# LIGHT-BY-LIGHT SCATTERING CROSS-SECTION MEASUREMENTS AT THE LHC\*

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Received 1 August 2022, accepted 28 September 2022, published online 14 December 2022

Light-by-light (LbyL) scattering,  $\gamma \gamma \rightarrow \gamma \gamma$ , is a rare Standard Model (SM) process, also proposed as a sensitive channel to study physics beyond the SM. In these proceedings, we perform a statistical combination of existing  $\gamma \gamma \rightarrow \gamma \gamma$  cross section measurements at the LHC with the aim of checking the consistency with different SM predictions. Using a simplified

<sup>\*</sup> Presented at the 29<sup>th</sup> International Conference on Ultrarelativistic Nucleus–Nucleus Collisions: Quark Matter 2022, Kraków, Poland, 4–10 April, 2022.

set of assumptions, we find the averaged result of  $115 \pm 19$  nb, consistent with SM predictions within two standard deviations. For the first time, we also consider the contribution from the  $\eta_b(1S)$  meson production to the diphoton invariant mass distribution.

DOI:10.5506/APhysPolBSupp.16.1-A123

# 1. Introduction

In these proceedings, we present an average of the integrated fiducial cross sections for the Pb+Pb  $(\gamma\gamma) \rightarrow Pb^{(*)}+Pb^{(*)}\gamma\gamma$  production at  $\sqrt{s_{_{NN}}} = 5.02$  TeV using the available LHC measurements from the ATLAS [1–3] and CMS [4] collaborations.

### 2. Measured and predicted light-by-light cross sections

In the offline ATLAS (CMS) analysis used as input measurements, events were selected with exactly two photons, each with transverse energy  $E_{\rm T} >$ 2.5 (>2.0) GeV and pseudorapidity  $|\eta| < 2.4$  that satisfied further selection requirements described in Refs. [3, 4]. To suppress nonexclusive background, characterized by a final state with larger transverse momenta,  $p_{\rm T}$ , and diphoton acoplanarities,  $A_{\phi} = (1 - |\Delta \phi^{\gamma\gamma}|/\pi)$ , the  $p_{\rm T}$  and  $A_{\phi}$  were required to be  $p_{\rm T}^{\gamma\gamma} < 1$  GeV and  $A_{\phi} < 0.01$ , respectively, and the invariant mass of the pair to be  $m_{\gamma\gamma} > 5$  GeV. The cross section for the Pb + Pb  $(\gamma\gamma) \rightarrow$  Pb<sup>(\*)</sup>+Pb<sup>(\*)</sup>  $\gamma\gamma$  process was measured in fiducial phasespace regions, defined by the above requirements, and reflecting the selection at the reconstruction level. A summary of the available measurements along with their total uncertainty, evaluated as the quadratic sum of the individual sources, is presented in Table 1.

The theoretical predictions for the Pb + Pb  $(\gamma\gamma) \rightarrow Pb^{(*)}+Pb^{(*)}\gamma\gamma$  production cross section at  $\sqrt{s_{_{NN}}} = 5.02$  TeV are calculated numerically at leading order (LO) in SuperChic [5]. An alternative LbyL numerical calculation is performed in Ref. [6], with the main difference of 2–4% relative to SuperChic originating from the implementation of the nonhadronic-overlap condition of the Pb ions. In either case, next-to-leading-order (NLO) QCD and QED corrections [7, 8] increase the cross section by a few percent, and are taken into account in the quoted total uncertainties of 10%. A summary of theoretical cross-section predictions with their uncertainties is shown in Table 2, separately for the phase-space regions defined in Table 1. Based on these numbers, the ratios of 77/101  $\approx$  0.76 and 77/50  $\approx$  1.54 are used as correction factors accounting for differences in the definition of phase-space regions between the two experiments.

