# MEASUREMENTS OF CHARGE-DEPENDENT CORRELATIONS WITH CMS\*

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Charge-dependent azimuthal anisotropy Fourier coefficients are measured with two- and three-particle correlations in pPb and PbPb collisions. The difference between positively- and negatively-charged particles for the second-order two-particle,  $v_2\{2\}$ , and three-particle,  $v_2\{3\}$ , coefficients for both pPb and PbPb, and third-order two-particle coefficient,  $v_3\{2\}$ , for PbPb, are presented. The observed results are challenging the hypothesis that attributes the charge-dependent azimuthal correlations in heavy-ion collisions to the chiral magnetic effect. In addition, the two-particle electric charge balance function is used as a probe to study the charge creation mechanism in high-energy heavy-ion collisions, for the first time in CMS. The balance function is constructed using like and unlike charged-particle pairs. The width of the balance function, both in relative pseudo-rapidity and relative-azimuthal angle, increases from more central collisions to peripheral ones. Narrowing and widening of these widths indicate late and early hadronization, respectively.

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## 1. Introduction

The ultra-relativistic nucleus–nucleus (AA) collisions can be used as a powerful tool to investigate the properties of the deconfined state of matter, known as quark–gluon plasma (QGP) [1, 2]. In quantum chromodynamics (QCD), the interactions of chiral quarks with the gluonic fields can produce a chiral imbalance due to the quantum anomaly by violating the local P and CP symmetries [3]. The strong magnetic field created from non-central heavy-ion collisions induces a current, and the electric charge separation is known as the chiral magnetic effect (CME) [4]. Also, a similar process due

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to the chiral charge separation along the magnetic field is known as the chiral separation effect (CSE). The combined effect of CME and CSE leads to a long-wavelength collective excitation, known as a chiral magnetic wave (CMW) [5]. There will be a development of electric quadrupole moment due to the CMW, resulting in the positive (negative) charges accumulated away from (close to) the reaction plane. The charge-dependent variation of Fourier azimuthal anisotropy  $(v_2)$  is calculated for the positive and negative particles and a linear dependence on the observed event charge asymmetry,  $A_{\rm ch} = (N_+ + N_-)/(N_+ - N_-)$ .  $N_+$  and  $N_-$  are the number of positively- and negatively-charged particles in each event. The elliptic azimuthal anisotropy as a function of charge asymmetry can be written as

$$v_{2,\pm} = v_{2,\pm}^{\text{base}} \mp r A_{\text{ch}} \,.$$
 (1)

In Eq. (1),  $v_{2,\pm}^{\text{base}}$  represents the value in the absence of a charge quadrupole moment from the CMW, and r is the slope parameter. The  $A_{ch}$  value is also corrected due to the detector inefficiency and acceptance, called true event charge asymmetry  $(A_{ch}^{\text{true}})$ .

## 2. Results

This article presents the charge-dependent azimuthal anisotropy in pPb and PbPb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV using the CMS detector [6].

The  $v_2$  of the positively- and negatively-charged hadrons are calculated using the two-particle cumulant method [7]. We observed a significant nonzero  $v_2$ , and the values are increasing  $(h^-)$  or decreasing  $(h^+)$  with  $A_{\rm ch}$  [8]. Figure 1 (right plot) shows the  $\langle p_{\rm T} \rangle$  as a function of charge asymmetry, we observed similar behaviour in average  $p_{\rm T}$  as  $v_2$ .

The normalized slope parameter  $r^{\text{norm}}$  is calculated by fitting a linear function for the  $v_2$  and average  $p_{\text{T}}$  as shown in figure 2. This behaviour is qualitatively consistent with the expectation of the local charge conservation (LCC) effect [9] from resonance decays. Since the  $v_n$  has a strong dependence on particle  $p_{\text{T}}$ , a correlation between the  $p_{\text{T}}$ -averaged  $v_n$  and  $A_{\text{ch}}$  can also be induced by the LCC mechanism.

