

GEOMETRIC OBJECTS RELATED TO THE POTENTIAL OF ELECTRIC CHARGES

J. MOZRZYMAS

Institute of Theoretical Physics, University of Wrocław
Plac Maksa Borna 9, 50-204 Wrocław, Poland

Dedicated to the memory of Professor Jan Rzewuski

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We derive explicit formulas for curvature and torsion of a line of the field of n electric charges. These formulas show that in general the torsion of a field line is not zero if $n \geq 3$. We also propose a geometric interpretation of the derived formulas. In the second part of the paper we present an outline of a new description of equipotential surfaces of two and three electric charges. In this description the golden section appears in a natural way when two electric charges are equal. This approach also relates an equipotential surface of three charges to the classic cubic surface containing twenty seven straight lines.

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1. Introductory remarks

This paper consists of two parts. In the first part (Section 2) we follow [1] in deriving explicit formulas for curvature and torsion of a line of the field of n electric charges. In particular, these formulas imply that for two charges the torsion is always zero, and in general, is not zero when $n \geq 3$. Section 3 gives a geometric interpretation of the formulas derived in Section 2.

The second part (Sections 4 and 5) contains an outline of a new description of equipotential surfaces of two and three electric charges. This description is based on the notion of an auxiliary line L_2 and an auxiliary surface L_3 . The auxiliary objects L_2 and L_3 give a geometric insight into a natural parametrization of equipotential surfaces of two and three electric charges. In Section 4 we derive formulas which show that the famous golden section, [2], appears when two charges are equal. In Section 5, making use of the finite Fourier transform $F(3)$, [3], we relate the auxiliary surface L_3 to the classic cubic surface containing twenty seven straight lines, [4-6].

2. Curvature and torsion of a field line of n electric charges

We use the following notation:

q_i ($i = 1, 2, \dots, n$) — i -th electric charge,

$\vec{r}(x, y, z)$ — position vector of the point (x, y, z) ,

$\vec{r}_i = (x_i, y_i, z_i)$ — position vector of q_i ,

$R_i = |\vec{R}_i| = |\vec{r} - \vec{r}_i|$ — distance between (x, y, z) and q_i .

Then the potential Φ and the electric field \vec{E} of q_1, q_2, \dots, q_n are given by the classic formulas

$$\Phi = \sum_{i=1}^n q_i R_i^{-1}, \quad (2.1)$$

$$\vec{E} = -\text{grad } \Phi. \quad (2.2)$$

Let

$$\vec{r}(t) = (x(t), y(t), z(t)) \quad (2.3)$$

be a line parametrized by t . It is well known that the line $\vec{r}(t)$ is by definition a line of the field \vec{E} if $\vec{E}(\vec{r}(t))$ is tangent to the line $\vec{r}(t)$ for any value of the parameter t :

$$\vec{r}'(t) = \vec{E}(\vec{r}(t)), \quad (2.4)$$

where

$$\vec{r}'(t) = \frac{d\vec{r}}{dt}. \quad (2.5)$$

The curvature (first curvature) k_1 and torsion (second curvature) k_2 of a line $\vec{r}(t)$ are given by the well known formulas (see, for example, [7, 8]):

$$k_1 = \frac{|\vec{r}' \times \vec{r}''|}{|\vec{r}'|^3}, \quad (2.6)$$

$$k_2 = -\frac{[\vec{r}', \vec{r}'', \vec{r}''']}{|\vec{r}' \times \vec{r}''|^2}, \quad (2.7)$$

where $[\vec{r}', \vec{r}'', \vec{r}''']$ denotes the mixed product of \vec{r}' , \vec{r}'' and \vec{r}''' :

$$[\vec{r}', \vec{r}'', \vec{r}'''] = (\vec{r}' \times \vec{r}'') \cdot \vec{r}'''. \quad (2.8)$$

