# SIMPLIFYING 5-POINT TENSOR REDUCTION* 

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The 5-point tensors have the property that after insertion of the metric tensor $g^{\mu \nu}$ in terms of external momenta, all $g^{\mu \nu}$-contributions in the tensor decomposition cancel. If furthermore the tensors are contracted with external momenta, the inverse 5 -point Gram determinant ()$_{5}$ cancels automatically. If the remaining 4-point sub-Gram determinant $\binom{s}{s}_{5}$ is not small then this approach appears to be particularly efficient in numerical calculations. We also indicate how to deal with small $\binom{s}{s}_{5}$. Explicit formulae for tensors of degree 2 and 3 are given for large and small (sub-) Gram determinants.

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## 1. Introduction

In [1] we have worked out an algebraic method to present one-loop tensor integrals in terms of scalar one-loop 1-point to 4-point functions. The tensor integrals are defined as

$$
\begin{equation*}
I_{n}^{\mu_{1} \cdots \mu_{R}}=\int \frac{d^{d} k}{i \pi^{d / 2}} \frac{\prod_{r=1}^{R} k^{\mu_{r}}}{\prod_{j=1}^{n} c_{j}} \tag{1.1}
\end{equation*}
$$

with denominators $c_{j}$, having chords $q_{j}$,

$$
\begin{equation*}
c_{j}=\left(k-q_{j}\right)^{2}-m_{j}^{2}+i \epsilon \tag{1.2}
\end{equation*}
$$

[^0]For the tensor decomposition we use Davydychev's approach [2], recursion relations as given in [3] and make detailed use of modified Cayley determinants introduced for this purpose in [4]. All these techniques we ask the reader to look up in [1]. The following linear combinations of the chords have proven as particularly useful

$$
\begin{align*}
Q_{s}^{\mu} & =\sum_{i=1}^{5} q_{i}^{\mu} \frac{\binom{s}{i}_{5}}{()_{5}}, \quad s=0 \ldots 5,  \tag{1.3}\\
Q_{s}^{t, \mu} & =\sum_{i=1}^{5} q_{i}^{\mu} \frac{\binom{s t}{i t}_{5}}{\binom{t}{t}_{5}}, \tag{1.4}
\end{align*} \quad s, t=1 \ldots 5 . .
$$

## 2. Explicit examples

According to [2] we write the tensor of rank 2 as $\left([d+]^{l}=d+2 l\right.$, $d=4-2 \varepsilon)$

$$
\begin{equation*}
I_{5}^{\mu \nu}=\sum_{i, j=1}^{4} q_{i}^{\mu} q_{j}^{\nu} \nu_{i j} I_{5, i j}^{[d+]^{2}}-\frac{1}{2} g^{\mu \nu} I_{5}^{[d+]} \tag{2.1}
\end{equation*}
$$

Inserting

$$
\begin{equation*}
\frac{1}{2} g^{\mu \nu}=\sum_{i, j=1}^{5} \frac{\binom{i}{j}_{5}}{()_{5}} q_{i}^{\mu} q_{j}^{\nu} \tag{2.2}
\end{equation*}
$$

and using the recursion relation

$$
\begin{equation*}
\nu_{i j} I_{5, i j}^{[d+]^{2}}=-\frac{\binom{0}{j_{5}}_{5}}{()_{5}} I_{5, i}^{[d+]}+\sum_{s=1, s \neq i}^{5} \frac{\binom{s}{j}_{5}}{()_{5}} I_{4, i}^{[d+], s}+\frac{\binom{i}{j}_{5}}{()_{5}} I_{5}^{[d+]}, \tag{2.3}
\end{equation*}
$$

we see that the $g^{\mu \nu}$ cancels with the result

$$
\begin{gather*}
I_{5}^{\mu \nu}=I_{5}^{\mu} Q_{0}^{\nu}-\sum_{s=1}^{5}\left\{Q_{0}^{s, \mu} I_{4}^{s}-\sum_{t=1}^{5} Q_{t}^{s, \mu} I_{3}^{s t}\right\} Q_{s}^{\nu},  \tag{2.4}\\
I_{5}^{\mu}=E Q_{0}^{\mu}-\sum_{s=1}^{5} I_{4}^{s} Q_{s}^{\mu}, \quad E=\frac{1}{\binom{0}{0}_{5}} \sum_{s=1}^{5}\binom{s}{0}_{5} I_{4}^{s} . \tag{2.5}
\end{gather*}
$$

