

J/ψ POLARIZATION IN $p + p$ COLLISIONS AT $\sqrt{s} = 200$ GeV IN STAR*

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In this paper, J/ψ polarization at mid-rapidity in $p + p$ collisions at $\sqrt{s} = 200$ GeV measured in the STAR experiment at RHIC is reported. J/ψ production is analyzed via the dielectron decay channel. J/ψ polarization is extracted from the decay angular distribution in the helicity frame. The J/ψ polarization is measured at transverse momentum range (2–6) GeV/ c and is found to be consistent with NLO⁺ Color Singlet Model (NLO⁺ CSM), Color Octet Model (COM) predictions and with no polarization within current uncertainties.

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1. Introduction

A number of models with different J/ψ production mechanisms are able to describe the measured J/ψ production cross section reasonably well. It suggests that other observables are needed to discriminate between different J/ψ production models. J/ψ spin alignment, commonly named as J/ψ polarization, can be used as such an observable since various models have different p_T dependent predictions on the J/ψ polarization.

The Color Evaporation Model (CEM) describes J/ψ cross sections measured in many experiments reasonably well but has no prediction power regarding the J/ψ polarization. The NLO⁺ Color Singlet Model (CSM) [1] predicts longitudinal J/ψ polarization in the helicity frame at low and mid p_T at mid-rapidity. These recent calculations of the CSM for the yield differential in p_T show better agreement with RHIC data at low and mid p_T than earlier CSM calculations and the result for J/ψ polarization is in good agreement with the PHENIX data. Non-relativistic quantum chromodynamics (NRQCD) effective theory, also known as Color Octet Model (COM)

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predicts transverse polarization at high J/ψ p_T , $p_T > 5$ GeV/ c which is in disagreement with the polarization measurement from the CDF experiment at FermiLab [2]. On the other hand, a tuned NRQCD predicts longitudinal J/ψ polarization at $1.5 < p_T < 5$ GeV/ c at mid-rapidity and is able to qualitatively describe the PHENIX J/ψ polarization measurements as well as the cross section measurements [3].

The measurement of J/ψ polarization at $p_T > 5$ GeV/ c is expected to have discrimination power against different models of J/ψ production since, e.g., NLO⁺ CSM and COM predict different polarization in that J/ψ p_T region.

1.1. Decay angular distribution

In this study, J/ψ polarization is analyzed via the angular distribution of the electron decay from charmonium in the helicity frame [4]. The angular distribution is derived from the density matrix elements of the production amplitude using parity conservation rules. Polar angle θ is the angle between the positron momentum vector in the J/ψ rest frame and J/ψ momentum vector in the laboratory frame.

The angular distribution integrated over the azimuthal angle is parametrized

$$\frac{dN}{d\cos\theta} \propto 1 + \lambda \cos^2\theta, \quad (1)$$

where λ is the polarization parameter that contains both the longitudinal and transverse components of the J/ψ cross section. When $\lambda = 0$ there is no polarization, $\lambda = -1$ means full longitudinal polarization and $\lambda = 1$ corresponds to full transverse polarization.

2. Data analysis

In this analysis, data recorded in 200 GeV $p + p$ collisions in the STAR experiment in year 2009 is used. The analyzed data was sampled from an integrated luminosity of ~ 1.5 pb⁻¹ and was triggered by the STAR Barrel Electromagnetic Calorimeter (BEMC). The trigger required transverse energy deposited in a single BEMC tower ($\Delta\eta \times \Delta\phi = 0.05 \times 0.05$) to be within $2.6 < E_T \leq 4.3$ GeV.

J/ψ is reconstructed via its dielectron decay channel $J/\psi \rightarrow e^+e^-$ with branching ratio BR = 5.9%, and is required that at least one of electrons from J/ψ decay satisfies the trigger conditions. The Time Projection Chamber (TPC), Time-of-Flight (TOF) and BEMC detectors are used to reconstruct and identify electrons. The TPC provides information about dE/dx . Information from the TOF is very useful for electron identification and hadron rejection at lower momenta, where electron and hadron dE/dx bands overlap, $|1/\beta - 1| < 0.03$ ($\beta = v/c$) cut was applied at $p < 1.4$ GeV/ c .

