1 \right] + \dots. \tag{4}
\end{aligned}$$

If one looks to the $2_1^+ - 0_1^+$ energy difference in ^{36}Ca and ^{36}S , one may notice that a contribution to the energy of 0_1^+ from a coupling to $1/2_1^+$ in $A-1$ nuclei is almost exactly compensated by contributions to the energy of 2_1^+ from couplings to $1/2_1^+$ and $3/2_1^+$ states. The asymmetry in the position of 2_1^+ states in ^{36}Ca and ^{36}S comes from the balance between the contribution of the $5/2_1^+$ in the g.s. and contributions of higher lying states $7/2_1^+$, $3/2_2^+$, $5/2_2^+$ and $9/2_1^+$ in the 2_1^+ state. These two sets of couplings act differently in the g.s. and in the first excited state. It is interesting to notice that the effect is produced by continuum couplings which practically do not involve s -wave neutron/proton components. These components, which are present both in the 0_1^+ state and in the 2_1^+ state (*c.f.* Eq. (4)), mutually cancel out and do not contribute to the Thomas–Ehrmann shift in ^{36}Ca – ^{36}S mirror systems.

The proximity of fp shell leads to important excitations from sd to fp across the $N = 20$, $Z = 20$ shell closure. As a consequence, many new channels $[J_i^{\pi; A-1}; (l, j)]_{g.s.}^{J_i^{\pi; A}}$ become involved in the continuum coupling, leading to an enhanced spreading of contributions to the continuum-coupling energy correction. In general, an increased spreading of couplings leads to an increased fractionation of continuum-coupling contributions and the reduction of $E^{(\text{corr})}$ due to an enhanced interference of large number of channels.

Full SMEC calculations in $sdfp$ shells for ^{36}Ca and ^{36}S are beyond actual possibilities so we proceed by introducing the quenching factor Q_f in the

continuum-coupling energy correction [5]. SM calculations using WBT [12] and IOKIN [13] effective Hamiltonians, which include fp shell and allow for $2\hbar\omega$ excitations, revealed that an appropriate quenching factor should be used to account for admixture of intruder configurations (*c.f.* Fig. 2). Almost identical admixtures of intruder configurations have been found both in parent nuclei and in all important states of $A - 1$ nuclei.

The value of the quenching factor depends on different extensions of sd -shell effective interaction into a larger model space (*c.f.* Table II). It turns out that the excitation energies of 2_1^+ states in ^{36}Ca and ^{36}S are sensitive to the value of the quenching factor (see Fig. 5). Both absolute excitation energies and their difference diminish as the quenching factor decreases. For $Q_f = 0.651$, the SMEC excitation energy of 2_1^+ state reproduces exactly the experimental value in well-bound ^{36}S . This value for the quenching factor is close to the value found for WBT Hamiltonian (*c.f.* Table II) in $psdfp$ shells.

TABLE II

Fraction of $0\hbar\omega$ configurations in $2\hbar\omega$ SM calculations for $^{36}\text{Ca}/^{36}\text{S}$ ($A = 36$) and corresponding daughter nuclei ($A = 35$). Results correspond to 0_1^+ , 2_1^+ states of the parent nucleus ($A = 36$) and to averages over 7 (respectively 8) most important states in $T_z = 5/2$ (respectively $T_z = 3/2$) daughter nuclei, respectively.

Hamiltonian	space	$0_1^+(A = 36)$	$2_1^+(A = 36)$	$\langle T_z = 5/2 \rangle_{av}$	$\langle T_z = 3/2 \rangle_{av}$
IOKIN	sdfp	0.769	0.768	0.772	0.782
WBT	sdfp	0.787	0.789	0.795	0.796
WBT	psdfp	0.687	0.684	0.683	0.685

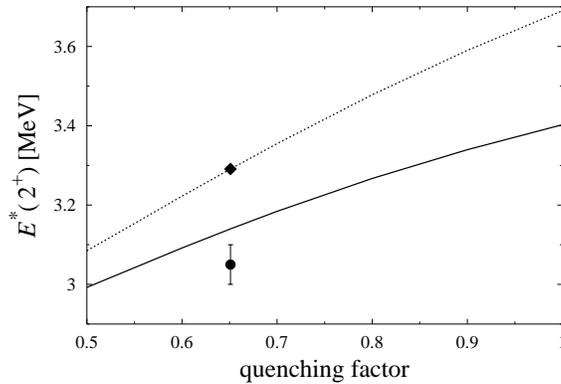


Fig. 5. SMEC excitation energies of 2_1^+ states in ^{36}Ca and ^{36}S are plotted as a function of the quenching factor. Experimental excitation energies of 2_1^+ state in ^{36}S and ^{36}Ca are depicted as diamond and circle, respectively.

For IOKIN and WBT Hamiltonians in $sdfp$ shells, one obtains higher values for Q_f . For those values, SMEC reproduces well an experimental difference of 2_1^+ energies in ^{36}Ca and ^{36}S nuclei, but the absolute values of 2_1^+ energies are too high. The latter problem could be easily resolved by mirror-symmetry conserving correction of the $V_{0d_{5/2};1s_{1/2}}^{T=1}$ monopole.

5. Conclusions

In this paper, we attempted to address two problems pertaining to the physics with exotic, weakly-bound nuclei: (i) What are the generic features of the continuum coupling in binding systematics of neutron-rich nuclei? (ii) Can one explain mirror-symmetry breaking effects in spectra by the asymmetry of threshold energies in mirror nuclei?

The answers to these questions are not simple and require further investigations. Studies in long chains of oxygen and fluorine isotopes revealed that the np residual coupling to the scattering continuum becomes strongly reduced with respect to the nn coupling in the vicinity of the neutron drip line. The np coupling is essential for an understanding of the anti-OES effect which is seen in odd- Z fluorine chain and leads to an apparent reduction of the gap parameter for neutrons. Hence, in weakly-bound nuclei close to the drip lines, the OES has three components: the first one originates from nucleonic pairing, the second one is the deformed mean-field effect [14], and the third one originates from the np coupling via scattering continuum. One should stress that the latter component is of a completely different nature than the singular behaviour of binding energies originating from the np correlations near $N = Z$ line [15].

Asymmetry in the spectra of mirror systems is another playground for the continuum shell model. In the mirror couple ^{36}Ca - ^{36}S , a relative shift of 2_1^+ excitation energies can be to a large extent directly related to the effect of the continuum coupling. However, in contrast to neutron-rich nuclei, the low one-proton separation energy in ^{36}Ca has little influence on the difference of excitation energies of 2_1^+ states in ^{36}Ca and ^{36}S . This is due to the Coulomb barrier which suppresses continuum-coupling effects in weakly-bound systems close to the proton drip line. In that respect, weakly-bound/unbound systems at the proton drip line are radically different from those at the neutron drip line. To understand dissimilarity of nuclear systems at the proton and neutron drip lines is a challenge for the nuclear structure theory.

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