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BOOTSTRAP TECHNIQUE TO STUDY CORRELATION BETWEEN NEUTRON SKIN THICKNESS AND THE SLOPE OF SYMMETRY ENERGY IN ATOMIC NUCLEI*

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We examine the proposed correlation between the neutron skin thickness and the slope of the symmetry energy for two doubly-magic nuclei $(^{100}$ Sn and 132 Sn) and find that a statistically significant correlation appears to exist for 132 Sn whilst no such correlation appears to exist for 100 Sn. Additionally, we present a new statistical tool based on random sampling to assess the confidence interval of Pearson's and Spearman's correlation coefficients in order to examine the aforementioned correlation. These estimators are then used to quantify the statistical correlations among the neutron skin thickness of atomic nuclei and the slope of the symmetry energy in the infinite nuclear medium.

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1. Introduction

The possible correlation observed between the neutron skin thickness in 208 Pb and the slope of the symmetry energy in the infinite nuclear mediums has attracted the interest of the scientific community [1]. The symmetry energy and its derivative are crucial ingredients for the Equation of State of massive astrophysical objects such as neutron stars [2, 3]. Both quantities

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are, strictly speaking, pseudo-observables as they cannot be measured directly. They are typically extracted from other measurements on atomic nuclei [4] or more fundamental theoretical calculations [5, 6].

The neutron skin of heavy nuclei has been identified as a possible candidate to assess the slope of the symmetry energy in the nuclear medium L_0 [7]. A series of theoretical studies have shown that using the suggested correlation between the neutron skin thickness Δr_{np} and L_0 , it would be possible to extract reliable confidence intervals for L_0 by performing accurate measurements of Δr_{np} . This has been a strong motivation for an intense experimental campaign to measure the skin quantity with the best possible accuracy [8, 9].

In the present article, we perform an analysis of the suggested correlation in two doubly-magic tin isotopes. Compared to previous works, we discuss a new statistical tool named Bootstrap analysis. This technique provides us with access to the simulated parent distribution of the estimator we apply on our data set and thus provides access to the confidence interval without assuming a Gaussian error.

The article is organised as follows: in Sec. 2, we revise the definition of neutron skin and in Sec. 3, we introduce the Bootstrap method. Our results are discussed in Sec. 4 and in Sec. 5, we present our conclusions.

2. Neutron skin

The neutron skin thickness $\Delta r_{np} = \langle r_n^2 \rangle^{1/2} - \langle r_p^2 \rangle^{1/2}$ is defined as the difference between the radial extension of the neutron matter density against the proton matter density. This quantity is extracted using a standard procedure based on the parametrisation of the nuclear density via a 2 parameter Fermi function (2pF) [10] given by

$$\rho_q(r) = \frac{\rho_{0q}}{1 + \exp\left[(r - C_q)/a_q\right]},$$
(1)

where q = n, p is an index representing either a neutron (n) or a proton (p) respectively. C_q , a_q are parameters fixed on the self-consistent nuclear density extracted from a fully self-consistent Hartree–Fock–Bogoliubov (HFB) calculation [10, 11]. The values ρ_{0q} are merely determined by the number of neutrons (N) and protons (Z) in a given nucleus.

Following Refs. [7, 12], we consider in our calculations different nuclear functionals, including both non-relativistic — Skyrme [13] and Gogny [14] functionals — and relativistic ones [15] as well. The Skyrme functionals have been selected from Ref. [16] to span as large a range of variation of the L_0 parameter as possible. Moreover, we have also used different functionals with different density-dependent forms [17] to enrich our model space. For the Gogny interaction, we used D1S [14] and D1M [18]. For the relativistic calculations, we used the DD-ME Lagrangians [19].

3. Statistical tools

Non-Parametric Bootstrap (NPB) is a statistical method introduced by Efron in 1979 [20] as a statistical tool to evaluate the bias of some particular statistical estimators. NPB is based on the simple assumption that any experimental data set contains information about its parent distribution thus, if our data set is sufficiently large, we can simply replace the original parent distribution via the empirical one obtained from the sample. The latter is then approximated using a Monte Carlo method and performing a re-sampling of the original data set. By exploring a large subset of the possible combinations (with repetitions), we build out from our original data set and effectively obtain the empirical distribution for the given estimator. We generate new data sets by randomly sampling from the empirical distribution. For the new data sets, we calculate the corresponding estimator and repeat this process a sufficiently large number of times in order to minimize bias. For further details, see Ref. [21].