From the above formulas it follows that the curvature k_1 and torsion k_2 can be written in the form

$$k_1 = \frac{|\vec{E} \times \vec{F}|}{|\vec{E}|^3}, \quad (2.9)$$

$$k_2 = -\frac{[\vec{E}, \vec{F}, \vec{G}]}{|\vec{E} \times \vec{F}|^2}, \quad (2.10)$$

where

$$\vec{F} = \sum_{i=1}^n b_i \vec{R}_i, \tag{2.11}$$

$$\vec{G} = \sum_{i=1}^n c_i \vec{R}_i, \tag{2.12}$$

$$a_i = q_i R_i^{-3}, \tag{2.13}$$

$$b_i = a'_i = -3R_i^{-2} a_i A_i, \tag{2.14}$$

$$c_i = a''_i = -3R_i^{-2} a_i [(-5R_i^{-2} A_i + A)A_i + B_i + E^2], \tag{2.15}$$

$$A_i = \vec{R}_i \vec{E}, \tag{2.16}$$

$$B_i = \vec{R}_i \vec{F}, \tag{2.17}$$

$$A = \sum_{i=1}^n a_i. \tag{2.18}$$

It is easy to verify that the vector $\vec{E} \times \vec{F}$ and the mixed product $[\vec{E}, \vec{F}, \vec{G}]$ can be written in the form

$$\vec{E} \times \vec{F} = \sum_{\substack{i,1=1 \\ (i < j)}} \mu_{ij} \vec{S}_{ij}, \tag{2.19}$$

$$[\vec{E}, \vec{F}, \vec{G}] = \sum_{\substack{i,j,k=1 \\ (i < j < k)}} \sigma_{ijk} V_{ijk}, \tag{2.20}$$

where

$$\mu_{ij} = \begin{vmatrix} a_i & a_j \\ b_i & b_j \end{vmatrix}, \tag{2.21}$$

$$\sigma_{ijk} = \begin{vmatrix} a_i & a_j & a_k \\ b_i & b_j & b_k \\ c_i & c_j & c_k \end{vmatrix}, \tag{2.22}$$

$$\vec{S}_{ij} = \vec{R}_i \times \vec{R}_j, \tag{2.23}$$

$$V_{ijk} = [\vec{R}_i, \vec{R}_j, \vec{R}_k]. \tag{2.24}$$

3. A geometric interpretation of the formulas (2.9)–(2.10) as (2.19)–(2.20)

We introduce the following terminology:

- $a_i \vec{R}_i$ — the first scaling of \vec{R}_i ,
- $b_i \vec{R}_i$ — the second scaling of \vec{R}_i ,
- $c_i \vec{R}_i$ — the third scaling of \vec{R}_i ,

$$\left. \begin{aligned} \vec{E} &= \sum_{i=1}^n a_i \vec{R}_i \\ \vec{F} &= \sum_{i=1}^n b_i \vec{R}_i \\ \vec{G} &= \sum_{i=1}^n c_i \vec{R}_i \end{aligned} \right\} \begin{array}{l} \text{Resulting vectors of the first, second} \\ \text{and third scaling of } \vec{R}_1, \dots, \vec{R}_n. \end{array}$$

- $P(\vec{A}, \vec{B})$ — parallelogram spanned on the vectors \vec{A} and \vec{B} ,
- $P(\vec{A}, \vec{B}, \vec{C})$ — parallelepiped spanned on the vectors \vec{A} , \vec{B} and \vec{C} ,
- $\vec{A} \times \vec{B}$ — surface vector of $P(\vec{A}, \vec{B})$,
- $[\vec{A}, \vec{B}, \vec{C}]$ — volume of $P(\vec{A}, \vec{B}, \vec{C})$,
- $|\vec{S}_{ij}| = |\vec{R}_i \times \vec{R}_j|$ — elementary surface,
- $V_{ijk} = [\vec{R}_i, \vec{R}_j, \vec{R}_k]$ — elementary volume,
- $\mu_{ij} \vec{S}_{ij}$ — surface scaling of \vec{S}_{ij} ,
- $\sigma_{ijk} V_{ijk}$ — volume scaling of V_{ijk} ,
- $NES(n) = \binom{n}{2}$ — number of elementary surfaces,
- $NEV(n) = \binom{n}{3}$ — number of elementary volumes,