The charge-dependent third-order harmonic coefficients were measured in PbPb collisions, and their magnitude was smaller than that of the secondorder coefficient [8]. Figure 3 represents the normalized difference in  $v_n$ , for n = 2 and n = 3, as a function of true event charge asymmetry. The normalized slope parameter of  $v_3$ ,  $r_3^{\text{norm}}$ , agrees well with  $r_2^{\text{norm}}$  within statistical uncertainties.



Fig. 1. The left plot shows the  $v_2$  as a function of  $A_{ch}$  and the right plot shows the average  $p_T$  with charge asymmetry for the same and opposite charge pairs in *p*Pb (top plot) and PbPb (bottom plot), respectively.



Fig. 2. The left plot shows the normalized  $v_2$  as a function of  $A_{ch}$  and the right plot shows the normalized average  $p_T$  with charge asymmetry in *p*Pb and PbPb collisions, respectively.

We measured the  $r_2^{\text{norm}}$  and  $r_3^{\text{norm}}$  as a function of centrality in PbPb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV. Over all the centrality ranges, the slope parameter is found consistent within the statistical uncertainties. The CMW effect is expected to be approximated by the second-order event plane with respect to the reaction plane in nucleus–nucleus collisions but highly suppressed for the third-order event plane [10]. The values for  $r_2^{\text{norm}}$  and  $r_3^{\text{norm}}$  are independent and similar, and this indicates that an underlying physics mechanism is unrelated to the CMW effect and, instead, can be qualitatively explained by the LCC effect.



Fig. 3. The normalized difference of  $v_n$  for the same and opposite charge pairs,  $(v_n^+ - v_n^-)/(v_n^+ + v_n^-)$  in PbPb collisions in 30–40% centrality.



Fig. 4. The linear slope parameters,  $r_2^{\text{norm}}$  and  $r_3^{\text{norm}}$ , as functions of the centrality class in PbPb collisions. Average  $N_{\text{trk}}^{\text{offline}}$  values for each centrality class are indicated on the top axis. Statistical and systematic uncertainties are indicated by the error bars and shaded regions, respectively.

Figure 5 shows the normalized slope parameters for  $v_2$  and  $\langle p_T \rangle$  in *p*Pb and PbPb collisions, respectively. We observed that normalized slope parameter  $r^{\text{norm}}$  values, for the both  $v_2$  and  $\langle p_T \rangle$  are found to have a weak dependence on the event multiplicity for both *p*Pb and PbPb collisions, with values for  $\langle p_T \rangle$  approximately half of those for  $v_2$ . Normalized slope parameters are larger in *p*Pb than PbPb collisions for the overlapping multiplicity region, which is not expected in the CMW context and may indicate a collision system dependence on the LCC.



Fig. 5. The linear slope parameters,  $r_2^{\text{norm}}$  and  $r_{\langle p_T \rangle}^{\text{norm}}$ , as functions of event multiplicity in *p*Pb and PbPb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV. Statistical uncertainties are indicated by the error bars and systematic uncertainties are indicated by shaded regions.

## 3. Summary

In this article, we discussed the charge-dependent Fourier coefficients of the azimuthal anisotropy as a function of the charge asymmetry in *p*Pb and PbPb collisions at  $\sqrt{s_{NN}} = 5.02$  TeV. A significant value of normalized slope parameters of  $v_2$  coefficient is observed in *p*Pb collisions, as well as in PbPb collisions over a wide range of charged particle multiplicities. A similar slope parameter for  $v_2$  and  $v_3$  coefficients are observed for the various centrality in PbPb collisions. A significant value of event-averaged transverse momenta of positively- and negatively-charged particles is observed in both *p*Pb and PbPb collisions. The above-presented results using the CMS phase-space window for  $\sqrt{s_{NN}} = 5.02$  TeV, are not expected from the chiral magnetic wave as the dominant physics mechanism. In addition, they are qualitatively consistent with predictions based on the LCC effect model. This indicates that the CMW is not the main source for the charge-dependent azimuthal anisotropies found in *p*Pb and PbPb collisions.

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