A NEW ISOMERIC STATE IN 134Pr AND EXCITED STATES OF 134Ce

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Radioactive sources of ¹³⁴Pr and ¹³⁴Nd in equilibrium with the ¹³⁴Pr daughter obtained by spallation reaction have been studied with Ge(Li) detectors. A new isomeric state has been found in ¹³⁴Pr. Gamma-gamma coincidence have been measured using two 25 cm³ Ge(Li) detectors. On the basis of gamma intensities, energy sums, and coincidences information decay schemes are proposed for the isomeric and ground state of ¹³⁴Pr.

1. Introduction

Several papers [1-4] have been devoted to the excited states of 134 Ce, one of the nuclei at the edge of the expected region of deformation [5]. Three of these [4, 3] dealt with the (HI, xn) and (α, xn) reactions, and the ground-state quasi-rotational band was revealed. The object of the present work was to obtain further information on the low lying levels excited in the beta decay of 134 Pr. In our studies of the excited states of 134 Ce populated in the decay of 134 Pr we encountered some difficulty in estimating the relative intensity of gamma transitions observed in singles spectra. It turned out that there is a group of gamma-rays whose relative intensity depends on the source history. A more careful analysis of the spectra obtained from 134 Pr and 134 Pr in equilibrium with 134 Nd indicates the existence of an isomeric state in 134 Pr.

2. Source preparation

The sources used in this experiment were obtained by spallation reaction from a $(NH_4)_2$ [Gd DTPA] \cdot H_2O^1 target irradiated with 660 MeV protons on the external beam of the JINR synchrocyclotron. The radioactive elements produced in this reaction were separated from the target material by sorption on the cation exchange resin Dowex

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¹ DTPA — Diethylentriaminpentaacetate.

 50×8 as described in Ref. [6] and according to their atomic number Z by the cation exchange chromatographic method. A 0.5 M buffer solution of ammonium alpha-hydroxy-isobutyrate was used as the eluent. Then the ¹³⁴Pr and ¹³⁴Nd source were prepared by electromagnetic isotope separation of the praseodymium and neodymium fractions using a surface ionization ion source [7]. The admixture of the neighbouring isotopes in the source were estimated to be negligibly small from the singles gamma-ray spectra.

3. Singles gamma-ray spectra

The singles gamma-ray spectra were studied with a 38.5 cc Ge(Li) detector with a resolution of 3.7 keV at 1333 keV. The gamma-ray energies of the larger peaks were determined by counting the spectra simultaneously with several well known standards (60Co,

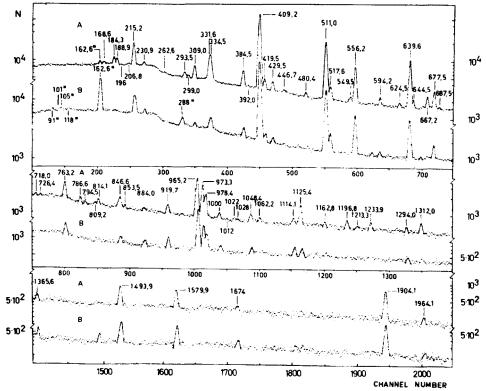


Fig. 1. Singles gamma-ray spectra obtained by adding the counts of 4 measurements: A) the gamma-ray spectrum of ¹³⁴Pr, B) the gamma-ray spectrum of ¹³⁴Nd in equilibrium with the ¹³⁴Pr daughter. Gamma transitions corresponding to ¹³⁴Nd are labeled by crosses

⁵⁴Mn, ⁸⁸Y, ¹³³Ba, ¹⁶⁹Yb, ²⁰⁷Bi) using a computer coupled system of data handling [8]. These peaks, in turn, were used to determine the energies of the weaker peaks in spectra taken without the standard sources. The experimental measurements were usually started 25–30 min after the end of the irradiation. In order to obtain the decrease of the intensities

of the gamma-lines, a few spectra from each source were recorded in succession for a 10-15 min period of accumulation. Measurements were carried out with two different ¹³⁴Pr sources. Figure 1 shows typical gamma-ray spectra observed with the Ge(Li) detector. The upper (A) spectrum was obtained in an experiment on mass-separated sources of the praseodymium fraction. The lower one (B) corresponds to the sources prepared by