Making use of this terminology we obtain the following geometric interpretation of the formulas (2.9)–(2.10) and (2.19)–(2.20):

- (i) The curvature k_1 is equal to the ratio of the surface $|\vec{E} \times \vec{F}|$ to the third power of the length of \vec{E} .
- (ii) The torsion k_2 is equal to minus the ratio of the volume $[\vec{E}, \vec{F}, \vec{G}]$ to the square of the surface $|\vec{E} \times \vec{F}|$.
- (iii) The volume $[\vec{E}, \vec{F}, \vec{G}]$ consists of $\binom{n}{3}$ scaled elementary volumes $\sigma_{ijk} V_{ijk}$.
- (iv) The surface vector $\vec{E} \times \vec{F}$ consists of $\binom{n}{2}$ scaled elementary surface vectors $\mu_{ij} \vec{S}_{ij}$.
- (v) $NEV(n) - NES(n) = \begin{cases} -2 & \text{if } n = 3, 4 \\ 0 & \text{if } n = 5 \\ > 0 & \text{if } n > 5 \end{cases}$

(vi) If $n = 1$ (one electric charge q_1), then $k_1 = k_2 = 0$ because the vectors \vec{E} , \vec{F} , \vec{G} are then colinear:

$$\begin{aligned} \vec{E} &= a_1 \vec{R}_1, \\ \vec{F} &= b_1 \vec{R}_1, \\ \vec{G} &= c_1 \vec{R}_1. \end{aligned}$$

(vii) If $n = 2$ (two electric charges q_1 and q_2), then in general $k_1 \neq 0$, but $k_2 = 0$ because the vectors \vec{E} , \vec{F} , \vec{G} are then linear combinations of \vec{R}_1 and \vec{R}_2 , and so are linearly dependent.

4. The golden section and an equipotential surface of two electric charges

In this Section we consider the surface of a fixed potential λ (the λ -potential surface) of two electric charges q_1 and q_2 . Then the formula (2.1) takes the form

$$\lambda R_1 R_2 = q_1 R_2 + q_2 R_1. \tag{4.1}$$

We choose the position vectors \vec{r}_1 and \vec{r}_2 in the following way

$$\begin{cases} \vec{r}_1 = (0, 0, a), \\ \vec{r}_2 = (0, 0, 0). \end{cases}$$

Then the λ -potential surface is invariant under any rotation round the z -axis, and so we can restrict ourselves, for example, to the plane $x = 0$. For this choice

$$\begin{aligned} R_1 &= s = (w^2 - 2az + a^2)^{1/2}, \\ R_2 &= w = (y^2 + z^2)^{1/2}, \end{aligned} \tag{4.3}$$

and the equation (4.1) can be rewritten in the form

$$uv = D_1 v + D_2 u, \tag{4.4}$$

where

$$u = \frac{s}{a}, \quad v = \frac{w}{a}, \tag{4.5}$$

$$D_i = \frac{q_i}{a\lambda} \quad (i = 1, 2). \tag{4.6}$$

Introducing new variables

$$\begin{aligned} \xi &= \frac{1}{2}(u + v), \\ \eta &= \frac{1}{2}(u - v), \end{aligned} \tag{4.7}$$

and notation

$$\begin{aligned} E_1 &= \frac{1}{2}(D_1 + D_2), \\ E_2 &= \frac{1}{2}(D_1 - D_2), \end{aligned} \quad (4.8)$$

we obtain from (4.4) the following equation of hyperbola

$$(\xi - E_1)^2 - (\eta - E_2)^2 = D_1 D_2. \quad (4.9)$$

We call this hyperbola the auxiliary line L_2 of the λ -equipotential surface of q_1 and q_2 . Let

$$P^2 = \begin{cases} D_1 D_2 & \text{if } q_1 q_2 > 0, \\ -D_1 D_2 & \text{if } q_1 q_2 < 0. \end{cases} \quad (4.10)$$