The only hypothesis behind NPB is that the parent distribution has finite moments and the data of the sample are independent. In the present implementation, we use equal weighting for each data point but more evolved bootstrap techniques based on Bayesian inference allow such an assumption to be bypassed [22].

The use of NPB helps us in highlighting the possible presence of outlier points that may artificially drive the correlation. Since we allow repetitions, it is not guaranteed that in each re-sampled set we find *all* the original data points. This means that the way correlations are calculated for each data set may differ. Having a peaked distribution would mean that the result is quite robust and independent of the particular choice of the data, a larger distribution would mean that a careful choice of the input data set may lead to different conclusions. In our case, this would be reflected by a higher uncertainty on the resulting error bars of the estimator.

For the current article, we apply the NPB approach to the calculations of the Pearson (r) and Spearman (ρ) coefficients [23]. The first is a statistical test to verify a possible linear correlation amongst the data, while ρ is a non-parametric measure of rank correlation. It assesses how well the data can be described using a monotonic function. They are both constructed to vary from -1 to +1 with the same meaning: for |r| = 1 ($|\rho| = 1$), there is a clear linear (monotonic) correlation among the data and for r = 0 ($\rho = 0$), there is no such correlation. We refer to Ref. [23] for an extensive discussion of the properties of the two estimators.

4. Results

In Fig. 1 (a), (b) we present the values of the neutron skin thickness in 100 Sn (panel (a)) and 132 Sn (panel (b)) as a function of the slope of the symmetry energy L_0 . The latter is extracted as a density derivative of the symmetry energy parameter calculated at saturation density $\rho_0 \sim 0.16 \text{ fm}^{-3}$ [16]. We have used a total of 60 functionals spanning a very large interval of L_0 . We observe that in the case of 100 Sn, a small proton skin develops, mainly due to the Coulomb repulsion since it is roughly independent of the particular choice of the L_0 parameter. In 132 Sn, a quite extended neutron skin develops and in this case, we identify a clear trend as a function of L_0 .



Fig. 1. Neutron skin thickness Δr_{np} calculated for ¹⁰⁰Sn (panel (a)) and ¹³²Sn (panel (b)) as a function of L_0 for different functionals. See the text for details.

To quantify the possible correlation, we apply both the Pearson and Spearman tests on these two sets of data. The results are reported in Table I.

TABLE I

The Pearson and Spearman coefficients for 100 Sn and 132 Sn. The errors refer to the 95% confidence limits derived by inspecting the distribution of each estimator.

100 Sn	^{132}Sn
$r = 0.03^{+0.47}_{-0.51}$	$r = 0.91^{+0.05}_{-0.06}$
$\rho = -0.17^{+0.38}_{-0.36}$	$\rho = 0.85^{+0.08}_{-0.11}$

The error bars are defined as 95% quantile of the distribution of the estimators for \hat{r} and $\hat{\rho}$ extracted via the NPB and illustrated in Fig. 2. These histograms have been derived by performing 30,000 bootstrap simulations and performing for each of them the corresponding Pearson or Spearman correlation coefficient calculation.



Fig. 2. Histograms of the \hat{r} and $\hat{\rho}$ estimators as derived via the NPB approach for ¹⁰⁰Sn (left panel) and ¹³²Sn (right panel). See the text for details.

We conclude that for ¹³²Sn, there is an observable linear correlation. For ¹⁰⁰Sn, there is no such correlation shown so the null-hypothesis of noncorrelated data cannot be rejected. It is worth recalling that the statistical test evaluates the correlation amongst the data set but it does not establish if such a correlation is related to a physical phenomenon or it is simply an artefact of the model used to extract the data.

5. Conclusions

We have applied a new statistical tool named Non-Parametric Bootstrap to assess the error bars of Pearson and Spearman correlation estimators. The main advantage of the NPB method is to produce robust confidence intervals that allow us to make the result on the estimator independent of the particular choice of the input data set used.

We have applied such a method to two doubly-magic nuclei and we have observed that we have a relatively high statistical correlation between the resulting neutron skin in ¹³²Sn and the slope of the symmetry energy, while we have no significant statistical correlation between the proton skin of ¹⁰⁰Sn and L_0 .

In the near future, we plan to apply such a technique to a large variety of nuclei to assess the evolution of the statistical correlation as a function of the isospin asymmetry.

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