PHENIX RESULTS ON COLLISION ENERGY-DEPENDENT LÉVY HBT CORRELATIONS FROM $\sqrt{s_{NN}} = 15$ TO 200 GeV^{*}

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Different regions on the QCD phase diagram can be investigated by varying the collision energy and the centrality in heavy-ion collisions. In our latest measurements at the PHENIX experiment at RHIC, we utilize Lévy-type sources to describe the measured HBT correlation functions. In this paper, we report on the current status of the analysis of the centrality and beam energy dependence of the Lévy source parameters in Au+Au collisions from $\sqrt{s_{NN}} = 15$ GeV to $\sqrt{s_{NN}} = 200$ GeV.

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1. Introduction

To experimentally explore the phase diagram of strongly interacting matter, one has to investigate heavy-ion collision data from different center-ofmass collision energies. The aim of the Beam Energy Scan program at the Relativistic Heavy Ion Collider (RHIC) was to collect a huge variety of data from different energies and different kinds of colliding nuclei. The second phase of this program, the BES-II, is still ongoing. Complemented with the fixed-target program at STAR, the beam energy scan will cover a wide range in baryochemical potential, spanning from about $\mu_{\rm B} = 20$ MeV up to $\mu_{\rm B} = 720$ MeV [1]. For the analysis reported in this paper, we used BES-I data recorded by the PHENIX experiment at six different collision energies. Figure 1 illustrates the QCD phase diagram and the data recorded by PHENIX during BES-I. The datasets used in our analysis are highlighted in the table.

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Fig. 1. QCD phase diagram and the RHIC Beam Energy Scan.

Having a huge variety of data is necessary but not sufficient for exploring the phase diagram. We also need experimental observables that provide us with information on the nature of the phase transition. One of these useful observables is the quantum-statistical (also called Bose–Einstein or HBT) correlations of identical bosons. With the help of these correlation functions, we can gain information on the space-time structure of the particle-emitting source in heavy-ion collisions. To describe the measured correlation functions, one has to utilize a certain type of source function. In the previous years, the usual practice was to use a Gaussian-type of source but the latest results from different experiments (PHENIX [3], CMS [4]) showed that we need to go beyond the Gaussian approximation. Different scenarios (e.q.anomalous diffusion) can lead to the appearance of Lévy-type sources, as described in Refs. [3, 5]. The symmetric Lévy distribution is the generalization of the Gaussian, and it is characterized by two main parameters, the Lévy scale R and the Lévy exponent α . It is defined by the following expression:

$$\mathcal{L}(\alpha, R, \boldsymbol{r}) = \frac{1}{(2\pi)^3} \int \mathrm{d}^3 \boldsymbol{q} e^{i\boldsymbol{q}\boldsymbol{r}} e^{-\frac{1}{2}|\boldsymbol{q}R|^{\alpha}} \,. \tag{1}$$

The R parameter is in connection with the physical size of the source, and α parameter determines the shape. If the Lévy exponent takes the value of 2, we get back the Gaussian distribution, and in the case of $\alpha < 2$, the Lévy distribution exhibits a power-law behavior. If we neglect the final-state interactions and assume a spherically symmetric Lévy source, the correlation function takes the simple form of

$$C(Q) = 1 + \lambda \exp\left(-(RQ)^{\alpha}\right).$$
⁽²⁾

In our analysis, we defined the one-dimensional relative-momentum variable Q as the absolute value of the three-momentum difference in the LCMS frame, see details in [3]. In the framework of the Core–Halo model [6], the λ parameter (the strength of the correlation function) is in connection with the core–halo ratio

$$\lambda = (N_{\rm C} / (N_{\rm C} + N_{\rm H}))^2 .$$
(3)

In this framework, we split the source function into two parts. The core part contains the pions created directly from the freezeout (or from very short lived resonances), while the halo part contains the pions created in decays of long-lived resonances. The number of pions associated with the core and halo part are denoted by $N_{\rm C}$ and $N_{\rm H}$, respectively.

Besides describing data better than the Gaussian, a Lévy-type source has further importance in the search for the critical point. It can be shown that the Lévy exponent α parameter is connected to the critical exponent η known from statistical physics. In the case of a second order phase transition at the critical point, the spatial correlations are characterized by a powerlaw behavior with an exponent $-(d-2+\eta)$, where d is the dimension of the system. The power-law tail of a three-dimensional Lévy distribution is described with an exponent of $-1 - \alpha$, so just by simply comparing the two cases, we can see that in the case of three dimensions (and assuming infinite volume nuclear matter), the exponents are identical at the critical point. If we assume that the universality class of QCD is identical to that of the three-dimensional Ising model (or the random-field 3D Ising model), we can look up the value of the η exponent from the literature and find the rather small value of $\eta = 0.03631(3)$ [7] (or $\eta = 0.5 \pm 0.05$ in the case of the random-field 3D Ising model [8]). If the previous arguments about the connection between η and α are true, with the measurement of α at different center-of-mass collision energies, we may get closer to being able to locate the critical point on the phase diagram of strongly interacting matter.