In 2009, 72% of the TOF detector was installed. At higher momenta the BEMC can reject hadrons very efficiently. For momenta above 1.4 GeV/ c a cut of $E/p > 0.5$ was used, where E is energy deposited in a single BEMC tower (for electrons E/p ratio is expected to be ≈ 1).

2.1. J/ψ signal

The cuts mentioned above allowed us to obtain a very clear J/ψ signal with a signal to background ratio of 15 and very high significance of 26σ . Fig. 1 (a) shows the invariant mass distribution of all combinations of e^+e^- pairs in black (closed circles) and combinatorial background in grey/red (open circles). The combinatorial background is calculated using like-sign technique, from a sum of e^+e^+ and e^-e^- pairs. The J/ψ signal is obtained by subtracting the (red) histogram (background) from the black histogram (signal + background), see Fig. 1 (b). The number of J/ψ s calculated by counting bin entries in the invariant mass range (2.9–3.3) GeV/ c^2 is 772 ± 29 in the J/ψ p_T range (2–6) GeV/ c and $|y| < 1$. The same mass window is used for the polarization analysis for which the signal is split into 3 statistically comparable p_T bins.

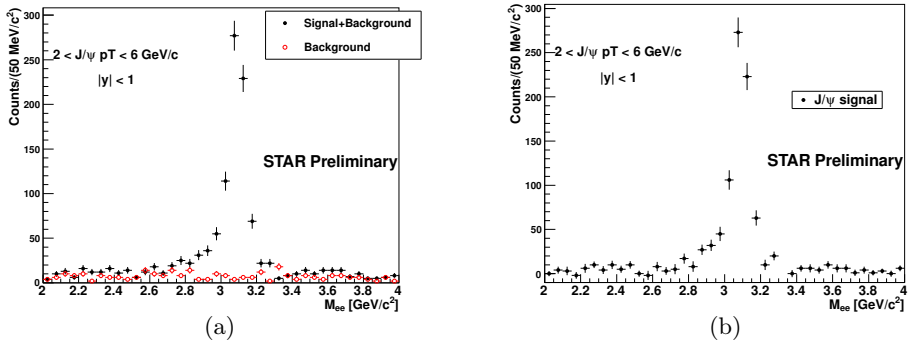


Fig. 1. Invariant mass distributions in J/ψ mass window (2.9–3.3) GeV/ c^2 , J/ψ p_T range (2–6) GeV/ c and $|y| < 1$. (a) shows signal + background in black (full circles) and like-sign background in grey/red (open circles). (b) shows J/ψ signal after the combinatorial background subtraction.

3. J/ψ polarization

The uncorrected $\cos\theta$ distribution is obtained using the same electron identification cuts as J/ψ signal, in the J/ψ p_T range of (2–6) GeV/ c and $|y| < 1$. The $\cos\theta$ distribution is divided into 3 J/ψ p_T bins: $2 < p_T < 3$ GeV/ c , $3 < p_T < 4$ GeV/ c and $4 < p_T < 6$ GeV/ c . The $\cos\theta$ distributions for the 3 p_T bins, with combinatorial background subtracted are presented in Fig. 2 (a), 2 (c), 2 (e) respectively.

