RECENT DEVELOPMENTS OF A MULTI-PHASE TRANSPORT MODEL*

Z.W. Lin

Department of Physics, East Carolina University, Greenville NC 27858-4353, USA

(Received February 14, 2014)

After the public release of A Multi-Phase Transport (AMPT) model in 2004 and detailed descriptions of its physics in a 2005 paper, the model has been constantly updated and developed to make it more versatile and to include more physical processes. This is an overview of recent developments of the AMPT model. Ongoing work to fix the violation of charge conservation in the code as well as possible directions for future work are also discussed.

DOI:10.5506/APhysPolBSupp.7.191 PACS numbers: 12.38.Mh, 24.10.Lx, 24.10.Jv, 25.75.Nq

1. Introduction

A Multi-Phase Transport model was constructed specifically for the study of relativistic heavy ion collisions. It intends to serve as a self-contained phenomenological model, since it incorporates essential stages of heavy ion collisions from the initial condition to final observables on an event-by-event basis, including the parton cascade, hadronization and the hadron cascade. The default version of the AMPT model was first constructed to predict particle yields and momentum spectra in heavy ion collisions at RHIC [1]. After large elliptic flow was discovered at RHIC, the string melting version of the AMPT model was introduced [2] in order to take into account the effects on flow from the whole partonic system in the overlap volume (instead of from only minijet partons in the default version).

To help illustrate the difference between the default version and the string melting version of the AMPT model, Fig. 1 shows the snapshot at the time of 8 fm/c of a central Au+Au event from the default version (upper panel) and from the string melting version (lower panel), where both calculations use an

^{*} Lecture presented at the XXXI Max Born Symposium and HIC for FAIR Workshop "Three Days of Critical Behaviour in Hot and Dense QCD", Wrocław, Poland, June 14–16, 2013.

identical HIJING initial condition such as the position of every participant nucleon and spectator nucleon. This is a side view where the two beams come from the left- and right-hand sides of the showed box respectively, and the full length of the box is 30 fm along each direction. The insert on the left-hand side of each panel shows the number of several particle species in the event as a function of time (up to the time of 8 fm/c for this figure). Black particles here represent pions which number in the left insert has been scaled down by a factor of 5, dark grey/red particles represent gluons (in the upper panel) or quarks (in the lower panel), while light grey/cvan particles in the lower panel represent antiquarks. We see in both panels that many partons still exist on the left- and right-hand sides (at large rapidities) while more hadrons have been formed in the middle (around mid-rapidity): this is a consequence of the time dilation that delays the dynamics at large rapidities when the global time is used to calculate particle interactions, as is the case for AMPT. In the lower string melting panel, we see many more partons and a delayed hadron phase compared to the upper panel for the default version, and this demonstrates the dominance of partons in early times in the string melting version. The full time evolutions of these events are available as animation files at the author's website [3].

The source code of the AMPT model was first publicly released online around April 2004, and a subsequent publication [4] provided detailed descriptions of the model such as the included physics processes and modeling assumptions. The model has since been applied to the study of many observables in heavy ion collisions. The AMPT model is also a useful test bed of different ideas; for example, the connection between triangular flow and initial geometrical fluctuations was discovered with AMPT [5]. Shown in Fig. 2 is the view along the beam axis for the transverse positions of particles in a peripheral Au+Au event at the early time of 1 fm/c from the string melting version of AMPT. Here, white particles that are mostly in the peripheral region represent spectator nucleons, and colored particles are mostly partons (quarks and antiquarks) that have already formed. We can clearly see that the shape of the parton system in the overlap region is not elliptical in the transverse plane, and for this particular event it looks quite like a triangle. As found in Ref. [5], subsequent partonic and hadronic scatterings convert the irregular geometry in each event into particle correlations in momentum space such as triangular flow.

2. Recent developments of the AMPT model

The "official" 2004 version of AMPT v1.11(default)/v2.11(string melting) and the 2008 version v1.21(default)/v2.21(string melting) are available at the OSCAR website [6]. Since then the AMPT model has been constantly updated and developed, sometimes at users' requests, to make it more versatile and stable and to include more physical processes. Several of these more recent "test" versions of the source codes are available at the author's website [3], where the AMPT Users' Guide and the readme file in the source code detail the changes made in each version.

The following summarizes some of the developments and updates of the model after the 2008 version v1.21/v2.21. Deuteron (d) interactions have been included in the hadron cascade, including inelastic collisions $d + M \leftrightarrow B + B$ (where M represents a meson and B represents a baryon), elastic collisions of deuterons, and corresponding interactions for anti-deuterons [7]. We have also extended the functionality of the code to allow users to (1) insert arbitrary user-defined hadrons at the start of hadron cascade, (2) obtain the transverse positions of all initial nucleons (participant and spectator nucleons) as well as their present and original flavor codes, (3) trigger events so that each event will have at least one initial minijet parton with transverse momentum above a user-defined value, (4) embed one back-to-back q/\bar{q} pair per event at the user-specified transverse momentum and transverse position at the start of parton cascade, (5) obtain the complete initial parton information before parton cascade as well as the complete parton collision

history, (6) set the nuclear shadowing strength anywhere between no shadowing and the default HIJING shadowing. In addition, we added deformed uranium-238 as a type of projectile and/or target nucleus [8]; since this extension was done in parallel with some of the above developments, it has not yet been incorporated into the posted AMPT versions [3].