Then the equation (4.9) can be rewritten in the following two forms

$$(\xi - E_1)^2 - (\eta - E_2)^2 = P^2 \quad \text{if } q_1 q_2 > 0, \quad (4.11)$$

$$(\eta - E_2)^2 - (\xi - E_1)^2 = P^2 \quad \text{if } q_1 q_2 < 0. \quad (4.12)$$

Here we restrict ourselves to the first case ($q_1 q_2 > 0$). The other case ($q_1 q_2 < 0$) is discussed in [9]. Let us note that if $q_1 q_2 > 0$, then D_1 and D_2 are always positive. Indeed, by the formula (4.1)

$$q_1 > 0, \quad q_2 > 0 \Rightarrow \lambda > 0, \quad (4.13)$$

$$q_1 < 0, \quad q_2 < 0 \Rightarrow \lambda < 0, \quad (4.14)$$

and so, by (4.6), $D_i > 0$ ($i = 1, 2$).

The formula (4.11) implies that the auxiliary line L_2 can be parametrized in terms of hyperbolic functions $\cosh \gamma$ and $\sinh \gamma$:

$$\begin{cases} \xi - E_1 = P \cosh \gamma, \\ \eta - E_2 = P \sinh \gamma, \\ P = (D_1 D_2)^{1/2}. \end{cases} \quad (4.15)$$

Taking into account the formulas (4.7), (4.8) and (4.15) we can write the old variables u and v in the form

$$\begin{cases} u = D_1 + P e^\gamma, \\ v = D_2 + P e^{-\gamma}. \end{cases} \quad (4.16)$$

The formula (4.16) and the geometric meaning of u and v imply that

$$\begin{cases} \gamma_1 \leq \gamma \leq \gamma_2, \\ D_1 + D_2 + 2P \geq 1. \end{cases} \quad (4.17)$$

The lowest value γ_1 and the highest value γ_2 of γ can be determined from the following two conditions

$$u_1 + 1 = v_1, \tag{4.18}$$

$$u_2 = v_2 + 1, \tag{4.19}$$

where

$$\begin{cases} u_i = u|_{\gamma=\gamma_i}, \\ v_i = v|_{\gamma=\gamma_i}, \\ i = 1, 2. \end{cases} \tag{4.20}$$

Then γ_i, u_i and v_i ($i = 1, 2$) are given by the formulas

$$\begin{aligned} e^{\gamma_1} &= \frac{1}{2P} \{D_2 - D_1 - 1 + [(D_2 - D_1 - 1)^2 + 4P^2]^{1/2}\}, \\ e^{\gamma_2} &= \frac{1}{2P} \{D_2 - D_1 + 1 + [(D_2 - D_1 + 1)^2 + 4P^2]^{1/2}\}, \end{aligned} \tag{4.21}$$

$$\begin{aligned} u_1 &= \frac{1}{2} \{D_1 + D_2 - 1 + [(D_2 - D_1 - 1)^2 + 4P^2]^{1/2}\}, \\ v_1 &= \frac{1}{2} \{D_1 + D_2 + 1 + [(D_2 - D_1 - 1)^2 + 4P^2]^{1/2}\}, \\ u_2 &= \frac{1}{2} \{D_1 + D_2 + 1 + [(D_2 - D_1 + 1)^2 + 4P^2]^{1/2}\}, \\ v_2 &= \frac{1}{2} \{D_1 + D_2 - 1 + [(D_2 - D_1 + 1)^2 + 4P^2]^{1/2}\}. \end{aligned} \tag{4.22}$$

Let us denote that au_1 is the minimal value s_{\min} of s , and av_2 is the minimal value w_{\min} of w .