A more direct approach is the use of the 5-point recursion in terms of 4-point functions [5]

$$
\begin{equation*}
I_{5}^{\mu_{1} \ldots \mu_{R-1} \mu}=I_{5}^{\mu_{1} \ldots \mu_{R-1}} Q_{0}^{\mu}-\sum_{s=1}^{5} I_{4}^{\mu_{1} \ldots \mu_{R-1}, s} Q_{s}^{\mu} . \tag{2.6}
\end{equation*}
$$

This formula in general takes into account some cancellations of the $g_{\mu \nu}$. Only $g_{\mu \nu}$-contributions from 4-points have still to be dealt with - and they are simpler to handle. This we demonstrate for the tensor of degree 3. As a special case of (2.6) we have

$$
\begin{equation*}
I_{5}^{\mu \nu \lambda}=I_{5}^{\mu \nu} \cdot Q_{0}^{\lambda}-\sum_{s=1}^{5} I_{4}^{\mu \nu, s} \cdot Q_{s}^{\lambda} . \tag{2.7}
\end{equation*}
$$

The corresponding 4 -point function reads ( $q_{5}=0$ )

$$
\begin{align*}
I_{4}^{\mu \nu, s} & =\sum_{i, j=1}^{4} q_{i}^{\mu} q_{j}^{\nu} \nu_{i j} I_{4, i j}^{[d+]^{2}, s}-\frac{1}{2} g^{\mu \nu} I_{4}^{[d+], s} \\
\nu_{i j} I_{4, i j}^{[d+]^{2}, s} & =-\frac{\binom{0 s}{j s}_{5}}{\binom{s}{s}_{5}} I_{4, i}^{[d+], s}+\frac{\binom{i s}{j s}}{\binom{s}{s}_{5}} I_{4}^{[d+], s}+\sum_{t=1}^{5} \frac{\binom{t s}{j s}}{\binom{s}{s}_{5}} I_{3, i}^{[d+], s t} \tag{2.8}
\end{align*}
$$

and with

$$
\begin{equation*}
\frac{\binom{i s}{s}_{5}}{\binom{s}{s}_{5}}=\frac{\binom{i}{j}_{5}}{()_{5}}-\frac{\binom{s}{i}_{5}\binom{s}{j}_{5}}{()_{5}\binom{s}{s}_{5}} \tag{2.9}
\end{equation*}
$$

we again observe the possibility to cancel $g^{\mu \nu}$ with the result

$$
\begin{align*}
I_{4}^{\mu \nu, s} & =Q_{0}^{s, \mu} Q_{0}^{s, \nu} I_{4}^{s}-\frac{()_{5}}{\binom{s}{s}_{5}} Q_{s}^{\mu} Q_{s}^{\nu} I_{4}^{[d+], s}-\sum_{t=1}^{5}\left\{Q_{t}^{s, \mu} Q_{0}^{s, \nu} I_{3}^{s t}+Q_{t}^{s, \nu} I_{3}^{\mu, s t}\right\} \\
I_{3}^{\mu, s t} & =-\sum_{i=1}^{4} q_{i}^{\mu} I_{3, i}^{[d+], s t}=Q_{0}^{s t, \mu} I_{3}^{s t}-\sum_{u=1}^{5} Q_{u}^{s t, \mu} I_{2}^{s t u} \tag{2.10}
\end{align*}
$$

## 3. Contracting the tensor integrals

Scalar expressions, contracting with chords, are

$$
\begin{align*}
q_{i_{1} \mu_{1}} \cdots q_{i_{R} \mu_{R}} I_{5}^{\mu_{1} \cdots \mu_{R}} & =\int \frac{d^{d} k}{i \pi^{d / 2}} \frac{\prod_{r=1}^{R}\left(q_{i_{r}} \cdot k\right)}{\prod_{j=1}^{5} c_{j}},  \tag{3.1}\\
g_{\mu_{1}, \mu_{2}} q_{i_{1} \mu_{3}} \cdots q_{i_{R} \mu_{R}} I_{5}^{\mu_{1} \cdots \mu_{R}} & =\int \frac{k^{2} d^{d} k}{i \pi^{d / 2}} \frac{\prod_{r=3}^{R}\left(q_{i_{r}} \cdot k\right)}{\prod_{j=1}^{5} c_{j}}, \tag{3.2}
\end{align*}
$$

etc.; (3.1) and (3.2) define the contraction of all tensor indices with chords and the direct contraction of two tensor indices, respectively. These are obtained in realistic matrix element calculations by constructing projection operators or by constructing scalar differential cross sections (Born $\times 1$-loop) before loop integration. As a result of these contractions the $1 /()_{5}$ cancels already.

