HISTORY OF $N(1680)^*$

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This paper describes my personal appreciation for some of the great research achievements of Mitya Diakonov, Vitya Petrov, and Maxim Polyakov, and how my own research career has followed the paths they opened. Among the topics where they have been the most influential have been the pursuit and study of the exotic pentaquark. The search for exotics may require a complementary approach, such as experimental and theoretical activity. Here, I would like to focus on a story of N(1680), a non-strange unitary partner of Θ^+ , in which I was involved.

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Dedication

This paper is dedicated to the memory of Mitya, Vitya, and Maxim. Perhaps now that they are no longer with us, we can better appreciate their scientific achievements.

1. Prehistory

I am an experimentalist and served in the Leningrad Nuclear Physics Institute (LNPI or PNPI now) High Energy Physics Division (HEPD) for 29 years. In particular, the reaction $\pi p \rightarrow p\pi \pi$ was an important part of the HEPD research. From my friends Senya Sherman (experimentalist) and Tolya Bolokhov (theorist), I had the opportunity to learn about an extraordinary student, Maxim, who was heavily involved in this activity, and he had several joint papers with seniors back to the end of 1980s.

Somewhere in 1997 or so (the time of the famous $\overline{10}$ paper [1]), Dick Arndt forwarded me a letter from Maxim asking what the Virginia Tech Scattering Analysis Interactive Dial-in (SAID) kaon–nucleon Partial-Wave

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Analysis (PWA) [2] gives for the exotic S = +1 baryon. Dick was very surprised, since the analysis gave nothing which may fit such a narrow exotic state. All four resonances in standard PWA are too heavy (mass > 1790 MeV) and too broad (width > 95 MeV). So they do not fit the predicted Θ^+ in [1]. This fact was a motivation for us to solve a standard PWA problem, and we do a modified PWA to look for narrow resonances [3].

It was a good surprise for me to see that the famous $\overline{10}$ paper [1] used our SAID sigma-term result [4] for the mass difference in anti-decuplet using the Gell-Mann–Okubo phenomenological mass formula.

The experimental results published in the top peer-review journals had problems for experimentalists and the main reason was English. I was a junior researcher at PNPI and was willing to write papers on $\pi^{\pm}p$ elastic scattering and proton-induced pion production to report the results of our group measurements at Gatchina. I translated the Russian versions into English and asked Mitya to polish them. Sure, Mitya ignored my silly staff, and one eve per paper was enough for him to make drafts which *Nucl. Phys. A* accepted in 1981 and then *J. Phys. G* accepted in 1988. It will be good to acknowledge Mitya for his help in translation, but our experimental physics was far away from his mainstream, and for that reason Mitya's name is still hidden in both papers.

In 2005 or so, Mitya gave a talk at a JLab seminar about pentaquark, and Dick Arndt, who was sitting next to me, asked where this guy got such English. My answer was simple — his mother (Nina Yakovlevna) was a professor of English literature in St. Petersburg State University and translated a poem by John Keats from English into Russian. Dick was satisfied with my answer.

I met Vitya for the first time at one of the LNPI Winter Schools at the end of the 1970s, but there is a background. In 1968, I measured the effciency of neutron detection with large scintillation counters, using a beam of neutrons from the Neutron Generator (that was my Master project). At the same time, a physicist and an engineer studied the reaction of transistors to neutrons using the same source of fast neutrons. The physicist's name was Tamara Markovna. One day, a strange man appeared looking for his wife. Somebody told me who he was, Yurii Viktorovich Petrov, a well-known theoretician. That is how I found Tamara Markovna's last name. They were Vitya's parents and Vitya was probably 13 years old at the time.

Obviously, the Instanton vacuum model developed by Mitya and Vitya is a very powerful model [5]. Before the $\overline{10}$ tasks came up, our friends focused on the nucleon-nucleon interaction and, in particular, on the dibaryon (see, for instance, Ref. [6]). In 1996 or 1997, I asked my postdoc, Serezha Pavlenko, to work with Mitya and Vitya because that was a time when many groups were looking for dibaryon. He was a brilliant young experimentalist and phenomenologist. In a short time, Serezha did many calculations and Mitya and Vitya were impressed with the Serezha results. Unfortunately, he passed away shortly after the stroke. In 1998, Serezha was 24 years old.