Let $D_1 = D_2 = 1$. Then the formula (4.22) implies the following proportions of the golden section, [2]:

$$\frac{s_{\min}}{a} = \frac{w_{\min}}{a} = \frac{1}{2}(1 + \sqrt{5}). \tag{4.23}$$

5. Twenty seven straight lines and the auxiliary surface L_3 of three electric charges

Let $C(3)$ denote the following cubic surface (in affine coordinates η_1, η_2, η_3):

$$\eta_1^3 + \eta_2^3 + \eta_3^3 + f = 0. \tag{5.1}$$

The surface $C(3)$ has remarkable properties. In particular $C(3)$ contains 27 straight lines, [4–6]. The configuration of the 27 straight lines of $C(3)$ is invariant with respect to the group G which is isomorphic to the Weyl group

of the exceptional Lie algebra E_6 , i.e. G has 51840 elements and contains a simple group of order 25920, [10].

According to the formula (2.1), the λ -potential surface (equipotential surface of a fixed potential λ) of three electric charges q_1, q_2 and q_3 is given by the equation

$$\lambda R_1 R_2 R_3 = q_1 R_2 R_3 + q_2 R_1 R_3 + q_3 R_1 R_2. \tag{5.2}$$

The auxiliary surface L_3 is by definition the surface obtained from (5.2) by its complexification (R_1, R_2 and R_3 are then complex variables). In [9] we show that

- (i) L_3 gives a geometric insight into a natural parametrization of the real λ -surface of q_1, q_2 and q_3 .
- (ii) L_3 can be related to $C(3)$ via the following change of variables

$$\begin{pmatrix} R_1 - d_1 \\ R_2 - d_2 \\ R_3 - d_3 \end{pmatrix} = \frac{\sqrt{3}}{(d_1 d_2 d_3)^{\frac{1}{3}}} \begin{pmatrix} d_1 & 0 & 0 \\ 0 & d_2 \varepsilon_3 & 0 \\ 0 & 0 & d_3 \varepsilon_3^2 \end{pmatrix} F(3) \begin{pmatrix} \eta_1 \\ \eta_2 \\ \eta_3 \end{pmatrix}, \tag{5.3}$$

where

$$F(3) = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 1 & 1 \\ 1 & \varepsilon_3 & \varepsilon_3^2 \\ 1 & \varepsilon_3^2 & \varepsilon_3 \end{pmatrix}, \tag{5.4}$$

$$d_j = \frac{q_j}{\lambda} \quad (j = 1, 2, 3), \tag{5.5}$$

$$\varepsilon_3 = e^{2\pi i/3}. \tag{5.6}$$

- (iii) L_3 , written in the new variables η_1, η_2, η_3 (defined by (5.3)) intersects $C(3)$ along a hyperbolic helix.
- (iv) The 27 straight lines of $C(3)$ intersect L_3 forming an ‘‘acupuncture’’ of L_3 .
- (v) There is no invertible inhomogeneous linear transformation which transforms L_3 onto $C(3)$.

Finally, let us note that the matrix $F(3)$, given by (5.4), is the so called finite Fourier transform which has remarkable properties collected in [3].

REFERENCES

- [1] J. Mozrzyimas, *Curvature and Torsion of a Current Line of the Field of n Electric Charges*, eds T. Lulek, W. Florek, S. Walcerz, *Symmetry & Structural Properties of Condensed Matter*, World Scientific 1995, p.235
- [2] S. Wajda, *Fibonacci & Lucas Numbers, and the Golden Section*, John Wiley & Sons 1989.
- [3] L. Auslander, R. Tolomieri, *Bull. Amer. Math. Soc.* **1**, 847 (1979).
- [4] Yu.J. Manin, *Kubičeskije Formy*, Nauka, 1972, in Russian.
- [5] B. Serge, *The Non-Singular Cubic Surfaces*, At the Clarendon Press 1942.
- [6] A. Henderson, *The Twenty-Seven Lines Upon the Cubic Surface*, Hafner Publishing Co. 1911.
- [7] M. Berger, B. Gostiaux, *Differential Geometry: Manifolds Curves and Surfaces*, Springer-Verlag 1988.
- [8] H.W. Guggenheimer, *Differential Geometry*, Mc Graw-Hill Book Co. 1963.
- [9] J. Jagielski, J. Mozrzyimas, *Lines and Surfaces Generated by Electric Charges*, to be published.
- [10] R. Harsthorne, *Algebraic Geometry*, Springer-Verlag 1977.