# HIGH $p_{\rm T}$ SUPPRESSION OF $\Lambda$ AND $K_{\rm s}^0$ IN Pb–Pb COLLISIONS AT $\sqrt{s_{NN}} = 2.76$ TeV WITH ALICE\*

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The nuclear modification factors  $R_{AA}$  and  $R_{CP}$  of  $\Lambda$  and  $K_s^0$  in Pb–Pb collisions at  $\sqrt{s_{NN}} = 2.76$  TeV measured by the ALICE experiment at the LHC are presented. In central collisions a strong suppression at high  $p_{\rm T}$  ( $p_{\rm T} \sim 8 \text{ GeV}/c$ ) with respect to pp collisions is observed similar for  $\Lambda$  and  $K_s^0$ . The  $p_{\rm T}$  region below is dominated by an enhancement of  $\Lambda$  over the suppressed  $K_s^0$ . The results are compared to those for charged particles and to  $\Lambda$  from lower collision energies.

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# 1. Introduction

The comparison of transverse momentum spectra from AA collisions to those from pp collisions is of importance for understanding the parton energy loss in the medium created in heavy-ion collisions. Especially, high momentum hadrons are diagnostic means to probe the medium before freezeout via their energy loss due to in-medium interactions since they are expected to be only weakly affected by flow. The relative energy loss can be quantitatively expressed by so-called nuclear modification factors  $R_{AA}$  (Eq. (1a)) and  $R_{CP}$ (Eq. (1b)) [1]

$$R_{AA} = \frac{1}{\langle N_{\rm bin} \rangle} \frac{(dN/dp_{\rm T})_{AA}}{(dN/dp_{\rm T})_{pp}}, \qquad (1a)$$

$$R_{\rm CP} = \frac{\left\langle N_{\rm bin}^{\rm peripheral} \right\rangle}{\left\langle N_{\rm bin}^{\rm central} \right\rangle} \frac{(dN/dp_{\rm T})_{\rm central}^{AA}}{(dN/dp_{\rm T})_{\rm peripheral}^{AA}}, \qquad (1b)$$

with  $\langle N_{\rm bin} \rangle$ , number of binary nucleon–nucleon collisions in (central or peripheral) AA collisions.

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In [2] the ALICE Collaboration reported a strong suppression of inclusive charged particles in central Pb–Pb collisions relative to pp at  $\sqrt{s_{NN}} = 2.76$  TeV in the region  $p_{\rm T} = 6-8$  GeV/c followed by a rise up to  $p_{\rm T} = 50$  GeV/c. The modification of the spectra of identified particles and especially the comparison of baryons and mesons may allow to disentangle differences between the energy loss of quarks and gluons. A difference in the modification pattern of baryons and mesons due to enhanced baryon production at intermediate  $p_{\rm T}$  ( $p_{\rm T} = 2-4$  GeV/c) is commonly referred to as the baryon-to-meson anomaly. In this contribution, the  $\Lambda$  baryon and the  $K_{\rm s}^0$  meson are studied in order to give a more detailed view of the parton energy loss at LHC energies.

### 2. Analysis

The  $\Lambda$  and  $K^0_{\rm s}$  are reconstructed employing a topological secondary vertex finder using tracking information from the TPC (Time Projection Chamber) and the ITS (Inner Tracking System). For each  $p_{\rm T}$  interval the yields are obtained by integrating over the invariant mass peak after the combinatorial background subtraction. The latter is based on a fit of the background with a polynomial of first and second order excluding the mass peak region. Systematic uncertainties related to the reconstruction efficiency cancel partially in the ratios  $R_{AA}$  and  $R_{CP}$ . The systematic errors given below are based on a conservative estimate of the uncertainty on the centrality dependence of the reconstruction efficiency. The  $p_{\rm T}$  spectra of A are corrected for feed-down using a preliminary centrality and  $p_{\rm T}$  dependent estimate of the contribution from  $\Xi$ . The uncertainty related to the centrality dependence of this estimate contributes mainly at low  $p_{\rm T}$  and is included in the systematic errors of  $R_{AA}$  and  $R_{CP}$ . The following results are extracted from two centrality classes in Pb–Pb collisions (0–5%, 60–80%) at  $\sqrt{s_{NN}} = 2.76$  TeV and from the corresponding pp reference data at  $\sqrt{s} = 2.76$  TeV.