2. Introduction

QCD gives rise to the hadron spectrum [7] and many $q\bar{q}$ and qqq have been observed [8]. However, $q\bar{q}q\bar{q}$ and $qqqq\bar{q}$ and other many-quark states are not forbidden but have not yet been observed. Recently, the LHCb Collaboration claimed evidence for four hidden-charm $qqqc\bar{c}$ states near open-charm decay thresholds for $\Sigma_c^+\bar{D}^0$ and $\Sigma_c^+\bar{D}^{*0}$ in the $\Lambda_b^0 \to P_{c\bar{c}}^+K^- \to J/\psi pK^-$ decay [9]. However, although there is no doubt about the LHCb observations, it is not clear whether they are compact multiquark or hadonic molecular states [10].

In the light-quark sector, Mitya, Vitya, and Maxim proposed a clearly exotic and narrow $\Theta^+(uudd\bar{s})$ state [11], yet to be unequivocally observed and identified. The original name of the pentaquark lying at the apex of $\overline{10}$ was Z^+ , then following Mitya's suggestion, we now call it Θ^+ (see Appendix). Using the Chiral Quark Soliton model [11], this particle along with other members of $\overline{10}$ has been proposed with a mass of $M_{\Theta^+} = 1.53$ GeV and a width less than 15 MeV [1] (Fig. 1). Due to the relatively low mass and simple decay channels to K^+n or K^0p , it has attracted the attention of many experiments at different facilities around the world (this renowned paper [1] has about 1000 citations collected by the Citation Index (CI) at inSpire-hep).

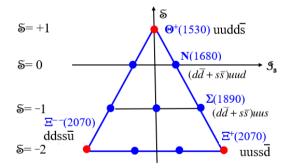


Fig. 1. The suggested anti-decuplet of baryons [1]. A weight diagram for antidecuplet formed by using two flavor conjugate diquarks $(\bar{q})_{qq}$ and one anti-quark \bar{q} [12]. Here $q\bar{q}$ denotes $u\bar{u}$, $d\bar{d}$, and $s\bar{s}$. The three corners of the triangle (in red) are exotic, which means that their quantum numbers require more than three quarks.

The first experimental evidence for Θ^+ came from the LEPS Collaboration at SPring-8 (investigating the $\gamma n \to K^+ K^- n$ reaction on ¹²C) [13] and the DIANA Collaboration from ITEP (investigating the $K^+ n \to K^0 p$ reaction using data from the Xe bubble chamber) [14].

The standard PWA (by construction) tends to miss a narrow resonance with $\Gamma_R < 20$ MeV or so [3]. The modified PWA assumes the existence of a narrow resonance by inserting a Breit–Wigner (BW) structure into one of the partial amplitudes and refitting the whole database. Such an analysis of the effect of the scanning mass M_R , full Γ_R , and partial width of the narrow resonance for each partial wave allows us to get a picture. The worse description of the full database that includes resonance with the corresponding M_R and Γ_R (in terms of $\Delta \chi^2$) does not support our hypothesis. The better description may result in the resonance that may exist, or the effect that can be due to various corrections (*e.g.*, thresholds), or both possibilities can contribute. Some additional checks are necessary: true resonance should provide the effect only in a particular partial wave, while a non-resonance source may show similar effects in various partial waves.

Using this kind of modified PWA, Arndt and co-workers performed a reanalysis of the existing KN database [2] to refit K^+N observables and consider the effect of a narrow state [3]. The existence of a Θ^+ in the P_{01} $(J^P = 1/2^+)$ state, with a mass of ~ 1545 MeV and a width of $\Gamma(\Theta^+) \leq 0.5$ MeV or so, was found to be possible (Fig. 2).

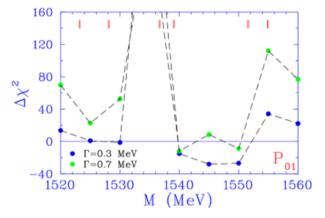


Fig. 2. Change of overall χ^2 due to insertion of a resonance into P_{01} for $M_R = 1520-1560$ MeV with $\Gamma_R = 0.3$ and 0.7 MeV, using the KN plane wave approximation [2]. Energies, where measurements exist, are marked by red vertical bars (for references, see [2]). The 1535 MeV values for $\Delta\chi^2$ are off-scale [3]. Lines are given to guide the eye.

3. My story of N(1680)

The mass difference in baryon multiplet is determined by the phenomenological Gell-Mann–Okubo mass formula (mixing is able to shift some masses). Originally, Mitya, Vitya, and Maxim claimed N(1710) as a non-strange member of $\overline{10}$, which was a suitable candidate for N^* [15]. The state N(1710), though listed in the PDG2002 baryon summary table [16] as a three-star resonance, was not seen in the PWA of the elastic scattering data πN by the SAID group [17].