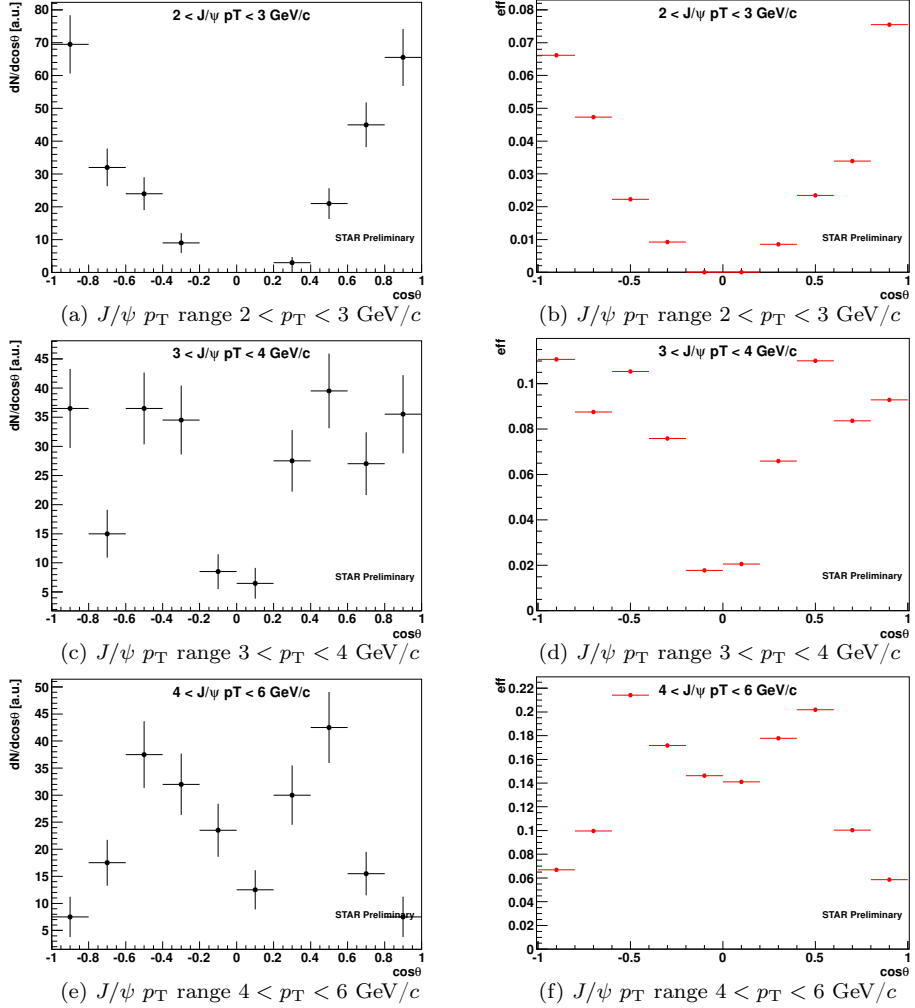


Fig. 2. Left plots — uncorrected $\cos\theta$ distribution with combinatorial background subtracted. Right plots — $\cos\theta$ efficiency.

3.1. Corrections

In order to get the $\cos\theta$ corrections, unpolarized Monte Carlo J/ψ s with uniform p_T and rapidity distributions are embedded into real events and the detector response is simulated. All cuts used in the analysis were applied and the $\cos\theta$ efficiency as a function of J/ψ transverse momentum was obtained. $\cos\theta$ distributions were also re-weighted according to the real J/ψ p_T and y shapes. Obtained in that way, corrections were applied to uncorrected $\cos\theta$ distributions from data in 1 GeV/c wide J/ψ p_T bins. Corrections include acceptance correction, tracking efficiency, electron identification efficiency

and trigger efficiency. The total J/ψ efficiency is shown *versus* J/ψ transverse momentum in Fig. 2 (b), 2 (d), 2 (f). The most critical factor is trigger efficiency. At least one of electrons from J/ψ decay must have p_T above 2.6 GeV/ c since it is required to satisfy the trigger conditions. It causes significant loss in number of observed J/ψ at lower J/ψ p_T and the efficiency decrease with decreasing $|\cos \theta|$. It is well visible in Fig. 2 (b), where we lose all entries at $\cos \theta \sim 0$. With increasing J/ψ p_T the trigger efficiency increase but because the trigger has also the upper threshold ($E_T \leq 4.3$ GeV) and due to the acceptance effect (single electron $p_T > 0.4$ GeV/ c and $|\eta| < 1$) we see drop of total efficiency at $|\cos \theta| \sim 1$, see Fig. 2 (f).

3.2. J/ψ polarization result

The corrected $\cos \theta$ distributions are fitted with: $\text{norm}(1 + \lambda \cos^2 \theta)$ (see Eq. 1) with no constraints on the fit parameters, see Fig. 3 (a), 3 (b), 3 (c), where ‘norm’ is a normalization factor, λ is the polarization parameter. Lines represent the most likely fits. Obtained results of the polarization parameter are: $\lambda = 0.46 \pm 0.40(\text{stat.}) \pm 0.34(\text{sys.})$, $\lambda = -0.32 \pm 0.20(\text{stat.}) \pm 0.16(\text{sys.})$, $\lambda = -0.17 \pm 0.29(\text{stat.}) \pm 0.11(\text{sys.})$ for J/ψ p_T ranges:

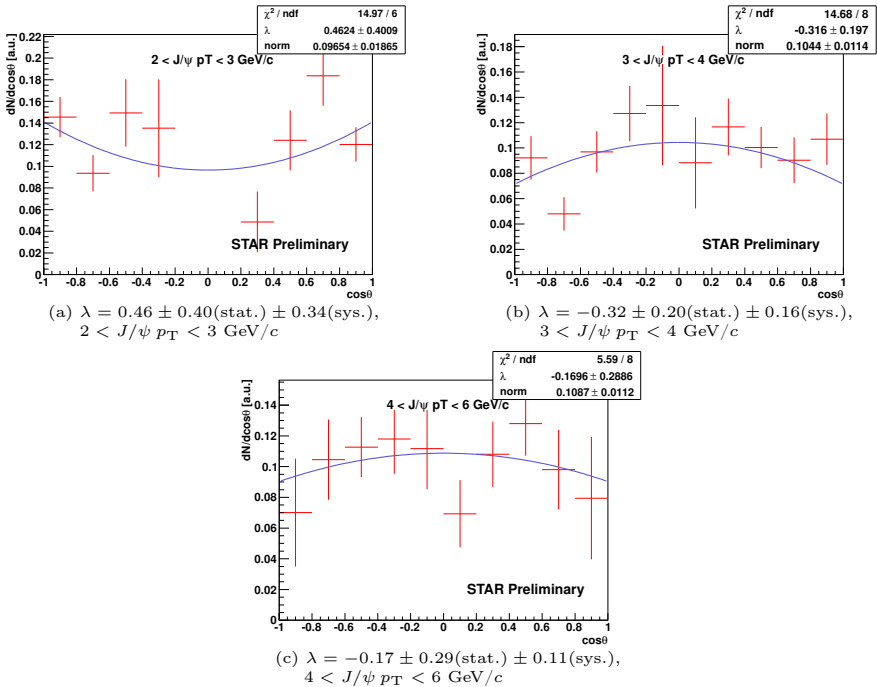


Fig. 3. Corrected $\cos \theta$ distributions with the fit: $\text{norm}(1 + \lambda \cos^2 \theta)$, errors are statistical.

(2–3) GeV/ c , (3–4) GeV/ c , (4–6) GeV/ c , respectively. Dominant sources of systematic uncertainties are: $\cos\theta$ binning and acceptance, J/ψ mass range, electron identification cut, trigger efficiency. Detailed work on systematic error estimation is in progress.

Polarization parameters as a function of J/ψ p_T are shown in Fig. 4. The STAR result (red star symbols) is compared with NLO⁺ CSM [1] (light grey/blue shaded area) and COM [3] (dark gray hatched area) model predictions and with the PHENIX result for J/ψ polarization at mid-rapidity (black symbols) [5]. The STAR result is consistent with the COM and CSM predictions within current experimental and theoretical uncertainties. The measurement is also consistent with the PHENIX data and extends the p_T reach to ~ 6 GeV/ c .

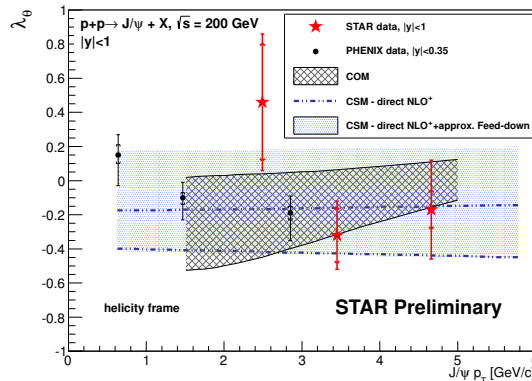


Fig. 4. STAR polarization parameter λ versus J/ψ p_T (red star symbols) compared with PHENIX result (black symbols) [5], NLO⁺ CSM (light grey/blue shaded area) [1] and COM [3] (dark gray hatched area) predictions.

4. Summary

In this report, the J/ψ polarization measurement from the STAR experiment at mid-rapidity is presented. The polarization parameter λ is extracted in the helicity frame in 3 J/ψ p_T bins. Within current uncertainties the obtained transverse momentum dependent λ parameter is consistent with NLO⁺ CSM and COM model predictions, and with no polarization.

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