

PROXIMITY EFFECTS IN CeCu<sub>6</sub>/Nb BILAYERS\*A. OTOP<sup>a,b</sup>, S. SÜLLOW<sup>a</sup>, M.B.S. HESSELBERTH<sup>b</sup> AND J. AARTS<sup>b</sup><sup>a</sup>Institut für Metallphysik und Nukleare Festkörperphysik

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We studied the proximity effect in order to probe the response of the heavy-fermion metal CeCu<sub>6</sub> to induced superconductivity. High-quality CeCu<sub>6</sub> films were grown by sputter deposition. Resistivity measurements show that the correlated electron state is formed. Bilayers were prepared of CeCu<sub>6</sub>(75 nm)/Nb( $d_{\text{Nb}}$ ), with the Nb thickness  $d_{\text{Nb}}$  between 10 nm and 50 nm. Measured were the superconducting transition temperature  $T_c$  and the parallel critical fields  $H_{c2}^{\parallel}$ . We find that the interface transparency is high, indicating that Fermi-velocity mismatch effects are not significant, but the coherence length of the induced superconductivity is small, which is probably due to the very low Fermi-velocity in the heavy fermion metal.

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**1. Introduction**

In simple normal metals such as Au or Pd, questions about their ability to sustain superconducting correlations or unbalanced spins are of much interest. Similar questions can be raised about a metallic ground state with strong intrinsic correlations such as in heavy fermion (HF) systems. Specifically, if a metal is in contact with a superconductor, it can be asked what the mechanism is for pair breaking of Cooper pairs diffusing out of the superconductor. For normal metals (N) the mechanism is dephasing due to a finite temperature; for ferromagnets (F) it is the effect of the exchange field; for a HF ground state the answer is not known, and could involve *e.g.* spin fluctuations. An intrinsic problem for such studies is that the material has to be available in thin film form. Here we present the first results of a

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study of the proximity effect in a system consisting of thin films of the HF material  $\text{CeCu}_6$  combined with thin films of a superconductor (S), in this case Nb. Among the different HF systems,  $\text{CeCu}_6$  is of particular interest since it does not order magnetically down to mK temperatures [1, 2].

## 2. Sample preparation

Sets of  $\text{CeCu}_6$  films and bilayers of  $\text{CeCu}_6/\text{Nb}$  were made by DC-magnetron sputtering. We used an ultra high vacuum system with a background pressure of the order of  $5 \times 10^{-10}$  mbar, and an Ar sputtering pressure of  $2.5 \times 10^{-3}$  mbar. Crystalline  $\text{CeCu}_6$  was grown as reported before [3], at a temperature of  $350^\circ \text{C}$  using Si-substrates with amorphous  $\text{Si}_3\text{N}_4$  buffer layers to prevent Cu diffusion at those temperatures.

The Nb was deposited on top of the  $\text{CeCu}_6$  after cooling the substrate holder with cold nitrogen gas to close to room temperature. Composition, thickness and crystallinity of the films were determined by Rutherford backscattering (RBS) measurements together with X-ray diffraction measurements at low and high angles. The RBS measurements show good agreement with the expected stoichiometry for  $\text{CeCu}_6$  and no diffusion is found either of Ce or Cu into the substrate or of Nb into the  $\text{CeCu}_6$ . Fig. 1 shows part of the RBS spectrum for  $\text{CeCu}_6$  (75 nm)/Nb(15 nm) on a Si/ $\text{Si}_3\text{N}_4$  substrate and a fit of the data without taking any diffusion into account.

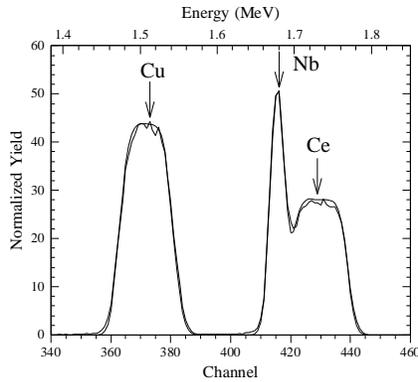


Fig. 1. RBS measurement using  $^4\text{He}$ -ions of 2 MeV on a sample Si/ $\text{Si}_3\text{N}_4$ / $\text{CeCu}_6$ (75 nm)/Nb(15 nm). The different elements are indicated. The thin smooth line is a fit to the measured curve.

Fig. 2(a) shows the electrical resistivity  $\rho$  as a function of temperature  $T$  for two single films of  $\text{CeCu}_6$  with different thickness. A clear maximum in  $\rho(T)$  is observed at  $T_{\text{max}} \approx 5 \text{ K}$ , similar to what is found for bulk material [4]

and in previous investigations on sputtered CeCu<sub>6</sub> films [3]. We decided to use 75 nm thick CeCu<sub>6</sub> layer, which with  $T_{\max} = 4$  K suggest only little deviation from the bulk properties.