Studies that have claimed to see this N(1710) state have given widely varying estimates of its mass (from 1680 MeV to 1740 MeV) and width (from 90 MeV to 500 MeV). Branching ratios have also been given with large uncertainties (10–20% for $N\pi$, 40–90% for $N\pi\pi$, etc.), apart from one that has been presented with greater precision ($6 \pm 1\%$ for $N\eta$) [16].

The π N PWA by the GW SAID group has shown that above the Roper resonance [18], W > 1500 MeV, $\sigma^{\text{tot}} \cong 2\sigma^{\text{el}} \cong 2\sigma^{\text{inel}}$ ($\sigma^{\text{tot}} = \sigma^{\text{el}} + \sigma^{\text{inel}}$). It means that it is almost pure diffraction, inelasticity $\eta \to \infty$ for P_{11} and $S \cong 0$ with $A \cong i/2$. The result is that phase δ is poorly defined and the P_{11} amplitude is spinning around Re A = 0 [19] (Fig. 3).

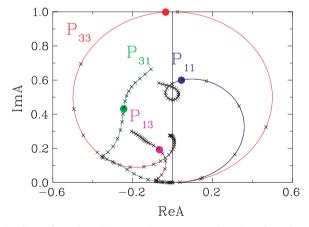


Fig. 3. Argand plots for the P partial-wave amplitudes for the πN case from threshold (1080 MeV) to W = 2500 MeV [19]. The crosses indicate 50 MeV steps in W. The solid circles correspond to the BW W_R .

Of course, the non-observation of a broad N(1710) state in πN elastic analyses (no pole position, no BW, no speed plot) could be due to a very small πN branching ratio. The standard procedure used in PWA may also miss narrow resonances with $\Gamma < 20$ MeV by construction [17]. Therefore, the non-strange true unitary partner of Θ^+ (if it is different from N(1710)and sufficiently narrow) could have eluded detection.

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Arndt and co-workers were considering the πN partial wave P_{11} , as this is the amplitude that is associated with resonances having $J^P = 1/2^+$ [20]. The character of χ^2 changes, $\Delta \chi^2$, after inserting a narrow resonance into the partial amplitude with a range of masses, widths, and branching fractions is illustrated in Fig. 4. Negative values of $\Delta \chi^2$ emerge more readily near $M_R = 1680$ MeV and 1730 MeV. We see that $\Delta \chi^2$ becomes negative only for $\Gamma_{\rm el} = (\Gamma_{\rm el}/\Gamma_{\rm tot}) \cdot \Gamma_{\rm tot}$ within the bounds of $\Gamma_{\rm el} \leq 0.5$ (0.3) MeV for $M_R =$ 1680 (1730) MeV (Fig. 4). The available data cannot reliably discriminate values of $\Gamma_{\rm el}$ below these bounds.

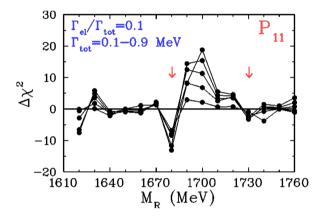


Fig. 4. Change of overall χ^2 due to insertion of a resonance into P_{11} for $M_R = 1660-1760$ MeV with $\Gamma_{\text{tot}} = 0.1$ (0.2) 0.9 MeV and $\Gamma_{\text{el}}/\Gamma_{\text{tot}} = 0.1$ using πN PWA [20]. Lines are given to guide the eye. The red vertical arrows indicate $M_R = 1680$ MeV and 1730 MeV.

The expected decay properties of N(1680) are essentially model-dependent and the Chiral Quark Soliton approach with violated $SU(3)_F$ [mixing $N_{\overline{10}} - N_8$] gives for $M_{N^*} = 1680$ (1730) MeV [20]:

- $\Gamma(N^* \to \pi \Delta) \approx 2.8 \ (3.5) \text{ MeV}.$
- $\Gamma(N^* \to K\Lambda) \approx 0.70 \ (1.56) \text{ MeV}.$
- $\Gamma(N^* \to \eta N) \approx 2$ MeV.
- $\Gamma(N^* \to \pi N) \approx 2.1$ (2.3) MeV.
- $\Gamma(N^* \to \text{tot}) \approx 10 \text{ MeV}.$

Then, the GRAAL Collaboration claimed evidence for a narrow (10 MeV) resonance state with a mass of 1675 MeV in the $\gamma n \rightarrow \eta n$ reaction (Fig. 5 (left)) [21]. By the way, in the case of a "neutron" target, the final-state interaction (FSI) is a problem. The effect is absent in the case of

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 $\gamma p \to \eta p$ (Fig. 5 (right)) and agrees with the prediction of the Chiral Quark Solution approach [22]. In exact SU(3)_F, the transition magnetic moment $\mu(p^* \to p)$ should vanish, since the U-spins are 3/2 for p^* , 1/2 for p, and 0 for the photon. With violation of SU(3)_F, this transition moment does not vanish but is still much smaller than $\mu(n^* \to n)$ [23].