To begin with, we have a look at
$q_{a \mu} q_{b \nu} I_{5}^{\mu \nu}=\left(q_{a} \cdot I_{5}\right)\left(q_{b} \cdot Q_{0}\right)-\sum_{s=1}^{5}\left\{\left(q_{a} \cdot Q_{0}^{s}\right) I_{4}^{s}-\sum_{t=1}^{5}\left(q_{a} \cdot Q_{t}^{s}\right) I_{3}^{s t}\right\}\left(q_{b} \cdot Q_{s}\right)$.
For $q_{n}=0, a=1, \ldots, n-1, s=1, \ldots n$

$$
\begin{align*}
& \left(q_{a} \cdot Q_{0}\right)=\sum_{j=1}^{n-1}\left(q_{a} \cdot q_{j}\right) \frac{\binom{0}{j}_{n}}{()_{n}}=-\frac{1}{2}\left(Y_{a n}-Y_{n n}\right)  \tag{3.4}\\
& \left(q_{a} \cdot Q_{s}\right)=\sum_{j=1}^{n-1}\left(q_{a} \cdot q_{j}\right) \frac{\binom{s}{j}_{n}}{()_{n}}=\frac{1}{2}\left(\delta_{a s}-\delta_{n s}\right) \tag{3.5}
\end{align*}
$$

and

$$
\begin{equation*}
\left(q_{a} \cdot I_{5}\right)=E\left(q_{a} \cdot Q_{0}\right)-\sum_{s=1}^{5} I_{4}^{s}\left(q_{a} \cdot Q_{s}\right) \tag{3.6}
\end{equation*}
$$

Further

$$
\begin{equation*}
\left(q_{a} \cdot Q_{0}^{s}\right)=\frac{1}{\binom{s}{s}_{5}} \Sigma_{a}^{2, s}, \quad\left(q_{a} \cdot Q_{t}^{s}\right)=\frac{1}{\binom{s}{s}_{5}} \Sigma_{a}^{1, s t} \tag{3.7}
\end{equation*}
$$

where the sums $\Sigma_{a}^{2, s}$ and $\Sigma_{a}^{1, s t}$ are given in [6]. Both are linear combinations of $\binom{s}{s}_{5}$ and Kronecker- $\delta^{\prime}$ 's, $\left(\delta_{a s}-\delta_{5 s}\right)$. Indeed the fact that there is no inverse ()$_{5}$ anymore is due to relations (3.4) and (3.5). The second scalar which can be constructed from the tensor of degree 2 is $g_{\mu \nu} I_{5}^{\mu \nu}$. Due to (2.4) we need to evaluate the following scalar products

$$
\begin{align*}
& \left(Q_{0} \cdot Q_{0}\right)=\frac{1}{2}\left[\frac{\binom{0}{0}_{5}}{()_{5}}+Y_{55}\right] \\
& \left(Q_{0} \cdot Q_{s}\right)=\frac{1}{2}\left[\frac{\binom{s}{0}_{5}}{()_{5}}-\delta_{s 5}\right] \\
& \left(Q_{0}^{s} \cdot Q_{s}\right)=-\frac{1}{2} \delta_{s 5} \\
& \left(Q_{t}^{s} \cdot Q_{s}\right)=0 \tag{3.8}
\end{align*}
$$

In this case the terms with $1 /()_{5}$ cancel and, not surprisingly, the result finally is

$$
\begin{equation*}
g_{\mu \nu} I_{5}^{\mu \nu}=\frac{Y_{55}}{2} E+I_{4}^{5} . \tag{3.9}
\end{equation*}
$$

To calculate $g_{\mu \nu} I_{5}^{\mu \nu \lambda}$ we need $g_{\mu \nu} I_{4}^{\mu \nu, s}$ and thus further scalar products, see (2.10)

