

THE UNQUENCHED QUARK MODEL*

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In this contribution, we briefly analyze the formalism of the unquenched quark model (UQM) and its application to the calculation of hadron spectra with self-energy corrections, due to the coupling to the meson–meson continuum. In the UQM, the effects of $q\bar{q}$ sea pairs are introduced explicitly into the quark model through a QCD-inspired 3P_0 pair-creation mechanism. The UQM formalism can be extended to include also the effects of hybrid mesons, *i.e.* hybrid loops. The main applications to spectroscopy and decays are analyzed.

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1. Introduction

The quark model can reproduce the behavior of observables such as the spectrum and the magnetic moments, but it neglects pair-creation (or continuum-coupling) effects. Above threshold, this coupling leads to strong decays; below threshold, it leads to virtual $q\bar{q}$ – $q\bar{q}$ (qqq – $q\bar{q}$) components in

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the hadron wave function and shifts of the physical mass with respect to the bare mass. The unquenching of the quark model for hadrons is a way to take these components into account.

Pioneering work on the unquenching of meson quark models was done by van Beveren and Rupp who used a t-matrix approach [1, 2], while Törnqvist and collaborators [3, 4] used their unitarized QM. These methods were used (with a few variations) by several authors to study the influence of the meson–meson (meson–baryon) continuum on meson (baryon) observables. As an example, we mention the study of the scalar meson nonet (a_0 , f_0 , *etc.*) of Ref. [2, 5] in which the loop contributions are given by the hadronic intermediate states that each meson can access. It is via these hadronic loops that the bare states become “dressed” and the hadronic loop contributions totally dominate the dynamics of the process. A very similar approach was developed by Boglione and Pennington in Ref. [6], in which they investigated the dynamical generation of the scalar mesons by initially inserting only one “bare seed”. The study of Ref. [7] demonstrates that the effect of the $q\bar{q}$ sea pairs in meson spectroscopy is simply a renormalization of the meson string tension. The strangeness content of the nucleon electromagnetic form factors was investigated in [8, 9]. Also Capstick and Morel in Ref. [10] analyzed baryon meson loop effects on the spectrum of nonstrange baryons. Eichten *et al.* explored the influence of the open-charm channels on the charmonium properties, using the Cornell coupled-channel model [11] to assess departures from the single-channel potential-model expectations. The flavor asymmetry of the proton was studied in the framework of the unquenched quark model (UQM) [9, 12–16], in which the effects of the quark–antiquark pairs were introduced into the constituent quark model (CQM) in a systematic way and the wave functions were given explicitly. The approach is a generalization of the unitarized quark model [1, 2, 4, 5].

In this contribution, we discuss some of the latest applications of the UQM to the study of meson observables.

2. UQM

2.1. Formalism

In the unquenched quark model for baryons [9, 12] and mesons [13–16], the hadron wave function is made up of a zeroth order qqq ($q\bar{q}$) configuration plus a sum over the possible higher Fock components, due to the creation of 3P_0 $q\bar{q}$ pairs. Thus, we have

$$|\psi_A\rangle = \mathcal{N} \left[|A\rangle + \sum_{BC\ell J} \int d\vec{K} k^2 dk \left| BC\ell J; \vec{K}k \right\rangle \frac{\langle BC\ell J; \vec{K}k | T^\dagger | A \rangle}{E_a - E_b - E_c} \right], \quad (1)$$

where $T^\dagger$ stands for the 3P_0 quark–antiquark pair-creation operator [13–16], A is the baryon/meson, B and C represent the intermediate state hadrons, E_a , E_b and E_c are the corresponding energies, k and ℓ the relative radial momentum and orbital angular momentum between B and C , and $\vec{J} = \vec{J}_b + \vec{J}_c + \vec{\ell}$ is the total angular momentum. It is worthwhile noting that in Refs. [13–17], the constant pair-creation strength in the operator (1) was substituted with an effective one, to suppress unphysical heavy quark pair-creation.

In the UQM [9, 12–16], the matrix elements of an observable $\hat{O}$ can be calculated as

$$O = \langle \psi_A | \hat{O} | \psi_A \rangle, \quad (2)$$

where $|\psi_A\rangle$ is the state of Eq. (1). The result will receive a contribution from the valence part and one from the continuum component, which is absent in naive QM calculations.

The introduction of continuum effects in the QM can thus be essential to study observables that only depend on $q\bar{q}$ sea pairs, like the strangeness content of the nucleon electromagnetic form factors [8, 9]. In other cases, continuum effects can provide important corrections to baryon/meson observables, like the self-energy corrections to meson masses [13–16] or the importance of the orbital angular momentum in the spin of the proton [12].

2.2. $c\bar{c}$ and $b\bar{b}$ spectra with self-energy corrections in the UQM

In Refs. [13–16], the method was used by some of us to compute the $c\bar{c}$ and $b\bar{b}$ spectra with self-energy corrections, due to continuum coupling effects. In the UQM, the physical mass of a meson,

$$M_a = E_a + \Sigma(E_a), \quad (3)$$

is given by the sum of two terms: a bare energy, E_a , calculated within a potential model [18], and a self energy correction,

$$\Sigma(E_a) = \sum_{BC\ell J} \int_0^\infty k^2 dk \frac{|M_{A \rightarrow BC}(k)|^2}{E_a - E_b - E_c}, \quad (4)$$

computed within the UQM formalism.

Our results for the self energies of charmonia [14, 16] and bottomonia [13, 15, 16] show that the pair-creation effects on the spectrum of heavy mesons are quite small. Specifically for charmonium and bottomonium states, they are of the order of 2–6% and 1%, respectively. The relative mass shifts, *i.e.* the difference between the self energies of two meson states, are of the order of a few tens of MeV. However, as QMs can predict the

meson masses with relatively high precision in the heavy quark sector, even these corrections can become significant. These results are particularly interesting in the case of states close to an open-flavor decay threshold, like the $X(3872)$ and $\chi_b(3P)$ mesons. For example, in our picture, the $X(3872)$ can be interpreted as a $c\bar{c}$ core [the $\chi_{c1}(2^3P_1)$], plus higher Fock components due to the coupling to the meson–meson continuum. In Ref. [16], we showed that the probability to find the $X(3872)$ in its core or continuum components is approximately 45% and 55%, respectively.

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