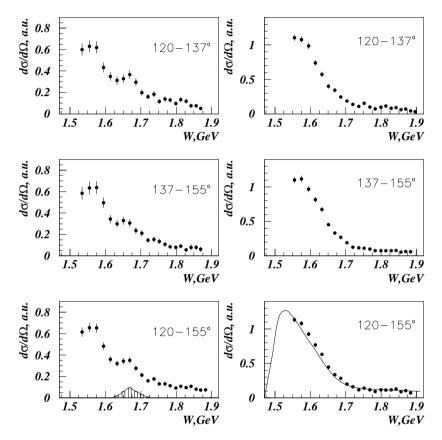


Fig. 5. Quasi-free ηn (left) and free ηp (right) photoproduction cross sections (dark circles) [21]. The dashed area (left plot) shows the simulated contribution of a narrow state at W = 1.675 MeV. The solid line (right plot) indicates the solution of E429 of the SAID $\gamma p \rightarrow \eta p$ PWA [24].

The GW SAID group performed a coupled channel analysis of the πN system that included elastic scattering of $\pi^{\pm}p$ and $\pi^{-}p \to \pi^{0}n$ with $\pi^{-}p \to \eta n$ [17]. The attractive factor for the $\pi^{-}p \to \eta n$ reaction is that it plays a role in isospin filtering. Unfortunately, the $\pi^{-}p \to \eta n$ data above 800 MeV (W = 1630 MeV) are not reliable for PWA (see, for instance, Fig. 6 (left)). Most of the NIMROD data do not satisfy requirements (systematics is 10%

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or more, momentum uncertainties are up to 100 MeV/c, and so on). That is the main reason why the SAID coupled channel analysis had no opportunity to look for N(1680) using the $\pi^- p \to \eta n$ reaction data (Fig. 6 (right)).

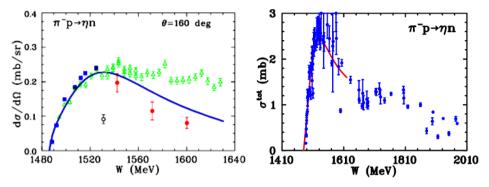


Fig. 6. Left: Excitation function for differential cross sections for the $\pi^- p \to \eta n$ reaction. Data came from BNL, RHEL, and Saclay. The solid blue line is from SAID PWA (WI08 solution) [26]. Right: Total cross sections for the $\pi^- p \to \eta n$ reaction. The solid red line is from SAID PWA [26]. All data are available in Ref. [25].

The EPECUR Collaboration at ITEP performed high quality measurements of the differential cross sections for elastic $\pi^{\pm}p$ over an energy range of p = 800-1300 MeV/c (W = 1.55-1.83 GeV) and for center-of-mass angles from 40 to 120 deg [27]. In total, approximately 10,000 new data points have been obtained. These data have been produced with a momentum resolution of ~ 1 MeV/c and with ~ 1% statistical uncertainties.

Following the analysis [20], $\Gamma(\pi N) \sim 0.5$ MeV, which stimulated us to look for N(1680) using high-quality EPECUR elastic $\pi^{\pm}p$ scattering data [28]. The bottom line is that two narrow structures observed in elastic $\pi^{-}p$ scattering can be explained by a combination of threshold effects and two narrow resonances, $S_{11}(1686)$ and $P_{11}(1720)$. Figure 7 presents the predictions for the $\pi^{-}p \to \eta n$ total cross sections. This model-dependent analysis shows that $S_{11}(1686)$ cannot be a member of $\overline{10}$, while $P_{11}(1720)$ may have a link to the results of Arndt *et al.* [20].

The new experimental proposal P102 for J-PARC addresses a measurement of the cross section of the $\pi^- p \to \eta n$ reaction at p = 850-1200 MeV/cwith a CsI (Tl) calorimeter [29]. The new data will help us solve the puzzle of N(1680).

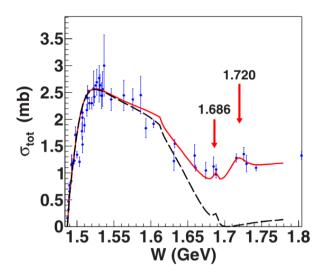


Fig. 7. Comparison of the measured πN data [25] and calculated $\pi^- p \to \eta n$ total cross sections. The solid red curve presents the calculations from the work [28]. The black dashed curve indicates the contribution of the S-wave.