Table 1. Summary of the fiducial LbyL cross-section measurements ( $\sigma_{\rm raw}^{\rm fid}$ ) at 5.02 TeV performed by the ATLAS and CMS collaborations. When applicable, they are further scaled by correction factors ( $\sigma_{\rm cor}^{\rm fid}$ ) to account for differences in the definition of phase-space regions, as described in the main text. Total uncertainties are shown. The symbol "—" means that no corresponding cross-section measurement currently exists. The cross sections marked with <sup>†</sup> are those used as input to the extraction of the averaged value of the Pb+Pb ( $\gamma\gamma$ )  $\rightarrow$  Pb<sup>(\*)</sup>+Pb<sup>(\*)</sup>  $\gamma\gamma$  process.

		ATLAS		CMS	
$\sqrt{s_{_{NN}}}$	Year (Lumi. $[nb^{-1}]$ )	$\sigma_{\rm raw}^{\rm fid}$ [nb]	$\sigma_{\rm cor}^{\rm fid}$ [nb]	$\sigma_{\rm raw}^{\rm fid}$ [nb]	$\sigma_{\rm cor}^{\rm fid}$ [nb]
$5.02\mathrm{TeV}$	2015 (0.39 - 0.48)	$70 \pm 29$ [1]	$108\pm45$	$120 \pm 55$ [4]	$91\pm42^\dagger$
	2018(1.73)	$78 \pm 15$ [2]	$120\pm23$		
	$2015{+}2018~(2.2)$	$120 \pm 22$ [3]	$120\pm22^{\dagger}$		

Table 2. Predicted cross sections for LbyL scattering at 5.02 TeV. Uncertainties take into account derivations from alternative approaches. The cross section marked with <sup>†</sup> is used as reference.

$\sqrt{s_{_{NN}}}$	Process	Accuracy	$\sigma_{\rm theo}^{\rm fid} [{\rm nb}]$	Phase-space region
$5.02\mathrm{TeV}$	$Pb+Pb(\gamma\gamma) \rightarrow Pb^{(*)}+Pb^{(*)}\gamma\gamma$	LO	$101 \pm 10$ [5]	$E_{\rm T} > 2.0 \text{ GeV},  \eta  < 2.4, m_{\gamma\gamma} > 5 \text{ GeV}, p_{\rm T}^{\gamma\gamma} < 1 \text{ GeV}, A_{\phi} < 0.01$
		LO	$103 \pm 10$ [6]	$E_{\rm T} > 2.0 \text{ GeV},  \eta  < 2.4, m_{\gamma\gamma} > 5 \text{ GeV}, p_{\rm T}^{\gamma\gamma} < 1 \text{ GeV}, A_{\phi} < 0.01$
		LO	$77 \pm 8^{\dagger}$ [5]	$E_{\rm T} > 2.5 \text{ GeV},  \eta  < 2.4, m_{\gamma\gamma} > 5 \text{ GeV}, p_{\rm T}^{\gamma\gamma} < 1 \text{ GeV}, A_{\phi} < 0.01$
		LO	$80 \pm 8$ [6]	$E_{\rm T} > 2.5 \text{ GeV},  \eta  < 2.4, m_{\gamma\gamma} > 5 \text{ GeV}, p_{\rm T}^{\gamma\gamma} < 1 \text{ GeV}, A_{\phi} < 0.01$
		LO	$50 \pm 5$ [5]	$E_{\rm T} > 3.0 \text{ GeV},  \eta  < 2.4, m_{\gamma\gamma} > 6 \text{ GeV}, p_{\rm T}^{\gamma\gamma} < 1 \text{ GeV}, A_{\phi} < 0.01$
		LO	$51 \pm 5$ [6]	$E_{\rm T} > 3.0 \text{ GeV},  \eta  < 2.4, m_{\gamma\gamma} > 6 \text{ GeV}, p_{\rm T}^{\gamma\gamma} < 1 \text{ GeV}, A_{\phi} < 0.01$

The background contribution of  $\eta_b(1S)$  meson production to the Pb + Pb  $(\gamma\gamma) \rightarrow Pb^{(*)}+Pb^{(*)}\gamma\gamma$  process is studied for the first time. Figure 1 shows a comparison of the theoretical calculations [6] of differential cross sections as a function of the invariant mass of the diphoton system for  $\gamma\gamma \rightarrow \gamma\gamma$  scattering, typically used to set limits on particle production beyond the Standard Model, and  $\eta_b(1S)$  meson production with decays into two photons in the final state. Although the height of the resonance peak is 1 nb, this contribution is therefore found to be insignificant in the context of the LbyL measurement because the experimental  $m_{\gamma\gamma}$  bins are typically very wide.