$$
\begin{align*}
\left(Q_{0}^{s} \cdot Q_{0}^{s}\right) & =\frac{1}{2\binom{s}{s}_{5}}\left[\binom{0 s}{0 s}_{5}+2\binom{s}{0}_{5} \delta_{s 5}\right]+\frac{1}{2} Y_{55} \\
\left(Q_{s} \cdot Q_{s}\right) & =\frac{1}{2} \frac{\binom{s}{s}}{()_{5}} \\
\left(Q_{t}^{s} \cdot Q_{0}^{s}\right) & =\frac{1}{2\binom{s}{s}_{5}}\left[\binom{t s}{0 s}_{5}-\binom{s}{s}_{5} \delta_{t 5}+\binom{s}{t}_{5} \delta_{s 5}\right], \\
\left(Q_{t}^{s} \cdot Q_{0}^{s t}\right) & =\frac{1}{2\binom{s}{s}_{5}}\left[-\binom{s}{s}_{5} \delta_{t 5}+\binom{s}{t}_{5} \delta_{s 5}\right] \\
\left(Q_{t}^{s} \cdot Q_{u}^{s t}\right) & =0 \tag{3.10}
\end{align*}
$$

which yields

$$
\begin{equation*}
g_{\mu \nu} I_{4}^{\mu \nu, s}=\frac{Y_{55}}{2} I_{4}^{s}+I_{3}^{s 5}+\frac{\delta_{s 5}}{\binom{s}{s}_{5}}\left[\binom{s}{0}_{5} I_{4}^{s}-\sum_{t=1}^{5}\binom{s}{t}_{5} I_{3}^{s t}\right] \tag{3.11}
\end{equation*}
$$

and finally

$$
\begin{equation*}
g_{\mu \nu} I_{5}^{\mu \nu \lambda}=-\frac{Y_{55}}{2} \sum_{s=1}^{5} I_{4}^{s} Q_{s}^{0, \lambda}+I_{4}^{5} Q_{0}^{5, \lambda}-\sum_{t=1}^{4} I_{3}^{5 t} Q_{t}^{5, \lambda} . \tag{3.12}
\end{equation*}
$$

It is remarkable that (3.11) is trivial again for $s \neq 5$. For $s=5$, however, the standard cancellation of propagators does not work and for this case (3.11) is indeed a useful result. For further contraction of (3.12) with a vector $q_{\lambda}$ again (3.7) can be applied.

## 4. Avoiding inverse 4-point Gram determinants

While in the above approach of taking scalar products of the tensors with chords the inverse ()$_{5}$ Gram determinant cancels already, there still remains the inverse $\binom{s}{s}_{5}$, sub-Gram determinant of the 4 -point functions. Therefore, we have to choose a different approach for the case if the latter becomes small. This approach consists in avoiding the inverse ()$_{5}$ already
in the 5 -point tensors from the very beginning and keeping only 4-point integrals in higher dimensions (i.e. integrals with only powers 1 of the scalar propagators), which for small $\binom{s}{s}_{5}$. should be evaluated in a different manner than by standard recursion, see [1]. If one does not want to reintroduce the inverse ()$_{5}$ to see the cancellation of the $g^{\mu \nu}$, one can explicitely see its cancellation also after taking contractions. For the tensor of degree 2 we refer to [6] for the contraction with two chords. For the self-contraction of the tensor indices no simpler result than (3.9) can be achieved anyway. For the tensor of degree 3 we present new results for the contraction with three chords and a self-contraction.

The tensor can be written as follows (see [1] (4.35)-(4.37))

$$
\begin{equation*}
I_{5}^{\mu \nu \lambda}=\sum_{i, j, k=1}^{5} q_{i}^{\mu} q_{j}^{\nu} q_{k}^{\lambda} E_{i j k}+\sum_{k=1}^{5} g^{[\mu \nu} q_{k}^{\lambda]} E_{00 k} \tag{4.1}
\end{equation*}
$$

with

$$
\begin{align*}
E_{00 k} & =\sum_{s=1}^{5} \frac{1}{\binom{0}{0}_{5}}\left[\frac{1}{2}\binom{0 s}{0 k}_{5} I_{4}^{[d+], s}-\frac{d-1}{3}\binom{s}{k}_{5} I_{4}^{[d+]^{2}, s}\right]  \tag{4.2}\\
E_{i j k} & =-\sum_{s=1}^{5} \frac{1}{\binom{0}{0}_{5}}\left\{\left[\binom{0 j}{s k}_{5} I_{4, i}^{[d+]^{2}, s}+(i \leftrightarrow j)\right]+\binom{0 s}{0 k}_{5} \nu_{i j} I_{4, i j}^{[d+]^{2}, s}\right\} . \tag{4.3}
\end{align*}
$$