2. Centrality dependence of the source parameters in Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV

It has been a matter of discussion for a while now that finite-size and finite-time effects play an important role in the phase transitions that occur in heavy-ion collisions [9]. It is obvious that we cannot talk about infinite volume nuclear matter in the case of such collisions, but we may be able to control the system size (with the help of the centrality of these collisions). For the previously mentioned reasons, it is important to investigate in detail the centrality dependence of the extracted source parameters. In this section, we present such an analysis in $\sqrt{s_{NN}} = 200 \text{ GeV Au}+\text{Au}$ collisions (recorded by PHENIX during the 2010 running period).

In our analysis, we measured one-dimensional two-pion HBT correlation functions (for both $\pi^+\pi^+$ and $\pi^-\pi^-$ pairs). After the experimental construction of the correlation functions with the help of the usual event-mixing method, we confirmed that the $\pi^+\pi^+$ and $\pi^-\pi^-$ data are consistent with each other, as expected theoretically. We combined the two in order to increase statistical precision. We used 17 different average transverse mass $(m_{\rm T})$ bins and 4 centrality intervals with 10% increments. The previously shown formula (Eq. (2)) does not take into account final-state interactions. In the fitting process, we incorporated the final-state Coulomb effect in our fit function, and introduced a linear background (which was consistent with a constant in most cases). The extracted fit parameters are shown in Fig. 2.



Fig. 2. (Color online) Centrality and $m_{\rm T}$ dependence of Lévy source parameters in $\sqrt{s_{NN}} = 200$ GeV Au+Au collisions. The different shades of gray/colors and marker styles are denoting the different centrality ranges. The auxiliary figures at the bottom show the relative systematic uncertainties.

The Lévy scale parameter, R, shows the distinctive features of a geometrical centrality dependence — its value decreases as we go from central to peripheral events. It also shows the usual decrease as we go to higher average transverse mass values, which is consistent with expectations from hydrodynamics. The strength of the correlation function, λ , decreases at small $m_{\rm T}$, which could be a possible sign of in-medium mass modification, for details, see Refs. [3, 10]. The Lévy exponent α parameter is between 0.5 and 2 in all centrality and $m_{\rm T}$ bins, so it is clearly far from both the Gaussian ($\alpha = 2$) and the conjectured CEP value ($\alpha \leq 0.5$ [11]).

During the fitting of the correlation functions, we observed that the parameters are strongly correlated, which motivated us to look for less correlated combinations. We indeed found such a combination empirically, defined as $\hat{R} = R/(\lambda (1+\alpha))$. This parameter is less correlated with the others, and it can be measured much more precisely. What is quite fascinating is the linear dependence of $1/\hat{R}$ on $m_{\rm T}$. The possible physical interpretation of this observation is still unclear. More details on the centrality dependence of the parameters can be found in Refs. [12, 13].

3. Excitation function of the source parameters

In the case of the $\sqrt{s_{NN}} = 200$ GeV data, the investigation of detailed centrality and $m_{\rm T}$ dependence becomes possible due to the adequate statistical precision. However, when we go down in energy, doing a similar analysis can become quite challenging due to the low statistics. To overcome this



Fig. 3. (Color online) Excitation functions of Lévy source parameters in Au+Au collisions. The error bars represent the statistical uncertainties, while the gray/green boxes show the systematic uncertainties.

issue, we chose to concentrate only on the excitation function of the parameters, and used only one fairly wide centrality range (0–30%) and just one $m_{\rm T}$ bin. At lower energies, we used wider $m_{\rm T}$ ranges while keeping the mean value, $\langle m_{\rm T} \rangle$ the same. This made it possible to perform our measurements on six different energies: 200, 62, 39, 27, 19, and 15 GeV. We observed that the choice of the $m_{\rm T}$ bin width can be a non-negligible source of systematic uncertainty. The results on the excitation functions of the Lévy source parameters can be seen in Fig. 3.

One main observation that we can make is that all the parameters show a weak non-monotonicity with collision energy. It is important to note, however, that the statistical and systematic uncertainties become quite large at lower energies. It is also important to note that the Lévy scale α parameter is above 1 at the whole energy range so we are still far from the conjectured critical point value. Among all parameters, the previously discussed scaling parameter \hat{R} shows the statistically most significant change with energy.

4. Summary and outlook

In these proceedings, we presented Lévy-type HBT measurements at PHENIX. We discussed the detailed centrality and $m_{\rm T}$ dependence of the Lévy source parameters in the case of the $\sqrt{s_{NN}} = 200$ GeV data, as well as the excitation functions of the parameters in one given $m_{\rm T}$ and centrality range. The shown results are still preliminary and more detailed investigations are currently ongoing. In the future, we also plan to finalize other Lévy HBT results on three-dimensional as well as three-particle correlation analysis.

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