# 3. Results

The  $R_{\rm CP}$  for  $\Lambda$  and  $K_{\rm s}^0$  of central (0–5%) to peripheral (60–80%) Pb–Pb collisions at  $\sqrt{s_{NN}} = 2.76$  TeV are shown in the left panel of Fig. 1. The results of  $K_{\rm s}^0$  agree with the  $R_{\rm CP}$  of charged kaons and of inclusive charged particles [2] for the whole  $p_{\rm T}$  range. Two  $p_{\rm T}$  domains can be distinguished for the  $\Lambda$ - $R_{\rm CP}$ : at high  $p_{\rm T}$  ( $p_{\rm T} > 8$  GeV/c) a common modification of  $\Lambda$  and  $K_{\rm s}^0$  with charged particles is observed, while at lower the  $\Lambda$ - $R_{\rm CP}$ shows a clear enhancement over the  $K_{\rm s}^0$ - $R_{\rm CP}$ , which seems to support the conjecture of a baryon-to-meson anomaly, as observed at RHIC (discussed in detail in [3]). In the right panel of Fig. 1 the ALICE results for  $\Lambda$ - $R_{\rm CP}$ are compared to that obtained by the STAR Collaboration for  $\Lambda + \overline{\Lambda}$  at  $\sqrt{s_{NN}} = 200$  GeV [4]. Within the uncertainties of the measurements, the results for  $R_{\rm CP}$  are comparable. It should be mentioned though that the ALICE results indicate that the  $\Lambda$ - $R_{\rm CP}$  enhancement is extended towards higher  $p_{\rm T}$ .



Fig. 1. The nuclear modification factor  $R_{\rm CP}$  of  $\Lambda$  and  $K_{\rm s}^0$  between central (0–5%) and peripheral (60–80%) Pb–Pb collisions at  $\sqrt{s_{NN}} = 2.76$  TeV. The boxes around the data points indicate the systematic error for  $K_{\rm s}^0$  and  $\Lambda$ . The uncertainty due to the calculation of  $\langle N_{\rm bin} \rangle$  is given by the gray boxes at the dotted line. Left panel: For the charged kaons (small filled circles) and the charged particles (filled triangles) only statistical errors are shown. Right panel: the comparison to the measurement for  $\Lambda + \overline{\Lambda}$  in Au–Au collisions at  $\sqrt{s_{NN}} = 200$  GeV measured by the STAR Collaboration [4]. The light gray box at unity indicates the error on  $\langle N_{\rm bin} \rangle$  for the STAR measurements.

The modification of the  $\Lambda$  and  $K_{\rm s}^0 p_{\rm T}$  spectra relative to pp at the same energy is shown in Fig. 2. In order to evaluate the modification, the results are presented together with the  $R_{AA}$  of charged particles [2]. The left panel shows  $R_{AA}$  for peripheral and the right panel for central collisions, respectively. As already seen in the measurement of  $R_{\rm CP}$ , the suppression of  $K_{\rm s}^0$  is comparable to that of charged particles for the whole  $p_{\rm T}$  range, in peripheral as well as in central collisions. In contrast, the  $\Lambda$  are differently suppressed at lower and higher  $p_{\rm T}$ : in particular only little nuclear modification,  $R_{AA} \approx 1$ , is observed for  $\Lambda$  at  $p_{\rm T} = 2-5$  GeV/c which results in an enhancement of the  $\Lambda$ - $R_{AA}$  over the  $K_{\rm s}^0$ - $R_{AA}$  as it is indicated by the  $R_{\rm CP}$ results. However, at large  $p_{\rm T}$ , modification is as strong as for the  $K_{\rm s}^0$ - $R_{AA}$ and hence as for that of charged particles.

Considering the  $R_{AA}$  of  $K_s^0$  in peripheral events (left panel of Fig. 2), it indicates a nearly constant and a rather moderate but significant suppression of  $R_{AA} \approx 0.6$ . Due to this moderate  $p_{\rm T}$  dependence of both, the  $\Lambda$  and  $K_s^0$  nuclear modification factors for peripheral collisions,  $R_{AA}$  is ex-

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pected to behave similarly to  $R_{\rm CP}$  for central collisions, only being scaled by some factor. This can be seen in the right panel in Fig. 2, where the  $\Lambda$ and  $K_{\rm s}^0 - R_{AA}$  for central events is depicted. A strong suppression at high  $p_{\rm T}$  with respect to pp collisions is observed for both hadrons which again is similar to the suppression of charged particles. The observation of a compatible suppression of strange and non-strange hadrons can be interpreted as a flavour independent modification at high  $p_{\rm T}$ . Corresponding results were obtained for D mesons supporting this interpretation [5].