Contraction of the tensor with three chords yields

$$
\begin{align*}
& q_{a \mu} q_{b \nu} q_{c \lambda} I_{5}^{\mu \nu \lambda}=\sum_{i, j, k=1}^{4}\left(q_{a} \cdot q_{i}\right)\left(q_{b} \cdot q_{j}\right)\left(q_{c} \cdot q_{k}\right) E_{i j k} \\
& +\sum_{k=1}^{4}\left[\left(q_{a} \cdot q_{b}\right)\left(q_{c} \cdot q_{k}\right)+\left(q_{a} \cdot q_{c}\right)\left(q_{b} \cdot q_{k}\right)+\left(q_{b} \cdot q_{c}\right)\left(q_{a} \cdot q_{k}\right)\right] E_{00 k} \tag{4.4}
\end{align*}
$$

Introducing $Y_{a}=Y_{a 5}-Y_{55}, D_{a}^{s}=\delta_{a s}-\delta_{5 s}$ and the following kinematical objects,

$$
\begin{align*}
P_{I 4}= & \frac{1}{8} \frac{1}{\binom{0 s}{0 s}_{5}}\left\{\binom{s}{s}_{5}\left[\binom{s}{0}_{5} Y_{a} Y_{b} Y_{c}+\binom{0}{0}_{5}\left(Y_{a} Y_{b} D_{c}^{s}+Y_{a} Y_{c} D_{b}^{s}+Y_{b} Y_{c} D_{a}^{s}\right)\right]\right. \\
& \left.+\binom{0}{0}_{5}\left[\binom{s}{0}_{5}\left(Y_{a} D_{b}^{s} D_{c}^{s}+Y_{b} D_{a}^{s} D_{c}^{s}+Y_{c} D_{a}^{s} D_{b}^{s}\right)+\binom{0}{0}_{5} D_{a}^{s} D_{b}^{s} D_{c}^{s}\right]\right\},(4.5)  \tag{4.5}\\
P_{Z 4}= & P_{I 4}-\frac{1}{12}()_{5}\left\{Y_{a} Y_{b} D_{c}^{s}+Y_{a} Y_{c} D_{b}^{s}+Y_{b} Y_{c} D_{a}^{s}\right\} \text { and }  \tag{4.6}\\
P_{I 3}= & \frac{1}{24} \frac{()_{5}}{\binom{0 s}{0 s}_{5}}\left\{\left[Y_{a} D_{b}^{s}+Y_{b} D_{a}^{s}\right]\left[\binom{0 s}{0 s}_{5} D_{c}^{t}-\binom{0 s}{0 t}_{5} D_{c}^{s}\right]+(a \leftrightarrow c)+(b \leftrightarrow c)\right\}, \tag{4.7}
\end{align*}
$$

we can write

$$
\begin{align*}
q_{a \mu} q_{b \nu} q_{c \lambda} I_{5}^{\mu \nu \lambda}= & \frac{d-2}{8} \frac{()_{5}}{\binom{0 s}{0 s}}\left(\delta_{a b} \delta_{a c} \delta_{a s}-\delta_{5 s}\right)(d-1) I_{4}^{[d+]^{2}, s} \\
& -\frac{1}{\binom{0}{0}_{5}}\left\{P_{I 4} I_{4}^{[d+], s}-P_{Z 4} Z_{4}^{[d+], s}+P_{I 3} I_{3}^{[d+], s t}\right. \\
& \left.+\frac{1}{3}\left[\Sigma_{c}^{1, s} R_{a b}+\Sigma_{b}^{1, s} R_{a c}+\Sigma_{a}^{1, s} R_{c b}\right]\right\}+F_{a b c}^{s} \tag{4.8}
\end{align*}
$$

symmetric in the indices $a, b, c$ and summation over $s, t$ assumed. Here

$$
\begin{align*}
R_{a b}= & -\frac{1}{\binom{0 s}{0 s}_{5}}\left\{\frac{1}{\binom{0 s}{0 s}_{5}} \Sigma_{b}^{2, s} \Sigma_{a}^{2, s t}(d-2) I_{3}^{[d+], s t}\right. \\
& \left.+\frac{1}{\binom{0 s t}{0 s t}_{5}} \Sigma_{b}^{2, s t}\left[\Sigma_{a}^{3, s t}(d-2) I_{3}^{[d+], s t}-\sum_{u=1}^{5} \Sigma_{a}^{2, s t u} I_{2}^{s t u}\right]\right\} \tag{4.9}
\end{align*}
$$

contains only 3 -point functions and no inverse $\binom{s}{s}_{5}$.

$$
\begin{equation*}
F_{a b c}^{s}=-\frac{1}{24} \frac{()_{5}}{\binom{0 s}{0 s}_{5}} \frac{\binom{s}{0}_{5}}{\binom{0}{0}_{5}}\left[Y_{c} D_{a}^{s} D_{b}^{s}+Y_{b} D_{a}^{s} D_{c}^{s}+Y_{a} D_{b}^{s} D_{c}^{s}\right] \tag{4.10}
\end{equation*}
$$

is a rational term obtained from an $\varepsilon$-expansion. The fact that no scalar products from (4.4) remain demonstrates that the $g^{\mu \nu}$ term has canceled.