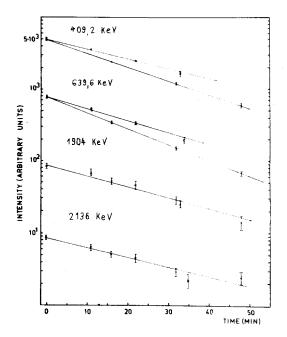


Fig. 2. Decay curves of the most intense gamma-rays associated with the ¹³⁴Pr (full circles) and ¹³⁴Pr in equilibrium with the ¹³⁴Nd parent (crosses), respectively

mass-separation of the neodymium fraction in equilibrium with the ¹³⁴Pr daughter. Gamma transitions corresponding to the ¹³⁴Nd isotope are labeled by crosses. Because of an admixture of praseodymium in the neodymium fraction (≤15%) these spectra can not be compared in detail. However, a comparison of the two spectra shows that the 184.3, 188.9, 667.2, 1125.4 and 1196.8 keV gamma transitions are greatly reduced in the spectrum (B) relative to the (A) one. A more careful analysis of the spectra permitted the gamma-rays to be divided into two groups:

- 1) gamma transitions showing an intensity decrease depending on the way of the source was prepared, and
- 2) gamma transitions for which the decrease of intensity is independent of the source used in the experiments. This is shown on the decay curves for the most intense gamma-lines in the low and high energy parts of the spectra (Fig. 2).

These facts suggest the existence of an activity due to an isomeric state in ¹³⁴Pr. The 11 min half-life given in Fig. 5 for this state was extimated as a mean value of the

TABLE I

Gamma-rays observed in the decay of ¹³⁴Pr and ¹³⁴Nd in equilibrium with the ¹³⁴Pr daughter. The gamma-ray intensity, calculated with a CDC-1604 A computer, were measured ~ 30 min after the end of chromatographic separation

Energy E_{γ}	ΔE _γ (keV)	Relative intensity					
		134Pr		134Nd-134Pr (in equilibrium)			
		I_{γ}	ΔI_{γ}	I_{γ}		ΔI_1	
168.6	0.5	0.5	0.1		weak		
184.3	0.5	1.9	0.3		weak		
188.9	0.5	1.4	0.2		weak		
206.8	0.5	0.3	0.2		weak		
215.1	0.2	9.7	1.0	6.8		1.0	
230.9	0.4	1.0	0.2	1.0		0.3	
293.5	0.5	1.2	0.3	0.3		0.1	
299.0	0.5	0.5	0.2		weak		
309.1	0.3	3.8	0.5	1.0		0.2	
331.6	0.6	7.4	1.0	1.9		0.5	
334.5	0.6	7.3	1.0	2.3		0.6	
384.5	0.4	4.9	0.6	2.2		0.3	
392.0	0.6	0.4	0.1		weak		
409.2	0.1	100	11	100		11	
417.5	0.5	2.5	0.4	1.5		0.4	
429.5	0.4	2.4	0.3	1.2		0.2	
446.7	0.6	0.4	0.2	0.3		0.1	
480.4	0.6	1.3	0.3	0.6		0.2	
517.6	0.5	3.1	1.0	2.1		1.0	
556.2	0.2	14.0	1.7	15.1		1.8	
594.2	0.3	1.7	0.3	0.7		0.2	
639.6	0.2	25.0	2.7	13.8		1.6	
644.5	0.4	1.8	0.3	1.0		0.3	
667.2	0.2	2.5	0.5	0.5		0.3	
677.5	0.2	5.1	0.8	2.7		0.5	
685.7	0.5	1.0	0.4	0.6		0.3	
718.0	0.5	1.1	0.5	0.0	weak	0.3	
763.2	0.2	4.4	0.6	2.4		0.4	
786.6	0.5	1.0	0.3	0.3		0.2	
794.5	0.5	0.4	0.2	0.5	weak		
809.2	0.5	0.4	0.2	0.4	Would	0.2	
814.1	0.5	1.1	0.2	0.4	weak	0.2	
846.6	0.5	2.8	0.4	1.2	Would	0.8	
965.2	0.3	14.1	1.6	15.0		1.8	
973.3	0.2	9.6	1.0	6.0		0.9	
973.3 978.4	0.2	5.7	0.7	3.4		0.6	
1000.0	0.2	1.0	0.7	0.9		0.3	
1062.2	0.5	0.5	0.2	0.9		0.3	
1125.4	0.3	4.2	0.2	1.0		0.1	
1162.8	0.2	0.5	0.6	0.6		0.3	

TABLE I (continued)

Energy E_{γ}	ΔE _γ (keV)	Relative intensity					
		¹³⁴ P _r		¹³⁴ Nd— ¹³⁴ Pr (in equilibrium)			
		I_{y}	ΔI_{γ}	I_{γ}	ΔI,		
1196.8	0.5	1.2	0.3	0.3	0.1		
1213.3	0.6	0.5	0.1	0.3	0.2		
1233.9	0.6	1.1	0.2	0.3	0.1		
1312.0	0.7	1.7	0.7	0.8	0.2		
1365.5	0.3	1.1	0.3	1.6	0.4		
1493.9	0.3	4.0	0.6	5.2	0.7		
1579.9	0.3	3.5	0.5	4.5	0.7		
1904.1	0.3	7.4	1.0	10.1	1.3		
1964.1	1.0	1.1	0.5	1.4	0.7		
2136.0	2	8.0	1.0	10.8	1.6		
2233		1.5	0.3	2.2	0.5		
- 2331		2.1	0.4	2.9	0.5		

half lives of the above-mentioned gamma transitions in the ¹³⁴Pr-¹³⁴Nd equilibrium sources (Fig. 1).