Fig. 2. The nuclear modification factor  $R_{AA}$  of  $K_s^0$  and  $\Lambda$  for peripheral (60–80%) Pb–Pb collisions (left panel) and for central (0–5%) Pb–Pb collisions at  $\sqrt{s_{NN}} = 2.76$  TeV (right panel). Definition of errors as in Fig. 1. In both panels the charged hadron  $R_{AA}$  is shown in addition. Vertical error bars indicate the systematic uncertainties and the horizontal ticks show the statistical errors.

Figure 3 shows the comparison of  $\Lambda$ - $R_{AA}$  to the measurements by the STAR Collaboration at  $\sqrt{s_{NN}} = 200$  GeV for the mentioned centralities [6]. At LHC energies we observe a much smaller modification than at RHIC energies. Taking into account that the STAR  $\Lambda$ - $R_{CP}$  is compatible to our measurement the significant difference between STAR and ALICE  $\Lambda$ - $R_{AA}$  at  $p_{T} \approx 3$  GeV/c may be driven by the pp references rather than by nuclear effects. As the STAR results are limited to the intermediate  $p_{T}$  region, no statement on high  $p_{T}$  suppression of  $\Lambda$  and  $K_{s}^{0}$  with respect to pp collisions from RHIC to LHC energies is possible yet.

Finally, in Fig. 4 the ALICE results for  $R_{AA}$  of  $\Lambda$  and  $K_s^0$  and charged pions in central collisions are compared to calculations from the HIJING/BB v2.0 model [7]. In the case of mesons, the results are well described by the model up to  $p_T = 10 \text{ GeV}/c$ . However, for  $\Lambda - R_{AA}$  the calculations exceed the data at high  $p_T$  and the calculated maximum is at lower  $p_T$ . Hence, the common flavour and strangeness content independent high  $p_T$  suppression of  $\Lambda$  and  $K_s^0$ , as seen in the data, is not reproduced by the model.



Fig. 3. The nuclear modification factor  $R_{AA}$  of  $\Lambda$  in comparison to the measurement for Au–Au collisions at  $\sqrt{s_{NN}} = 200$  GeV by the STAR Collaboration [6]. The left panel shows the result for peripheral (60–80%) and the right panel  $R_{AA}$  for central (0–5%) collisions. Definition of errors as in Fig. 1.



Fig. 4. The nuclear modification factor  $R_{AA}$  of  $\Lambda$ ,  $K_s^0$  and charged pions for central (0–5%) Pb–Pb collisions at  $\sqrt{s_{NN}} = 2.76$  TeV in comparison to HIJING/BB v2.0 model calculations [7]. Definition of errors as in Fig. 1.

#### 4. Summary

We have presented the measurements of the nuclear modification factors of  $\Lambda$  and  $K_s^0$  up to  $p_T = 16 \text{ GeV}/c$  in Pb–Pb collisions at  $\sqrt{s_{NN}} = 2.76 \text{ TeV}$ . For both particle species a strong suppression at high  $p_T$  ( $p_T > 8 \text{ GeV}/c$ ) in central collisions with respect to pp is observed. A similar suppression at high  $p_{\rm T}$  of both hadrons is found in the ratio of central-to-peripheral collisions, too. The nuclear modification of  $\Lambda$  and  $K_{\rm s}^0$  is compatible with the modification of charged particles at high  $p_{\rm T}$ . At lower  $p_{\rm T}$  ( $p_{\rm T} < 5 \,{\rm GeV}/c$ ) we observe an enhancement of the  $\Lambda$ - $R_{AA}$  with respect to the  $K_{\rm s}^0$ - $R_{AA}$ , which can be related to the baryon-to-meson anomaly. While the  $R_{\rm CP}$  is similar at RHIC and LHC energies, we find significantly less  $\Lambda$ - $R_{AA}$  enhancement relative to  $K_{\rm s}^0$ - $R_{AA}$  at intermediate  $p_{\rm T}$  in central and peripheral events, as compared to the STAR results. Regarding the comparison to model calculations with the HIJING/BB v2.0 model, we observe an agreement with our results for  $K_{\rm s}^0$  and charged pion  $R_{AA}$ , whereas the  $R_{AA}$  of  $\Lambda$  is not described by the model.

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