For the selfcontracted tensor we obtain

$$
\begin{aligned}
& q_{a \lambda} I_{5, \mu}{ }^{\mu \lambda}=-\frac{1}{\binom{0}{0}_{5}} \sum_{s=1}^{5}\left\{\Sigma_{a}^{1, s}\left[1+\frac{1}{2} \frac{1}{\binom{0 s}{0 s}_{5}}\left(\binom{s}{s}_{5} Y_{55}+2\binom{s}{0}_{5} \delta_{s 5}\right)\right]\right. \\
& \left.+()_{5}\left(Y_{a 5}-Y_{55}\right) \delta_{s 5}\right\} I_{4}^{[d+], s} \\
& -\frac{1}{\binom{0}{0}_{5}} \sum_{s, t=1}^{5} \Sigma_{a}^{1, s} \frac{\binom{t s}{0 s}_{5}}{\binom{0 s}{0 s}} \frac{Y_{5}}{\binom{05 t}{0 s t}_{5}}\left[\frac{d-2}{2}\binom{s t}{s t}_{5} I_{3}^{[d+], s t}+\frac{1}{2} \sum_{u=1}^{5}\binom{s t 0}{s t u}_{5} I_{2}^{s t u}\right] \\
& +\frac{\delta_{s 5}}{\binom{0}{0}_{5}} \sum_{t=1}^{5} \Sigma_{a}^{1, s} \frac{1}{\binom{0 s}{0 s}_{5}\binom{0 s t}{0 s t}_{5}}\left\{\left[\binom{0 s}{0 t}_{5}\binom{s t}{s t}_{5}-\binom{t s}{0 s}_{5}\binom{s t}{0 t}_{5}\right] \frac{d-2}{2} I_{3}^{[d+], s t}\right. \\
& +\sum_{u=1}^{5}\left[\binom{0 s}{0 t}_{5}\right. \\
& \left.\left.\binom{s t 0}{s t u}_{5}-\binom{t s}{0 s}_{5}\binom{s t 0}{0 t u}_{5}\right] \frac{1}{2} I_{2}^{s t u}\right\} \\
& -\frac{\delta_{t 5}}{\binom{0}{0}_{5}} \sum_{s=1}^{5} \Sigma_{a}^{1, s} \frac{1}{\binom{0 s}{0 s}_{5}\binom{0 s t}{0 s t}_{5}}\left\{\left[\binom{0 s}{0 s}_{5}\binom{s t}{s t}_{5}+\binom{t s}{0 s}_{5}\binom{t s}{0 s}_{5}\right] \frac{d-2}{2} I_{3}^{[d+], s t}\right.
\end{aligned}
$$

$$
\begin{align*}
& \left.+\sum_{u=1}^{5}\left[\binom{0 s}{0 s}_{5}\binom{s t 0}{s t u}_{5}+\binom{t s}{0 s}_{5}\binom{t s 0}{0 s u}_{5}\right] \frac{1}{2} I_{2}^{s t u}\right\} \\
& +\frac{1}{\binom{0}{0}_{5}} \sum_{s, t=1}^{5} \Sigma_{a}^{1, s} \frac{1}{\binom{0 s}{0 s}_{5}}\left\{\left[\binom{s}{t}_{5} \delta_{s 5}-\binom{s}{s}_{5} \delta_{t 5}\right] \frac{d-2}{2} I_{3}^{[d+], s t}-\binom{t s}{0 s}_{5} \frac{1}{2} I_{2}^{s t 5}\right\} \\
& +\frac{1}{\binom{0}{0}_{5}} \sum_{s, t=1}^{5}\left[\binom{t}{0}_{5}\left(\delta_{a s}-\delta_{5 s}\right)-\binom{s}{0}_{5}\left(\delta_{a t}-\delta_{5 t}\right)\right] I_{3}^{[d+], s t} . \tag{4.11}
\end{align*}
$$

Only the first two lines of (4.11) contain complete double sums while the remaining terms contribute only for specific values of $s, t, u$.
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