It should be pointed out that a majority of the gamma-lines seen in this experiment are associated with the postulated isomeric state.

4. Gamma-gamma coincidence measurements

Gamma-gamma coincidence measurements were carried out with two 25 cc Ge(Li) detectors oriented at 180°. A coincidence circuit with a time resolution of about 170 ns was used to select the coincidence events stored in a "Minsk-2" computer [9]. Figures 3 and 4 show coincidence spectra obtained simultaneously with eleven gating windows, seven of which were set on the photopeaks of interest. The last ones were set close to these photopeaks on the Compton background associated with the more energetic gamma transitions permitting the spectra to be corrected for the Compton distribution underlying the photopeaks in the windows. Because of the short half-life of the isotopes and the low efficiency of the detectors used in the experiment only the most intense and low-energy gamma transitions could be observed with satisfactory statistics. Figures 3 and 4 show the coincidence spectra obtained by adding the counts of eight measurements.

5. Decay scheme

The coincidence spectrum gated with the most intense 409.2 keV gamma transition (Fig. 4) revealed strong 409.2-639.6 and 409.2-556.2 keV cascades veryfying the existence of the 409.2, 1048.8 and 965.3 keV levels [2, 3, 4]. The latter is supported by the 965.2 keV crossover transition.

The spectra (Figs 3 and 4) gated with the 215.2, 556.2 and 639.6 keV transitions suggest the 215.1 keV transition to be in coincidence both with the 556.2 and 639.6 keV gamma-rays. Moreover, these spectra clearly indicate that the 215.1 keV transition is in coincidence with gamma-rays of energies 846.6 and 763.2 keV and the last one is in coincidence with

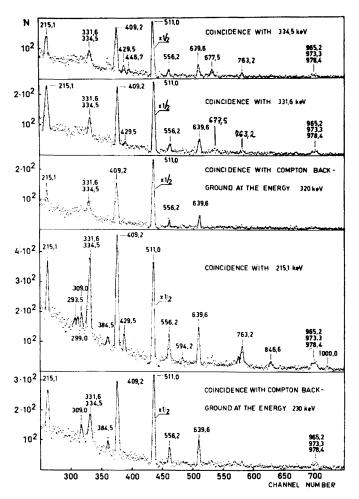


Fig. 3. Gamma-ray spectra in coincidence with the 215.1, 331.6 and 334.5 keV transitions. Spectra associated with gates set close to these photopeaks on the Compton background are also shown

the 639.6 keV transition which, in turn, seems to be in cascade with the 978.4 keV transition. These facts indicate that there are very likely two levels of energies 1812.0 and 2027.1 keV, respectively. The first of these is populated only by the 215.1 keV gamma-ray and depopulated by the above mentioned gamma-rays of energies 763.2 and 846.6 keV. The second one is supported by the 978.4 keV crossover transition and by a 677.5–384.5 keV cascade revealed in the spectra corresponding to the coincidence with the 384.5 and 556.2 keV transitions (Fig. 4) and confirmed by the energy balance.

The 1643.0 keV level, by wihich the 384.5 and 677.5 keV transitions cascade is additionally supported by two gamma-rays of energies 594.2 and 1233.9 keV. The latter, shown in Fig. 5 as a dashed line, is included in the decay scheme on the basis of the energy balance only.

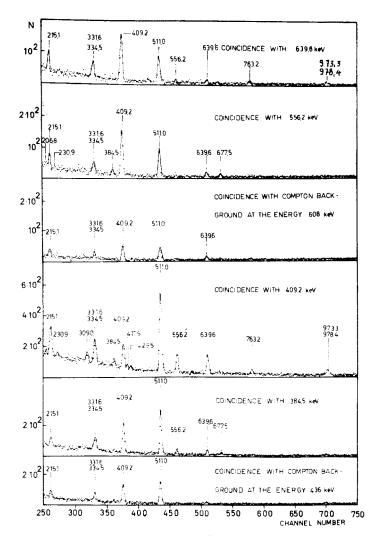


Fig. 4. Gamma-ray spectra in coincidence with the 384.5, 409.2, 556.2 and 639.6 KeV transitions. Gamma-spectra coincident with the Compton background are shown for comparison

In the spectra gated with the 409.2 and 639.6 keV transitions (Fig. 4) a peak in the 973-978 keV region is seen. Because this peak is relatively wide in the coincidence spectra and the 978.4 keV transition was assumed to be in cascade with the 639.6 keV transition, one can conclude that the 973.3 keV gamma-ray is in cascade with the 409.2 or 639.6 keV

transition (it is absent in the coincidence spectrum with the 556.2 keV transition) which locates a level at 1382.5 or 2022.1 keV.

