AC SUSCEPTIBILITY STUDIES OF Zn-DOPED MAGNETITE SINGLE CRYSTALS*

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We present selected results of systematic studies of AC susceptibility in the series of single crystalline Fe$_{3-x}$Zn$_x$O$_4$ ($x < 0.04$) and Fe$_{3(1-\delta)}$O$_4$ samples. Two sets of anomalies were found: the first one connected with the Verwey transition and the second consisting of two different effects, one at 26 K, visible only for stoichiometric samples, and the other at 50 K. While the temperature position of the first anomaly (26 K) does not depend on $f$, the effect at 50 K shifts to higher temperatures with increasing frequency. Both effects gradually move to lower temperatures and finally disappear with increasing $x$ and $\delta$. Qualitatively similar results were reported from Magnetic After Effect (MAE) technique, that records effects with the relaxation time $10^4$ higher than that characteristic for our technique.

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1. Introduction

The aim of the present report is to study the effect of Zn doping and non-stoichiometry on the low temperature (mainly $T < 70$ K) magnetic properties of magnetite single crystals as visualised by $\chi_{AC}$. By this means we intend to complement earlier literature reports [1] devoted only to pure magnetite.

Apart from its technological importance magnetite, Fe$_3$O$_4$, is known due to the Verwey transition at $T_V \approx 120$ K. At high $T$ magnetite is a disordered electron system in which “additional” (beyond Fe$^{3+}$ ion core) electrons transfer between adjacent octahedral Fe. Although, recently, the common picture of the transition as the freezing out of mobile strongly correlated electrons have been questioned [2, 3], the ionic ordering is rather well experimentally evidenced (see the discussion in [4]). With still existing controversies over Fe$^{3+}$–Fe$^{2+}$ exchange we have decided to study $\chi_{AC}$, the technique reported to have a close link with the electron exchange phenomenon [4].

2. Results and discussion

The in-phase component $\chi'$ measured at $f = 125$ Hz and $H_{AC} = 1$ Oe for all studied samples (see Fig. 1) falls considerably while $T$ approaches $T < T_V$ region, with the simultaneous increase in the loss $\chi''$ contribution as already reported [1]. Additionally there is the anomaly in $\chi$ below 60 K in all samples: $\chi'$ drops with the simultaneous peak in $\chi''$ (see Fig. 2), proving that the domain walls can not follow $H_{AC}$ changes. This effect is virtually the same for field cooled samples, and similar, although smaller, for field cooled

![Fig. 1. Dispersive part $\chi'$ of AC susceptibility for Fe$_{3-x}$Zn$_x$O$_4$ and Fe$_{3(1-\delta)}$O$_4$ single crystals.](image)
and subsequently in-field measured stoichiometric magnetite. As presented in Fig. 2 linear relation $\log(f)$ vs. $1/T$ holds, indicating thermally activated relaxation process.

![Graph](image)

Fig. 2. AC $\chi'$ and $\chi''$ for stoichiometric Fe$_3$O$_4$; the inset proves that the peak in $\chi''$ at $T = 50 - 70$K is a manifestation of relaxation effects described roughly by the formula $f = (1/\tau_0) \exp(-Q/\kappa T)$. Values of $Q$ equal to 98, 26 and 42 (in meV), and $\tau_0$ equal to $4.5 \times 10^{-13}$s, $6.7 \times 10^{-8}$s and $1.3 \times 10^{-9}$s were found for $x = 0, 0.0072$ and 0.011, respectively.

Interestingly, the qualitatively similar results have been noticed by MAE [6], the technique that records effects with the relaxation time $10^4$ higher than that characteristic for our technique. In particular, the relaxation gap was observed there between 30 and 50 K, i.e. at the $T$ range where our anomalies are found. Thus both our results and MAE data prove the existence of the relaxation effects at ca. 60 K for stoichiometric magnetite and at lower $T$ for doped and nonstoichiometric samples with very wide relaxation time spectrum.

In case of stoichiometric magnetite the drop in $\chi'$ at 60 K is further followed by the increase of $\chi'$, exhibiting characteristic subtle structure until the high-$T$ $\chi'$ value is reached, suggesting that the wall movement once more takes place at lower $T$. Since this anomaly is absent in the results of in-field measurements, the external magnetic field apparently curtails the processes that enable the wall movement at temperatures below 26 K. Moreover, the effect at $T = 26$ K is of different kind than that at 60 K since no $f$ dependence was seen. The corresponding peak is again observed in the results from MAE [4, 6]; there, however, the peak shows typical relaxation behaviour and is linked to intraatomic Fe$^{2+}$ electron excitations. If the relaxation times
of these processes were of the order of the characteristic observation time for MAE, i.e. 1 to 180 s, our technique should see the static arrangement resulting in $\chi'^{m} = 0$. Thus, the fact that we still observe the anomaly suggests that some more complicated phenomena are taking place.

In conclusion, we have measured temperature, frequency and $H_{AC}$ dependence of AC susceptibility of single magnetite based crystals. We have found the usual drop in $\chi'$ at $T_{y}$ followed at lower $T$ by the anomaly, that was explained by the slowing domain wall movement. The additional anomaly at $T = 26K$, observed only for stoichiometric magnetite, does not show the typical $f$ dependence, suggesting that some complicated effects are taking place. Since both anomalies are probably linked to the octahedral iron energy level fluctuations, AC susceptibility studies, together with MAE, may help to clarify the still open problem of the processes causing the Verwey transition.

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REFERENCES