4. Conclusion

What would Mitya, Vitya, and Maxim think of the current situation in particle physics? On the one hand, they would be justifiably proud of the robust successes of an instanton-vacuum approach, which did so much to place it on a firm footing. However, some of the problems that concern them remain unresolved, notably the exotic Θ^+ and its partners.

In fact, twenty years ago, in 2003, two experimental groups, LEPS and DIANA, announced the observation of a light and narrow exotic baryon with a mass in the vicinity of 1540 MeV, which was predicted by our friends. The history of this discovery and its theoretical interpretations can be found in the recent review [30].

Often memory does not keep details, and let me remind two events associated with Mitya, Vitya, and Maxim and Θ^+ ...

 In 2009, I organized a Workshop on Narrow Nucleon Resonances: Predictions, Evidences, Perspectives, Edinburgh, Scotland, June, 2009. https://gwdac.phys.gwu.edu/~igor/Edinburgh2009/index.htm. Unfortunately all three friends (Maxim was one of organizers) missed this event for some reasons, but many people attended and discussed physics. Finally, a photo taken in Maxim's apartment in Bochum, Germany, back to 2006 (Fig. 8).



Fig. 8. Sitting on sofa from left to right: Masha, Maxim's daughter, me, Maxim, Takashi Nakano. Slava Kuznetsov is sitting just in front of us, he was a chef of our dinner.

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Appendix

My archive kept a letter for the name Θ^+ ' that Mitya sent to the community [32]. The result is that the community accepted Mitya's suggestion.

From: Dmitri Diakonov To: Takashi Nakano, Anatoly Dolgolenko, Elton Smith, Valeri Koubarovski, Ken Hicks, Eugene Pasyuk,

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Stepan Stepanyan, Mark Strikman, Maxim Polyakov, Victor Petrov, Igor Strakovsky, James Bjorken, Stan Brodsky, Lev B. Okun, Leonid G. Landsberg, Hiroshi Toki. Date: Apr 12, 2003, 12:33PM Subject: exotic baryon name

Dear friends,

Now that, thanks to your efforts, the exotic S = +1 baryon seems to become gradually a reality and not just a fancy dream of a theorist, we have to think how to name the child to be born, hopefully.

The tentative name Z^{t} we all have used is, to my mind, unfortunate. First, I have noticed that the first reaction by people hearing about it is that the new particle has some association with the electroweak Z boson. This association is completely misleading and may create lots of confusion in future and cause us much trouble. Second, the tradition is that baryons are named by capital Greek letters: Delta, Lambda, Sigma, Xi, Omega. (Except the nucleon but one cannot change it.) Third, "Z" is pronounced very differently in different languages; even in English it is "zee" in American English and "zet" in British which is mostly used in Europe and elsewhere. Fourth, and probably most important, the name "Z" was imposed on us by the unsuccessful searches of exotic baryons in the far higher mass region, as summarized in the 1986 edition of Particle Data. This was "prehistoric" time, before the discovery of the Z boson, and in any case nothing has been found worth mentioning in Particle Data Listings since the 1986 edition. The 1540 MeV signal has nothing to do with the broad and always suspicious one-star candidates in the 1700 MeV region. I regret that we uncritically and unthoughtfully borrowed this queer name and used it in our paper with Petrov and Polyakov. But at that time it was so far from experiment!

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I suggest that the name for this baryon should satisfy the following criteria:

It must be a capital Greek letter, according to the tradition of naming baryons.
It must be distinct from anything used before and carry no associations with bosons.
Last but not least, the character must exist in LaTeX.

If you look into the list of upper case Greek letters used in LaTeX you'll find that there is one and only one character satisfying all criteria, and it is

THETA.

When I discovered it I realized that I kind of liked this name. It is a symmetric and "round" character like the Omega also sitting at the vertex of the big decuplet triangle and, like Omega, alluding to that it is a singlet. If you disregard the historic "tau-theta problem" of the last century's 50s, this character has been never used in particle physics, as far as I know. It carries no associations and hints that it is something really new, which it certainly is. Last, "Theta" is pronounced more or less in the same way in all languages I know.

Therefore, I suggest that we, experimentalists and theorists, will henceforth call the new candidate for the exotic baryon with strangeness +1 and mass around 1540 MeV.

 Θ^+ baryon.

Any objections?

With my best regards, Dmitri Diakonov.

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