The spectrum in coincidence with the 215.1 keV transition revealed a strong peak in the place of two unresolved gamma-rays of energies 331.6 and 334.5 keV. On the other hand, the spectra gated with windows set at these energies, seem to be identical (within

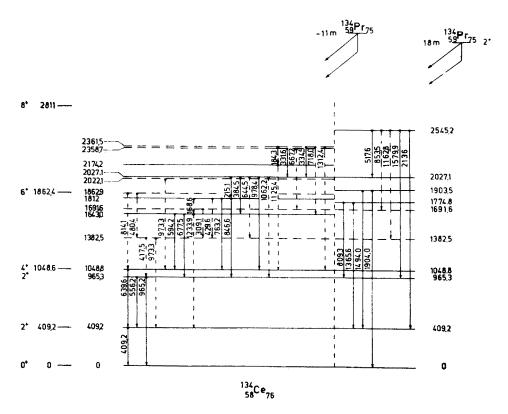


Fig. 5. Decay scheme proposed for 134 Pr. The decay scheme constructed only on the base of the energy sum rules and proposed for the ground state of 134 Pr is shown at the right. The gamma transitions associated with the isomeric state, the placement of which in the decay scheme is suggested by $\gamma - \gamma$ coincidence measurements, are labeled by full lines. Transitions labeled as dashed lines are included in the decay scheme on the basis of the energy sum rules

the statistical error limits). Basing on this information we assumed that the 2027.1 keV state is fed by these two transitions which depopulate two levels of energies 2358.7 and 2361.5 keV. The last one seems to be supported by three gamma-rays observed in single gamma spectra (Fig. 1).

The tentative 2174.2 keV state was incluted in the decay scheme on the basis of the energy sum rule and consistency of the relative intensity decay rate.

The next probable level at 1691.6 keV is based on a weak indication of a 309.1-667.2 keV coincidence (not shown here) and on intensity and energy sum rule.

The right part of the decay scheme shown in Fig. 5 was constructed on the basis of the energy sum rules only. Only these transitions are included for which the decay rates are independent of the source preparation method and for which the relative intensity (related to the 409.2 keV) is increased in the ¹³⁴Pr-¹³⁴Nd equilibrium source. It was concluded, therefore, that these transitions are associated with the ¹³⁴Pr ground state activity.

About 20% of the gamma transitions observed in the singles spectra are not included in the decay scheme shown in Fig. 5.

6. Discussion

A rough extimation of the relative intensity of the annihilation radiation indicates that the ¹³⁴Ce ground state, within the statistical error limits, is not populated in the β^+ decay from the ¹³⁴Pr ground state. Because the β^+ transition to the 2^+ first excited state of ¹³⁴Ce is the most intense one, the spin of the ground state of ¹³⁴Pr is restricted to the values 2 or 3. The 59th and 75th odd nucleons occupy in the neighbouring nuclei the $d_{5/2}$ and $s_{1/2}$ orbitals, respectively. Assuming the same orbitals for the odd nucleons in the odd-odd ¹³⁴Pr nucleus, one could expect that its ground state has a spin of 3 or 2 and positive parity [10]. The value of 2 has been confirmed experimentally by Akström et al. [11]. Assuming that the ground state of ¹³⁴Pr has an oblate shape this value is also in good agreement with calculations in Ref. [16]. If one assumes that the 973.3 and 639.6 keV gamma transitions are in cascade (Fig. 5), the above mentioned assignments can be supported in this experiment.

The branching ratio of the two components depopulating the second excited state in ¹³⁴Ce is consistent with the systematics for even-even nuclei in this mass region [12-15]. According to this systematics, it is believed that this state has spin 2⁺.

The extreme left-hand side of Fig. 5 shows the results obtained by in-beam spectroscopy [3, 4], where a quasi-rotational ground-state band up to $I^{\pi} = 8^{+}$ has been identified. The first two members of this band seen here are consistent with the 2^{+} , 4^{+} assignments [2, 3, 4].

The excited state at 1862.9 keV, tentatively proposed on the basis of energy differences, can be compared with the third member of the ground band of energies 1862.4 [3] and 1864.1 keV [4] observed in the works on (HI, xn) and (α, xn) reactions, respectively. This level would be depopulated by two 480.4 and 814.1 keV gamma-ray emission of approximately equal intensity. The first of these is not reported in the work on the reactions mentioned above.

This fact may suggest that these transitions either have not been properly located in the decay scheme, and the transition energy agreement of the 814.1 keV with those (813.9, 814.8 keV) observed in the reaction [3, 4] is accidental, or there exists an adjacent level of another nature.

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