3. Ongoing work and future directions

An ongoing work is to implement electric charge conservation in the AMPT model so that it can be used to reliably study charge-dependent variables such as charge fluctuations in heavy ion collisions. Currently, the AMPT model violates charge conservation due to two reasons. First, the hadron cascade has K^+/\bar{K}^- as explicit particles but not K^0/\bar{K}^0 , so the code changes K^0 to K^+ and changes \overline{K}^0 to K^- before hadron cascade in order to include neutral kaons in the hadron cascade, and after hadron cascade the code changes half of the final K^+ into K^0 and changes half of K^- to \bar{K}^0 . This however introduces a violation of charge conservation. Secondly, not all hadron reactions or resonance decays in the hadron cascade of AMPT are implemented for each possible individual isospin configuration; for some channels the isospin-averaged cross section is used instead and the electric charge of each final-state particle is selected randomly from all possible values (without considering the initial-state charge configuration). For example, each final-state pion could have +, 0, or - charge in the reaction $\pi\eta \to \pi\pi$ in AMPT, where, for example, $\pi^+\eta \to \pi^+\pi^0$ is allowed but $\pi^+\eta \to$ $\pi^+\pi^-$ should have been forbidden. To fix the charge violation problem in the AMPT model, we thus need to add K^0 and \overline{K}^0 as explicit particles in the hadron cascade; we also need to consider each isospin configuration for the hadronic channels and use isospin-dependent cross sections.

Regarding future directions of the AMPT model, from the author's point of view we may take two different approaches to further develop such a selfcontained phenomenological model for relativisitc heavy ion collisions. One approach is to develop key ingredients of the AMPT model within the same scheme of a partonic phase followed by a hadronic phase, while the other approach is to couple the AMPT model with a hydrodynamic model.

The first approach will include the following developments: update of parton distribution functions in nuclei including the nuclear shadowing effect, development of the dynamical parton recombination process (*e.g.* by using local parton density as the recombination criterion), incorporation of inelastic parton reactions, and the consideration of gluons in parton recombination. For example, a study of the AMPT model [9] has shown that the quark coalescence in the string melting version of AMPT starts too late, *i.e.*, it starts at energy densities that are too low ($\ll 1 \text{ GeV/fm}^3$) compared to

the equation of state from lattice QCD. Implementing parton recombination according to local energy density would make the effective equation of state in the AMPT model more realistic and also lead to a better success rate of parton recombination due to the higher parton density at hadronization.

The second approach aims to couple AMPT with a hydrodynamic model, and the resulted hybrid model will provide us a more direct link to QCD variables and properties. The AMPT model will provide the initial condition, including the event-by-event fluctuations and correlations, to the hydrodynamic model. After the hydrodynamic model calculates the evolution of the dense matter through the QCD phase transition, hadron interactions will be taken into account by the hadron cascade of the AMPT model. This hybrid model will be a self-contained model in that it will have self-consistent initial conditions for different collision systems at different energies, and a 3+1D hydrodynamic code will be ideal so that the hybrid model can address non-equilibrium evolution and observables at different rapidities.

The author acknowledges support from the Helmholtz International Center for FAIR, where part of the ongoing work was done at the Institut für Theoretische Physik at Goethe University.

REFERENCES

- [1] B. Zhang, C.M. Ko, B.-A. Li, Z.-W. Lin, *Phys. Rev.* C61, 067901 (2000).
- [2] Z.-W. Lin, C.M. Ko, *Phys. Rev.* C65, 034904 (2002).
- [3] http://personal.ecu.edu/linz/ampt/
- [4] Z.-W. Lin et al., Phys. Rev. C72, 064901 (2005).
- [5] B. Alver, G. Roland, *Phys. Rev.* C81, 054905 (2010) [*Erratum ibid.*, C82, 039903 (2010)].
- [6] https://karman.physics.purdue.edu/OSCAR
- [7] Y. Oh, Z.-W. Lin, C.M. Ko, *Phys. Rev.* C80, 064902 (2009).
- [8] M.R. Haque, Z.-W. Lin, B. Mohanty, *Phys. Rev.* C85, 034905 (2012).
- [9] B. Zhang, L.-W. Chen, C.M. Ko, J. Phys. G 35, 065103 (2008).



Fig. 1. Side view of a 200 AGeV Au+Au event at impact parameter b = 0 fm at the time of 8 fm/c. Upper panel: from the default version of AMPT. Lower panel: from the string melting version of AMPT (with identical initial condition for the event).



Fig. 2. View of the transverse plane along the beam axis at the time of 1 fm/c for a 200 AGeV Au+Au event at impact parameter b = 10 fm from the string melting version of AMPT.