Fig. 1. Differential cross sections for the signal (contribution from fermionic boxes [6]) Pb+Pb  $(\gamma\gamma) \rightarrow$  Pb<sup>(\*)</sup>+Pb<sup>(\*)</sup>  $\gamma\gamma$  and the intermediate  $\gamma\gamma \rightarrow \eta_b(1S) \rightarrow \gamma\gamma$ background production process as a function of the diphoton invariant mass. For the signal, the ATLAS and CMS kinematic requirements from Refs. [3] and [4], respectively, are adopted. For the background  $\eta_b(1S)$  process, the decay to a diphoton system is shown for the maximum and minimum values of diphoton decay rates of 0.46 and 0.17 keV, respectively.

# 3. Averaged cross-section measurement at the LHC

The cross-section measurements, described in Section 2 and denoted by <sup>†</sup> in Table 1, are used as input to an averaged cross section. We use the best linear unbiased estimator (BLUE) method to average these measurements. Systematic uncertainties are categorized and a simplified correlation scheme is assumed [9].

The averaged cross-section measurement for the LbyL process at 5.02 TeV is

$$\begin{aligned} \sigma_{\text{meas}}^{\text{fid}} &= 115 \pm 15 \; (\text{stat.}) \pm 11 \; (\text{syst.}) \pm 3 \; (\text{lumi.}) \pm 3 \; (\text{theo.}) \; \text{nb} \\ &= 115 \pm 19 \; \text{nb} \; , \end{aligned}$$

with a relative uncertainty of 17%. The statistical uncertainty is still found to be the dominant overall uncertainty. Figure 2 shows a summary of the Pb + Pb  $(\gamma\gamma) \rightarrow$  Pb<sup>(\*)</sup>+Pb<sup>(\*)</sup>  $\gamma\gamma$  measurements at 5.02 TeV and their comparison to the theory predictions from Table 2. The averaged cross section is consistent within about two standard deviations with the SM predictions. A related analysis [10] has shown that the recently observed tetraquark state X(6900) could, in principle, account for the excess seen in the Pb + Pb  $(\gamma\gamma) \rightarrow$  Pb<sup>(\*)</sup>+Pb<sup>(\*)</sup>  $\gamma\gamma$  data.



Fig. 2. The averaged Pb+Pb  $(\gamma\gamma) \rightarrow Pb^{(*)}+Pb^{(*)}\gamma\gamma$  cross-section value along with the individual cross-section measurements at 5.02 TeV from ATLAS and CMS. The theoretical predictions [5, 6] are computed at LO accuracy. The  $\sigma_{\text{theo}}^{\text{fid}}$  uncertainties used to compute Pb + Pb  $(\gamma\gamma) \rightarrow Pb^{(*)}+Pb^{(*)}\gamma\gamma$  are described in the text.

#### 4. Summary

Although an improved determination of the integrated fiducial Pb +  $Pb(\gamma\gamma) \rightarrow Pb^{(*)}+Pb^{(*)}\gamma\gamma$  cross section by approximately 10% could be potentially achieved relative to current measurements, further improvements are expected with the inclusion of existing or forthcoming LHC nuclear data.

This article is part of a project that has received funding from the European Union's Horizon 2020 research and innovation program under grant agreement No. 824093, better known as STRONG-2020. G.K.K. is supported by the Office of Nuclear Physics in the Department of Energy (DOE NP) under grant number DE-FG02-96ER40981. I.G.B. is supported by the National Science Centre, Poland (NCN) under grant number UMO-2020/37/B/ST2/01043 and by PL-GRID infrastructure.

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