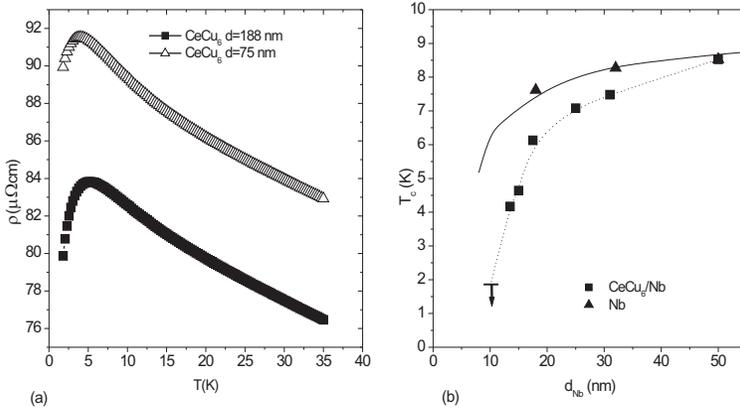


Fig. 2. (a) Resistivity as a function of temperature for CeCu<sub>6</sub> films of 75 nm and 188 nm. (b) Critical temperature as a function of Nb thickness  $d_{\text{Nb}}$  for bilayers CeCu<sub>6</sub>/Nb and single Nb films. The solid line represents the typical behavior for Nb single films from different deposition systems, triangles show values for Nb films made in the same system as the CeCu<sub>6</sub>/Nb bilayers, the dotted line is a guide to the eye.

### 3. CeCu<sub>6</sub>/Nb bilayers; results and discussion

The superconducting transition temperature  $T_c$  was determined from resistivity measurements on a set of CeCu<sub>6</sub>(75 nm)/Nb( $d_{\text{Nb}}$ ) with variable Nb thickness  $d_{\text{Nb}}$  in the range 10 nm to 50 nm. As seen in Fig. 2(b),  $T_c$  shows a strong suppression, with a critical thickness  $d_{\text{cr}}$  for onset of superconductivity reached around 12 nm. This would be equivalent to 24 nm in a trilayer configuration, and of a similar magnitude as found in Fe/Nb/Fe trilayers [5]. It has to be kept in mind however that single thin Nb films also show a decrease of  $T_c$  at small thickness (see Fig. 2(b)), which makes the correct determination of  $d_{\text{cr}}$  more difficult. The result shows, however, that the interface does allow particle exchange, opposite to what might be expected from the huge Fermi velocity mismatch in the system of order  $10^3$ .

Also the parallel critical field  $H_{c2}^{\parallel}$  behaves differently than expected for simple N/S systems; as shown in Fig. 3(a) no dimensional crossover (DCO) is found at any Nb thickness, in contrast to the behavior of *e.g.* Cu/Nb, as witnessed by a sample Cu(75 nm)/Nb(15 nm) prepared for comparison (Fig. 3(b)).

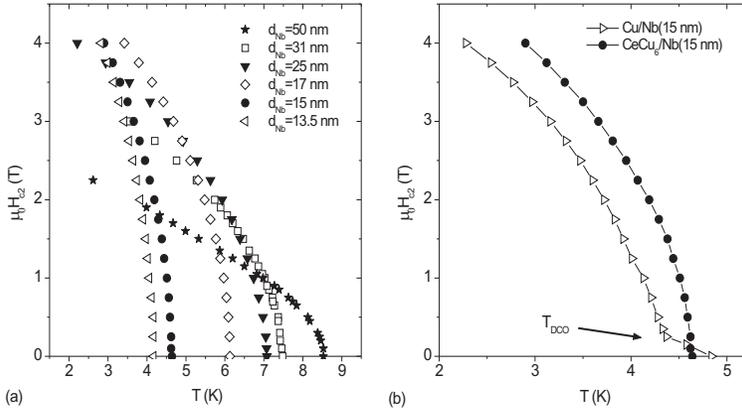


Fig. 3. Upper parallel critical field as a function of temperature for: (a) bilayers CeCu<sub>6</sub>(75 nm)/Nb,  $d_{Nb}$  as indicated; (b) bilayers of CeCu<sub>6</sub>(75 nm)/Nb(15 nm) and Cu(75 nm)/Nb(15 nm) with  $T_{DCO}$  indicated.

Absence of the DCO means that no appreciable amount of superconductivity leaks into the HF metal; together with the finite transparency this means either appreciable pair breaking by the HF metal, or a very small diffusivity due to the low Fermi velocity  $v_F$ . Whether a small  $v_F$  alone is able to explain the experiments is presently under investigation.

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