THERMAL EXPANSION OF THE SKUTTERUDITE SUPERCONDUCTOR PrOs$_4$Sb$_{12}$

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The recently discovered Pr-based skutterudite superconductor PrOs$_4$Sb$_{12}$ was studied by means of low-temperature measurements of the thermal expansion coefficient $\alpha$ along the cubic (100) direction of a single crystal sample in magnetic fields up to 8 T. Two superconducting phase transitions were found with critical temperatures $T_{c1} = 1.84$ K and $T_{c2} = 1.71$ K. Their hydrostatic pressure dependencies calculated using the Ehrenfest relation are $-250 \pm 50$ mK/GPa and $-450 \pm 70$ mK/GPa, respectively. For $B \geq 5$ T, we observe a thermodynamic phase transition which is presumably due to the onset of quadrupolar order. The transition temperature shifts with increasing $B$ towards higher temperatures. The strongly enhanced value of the Grüneisen parameter $\Gamma$ provides clear evidence for the occurrence of a heavy-fermion normal state. The pronounced change of $\Gamma$ at $T_c$ provides evidence for the formation of Cooper pairs out of the heavy quasiparticles.

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A large variety of different ground states have been found in the class of filled skutterudite compounds MT$_4$X$_{12}$ (M = lanthanide or actinide metal, T = Fe, Ru, Os, X = P, As, Sb); e.g., superconductivity [LaFe$_4$P$_{12}$, LaRu$_4$Sb$_{12}$], Kondo insulator (CeFe$_4$P$_{12}$, UFe$_4$P$_{12}$), heavy-fermion (CeFe$_4$Sb$_{12}$, YbFe$_4$Sb$_{12}$), and non-Fermi liquid (CeRu$_4$Sb$_{12}$) behaviour [1]. They crystallize in the cubic LaFe$_4$P$_{12}$ structure with TX$_3$ cages filled by lanthanide or actinide atoms [2].

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A new member of the class of filled skutterudites, PrOs$_4$Sb$_{12}$, exhibits heavy-fermion behaviour and superconductivity and is considered to be the first Pr-based heavy-fermion superconductor [3]. The susceptibility data are best described by a crystalline-electric field scheme with a non-magnetic $I_3$ ground state which causes an electric quadrupole moment separated from a $I_y$ triplet [1]. The large size of the jump in the specific heat at $T_c \sim 1.8$ K, $\Delta C/T_c$, which is of the same order of magnitude as the Sommerfeld coefficient $\gamma_0$ suggests heavy-fermion superconductivity in PrOs$_4$Sb$_{12}$ [3].

Here we report on ultrahigh resolution thermal expansion measurements on single crystal PrOs$_4$Sb$_{12}$ down to 50 mK and in fields up to 8 T. The linear thermal expansion coefficient $\alpha$ is defined as $\alpha = l^{-1} \partial l / \partial T$ where $l$ denotes the sample length. For a cubic system, the volume thermal expansion coefficient $\beta(T)$ is given by $3 \alpha(T)$. Two samples were mounted on top of each other to improve the resolution. The upper part of Fig. 1 shows the volume thermal expansion coefficient $\beta$ vs $T$. Two superconducting phase transitions lying close to each other separated by a small plateau are resolved with transition temperatures $T_{c1} = 1.84$ K and $T_{c2} = 1.71$ K (cf. inset of Fig. 1). This might be caused by different transition temperatures of both crystals. However, specific heat measurements on one sample also showed two transitions with similar $T_c$ values [1]. Further arguments in favor of two transitions are given by the calculation of the hydrostatic pressure dependence of $T_c$ by using the Ehrenfest relation, $\partial T_c / \partial p = V_{m0} T_c \Delta \beta / \Delta C$. Here $V_{m0} = 2.42 \times 10^{-4}$ m$^3$/mol and $\Delta C$ denote the molar volume and the jump height of $C$ at $T_c$, respectively. $\Delta \beta(T_c)$ is estimated by an equal-areas

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Fig. 1. (a) $T$ dependence of the volume thermal expansion coefficient $\beta$ (a) and Grüneisen parameter $\Gamma$; (b) for PrOs$_4$Sb$_{12}$. The arrows mark the superconducting phase transitions in $\beta$ vs $T$ enlarged in the inset.
construction. For the transition temperature $T_c$, we deduce a hydrostatic pressure dependence of $-250 \pm 50 \text{ mK/GPa}$, which is slightly higher than the value obtained by resistivity measurements under hydrostatic pressure [1]. The pressure dependence of $T_c$ is found to be $-450 \pm 70 \text{ mK/GPa}$. Assuming only one transition in $\beta$ and $C$ a value of $-510 \pm 100 \text{ mK/GPa}$ would be derived which is three times higher than that of the resistivity measurements under hydrostatic pressure. Thus, the splitting of the transition seems to be intrinsic, although its origin is still unknown. The occurrence of two transitions could hint to unconventional superconductivity with a multicomponent order parameter as found for UPt$_3$ [4].

In the lower plot of Fig. 1 the Grüneisen parameter $\Gamma vs T$ for PrOs$_4$Sb$_{12}$ is shown. The Grüneisen parameter is defined by $\Gamma = V_{mol}/\kappa_T \beta /C$, where the value $4.7 \times 10^{-12} \text{ Pa}$ for the isothermal compressibility $\kappa_T$ for PrFe$_4$P$_{12}$ was used [5]. The corresponding specific heat data are taken from [3]. The Grüneisen parameter for PrOs$_4$Sb$_{12}$ at $T = 6 \text{ K}$ exceeds the typical values for ordinary metals ($\Gamma \sim 1$) by a factor of 10 and decreases linearly upon lowering $T$. The highly enhanced Grüneisen parameter in the normal state provides strong evidence that PrOs$_4$Sb$_{12}$ forms a heavy-fermion state at low temperatures. The strong change of the Grüneisen parameter at $T_c$ gives clear evidence for the occurrence of heavy-fermion superconductivity formed by the heavy quasiparticles.

The magnetic field dependence of the linear thermal expansion coefficient $\alpha$ is shown in Fig. 2. For $B > 1 \text{ T}$, the superconducting jump is not resolvable anymore although the upper critical field is found to be $2.2 \text{ T}$ [3]. The minimum structure in $\alpha$ whose minimum temperature $T_{\text{min}}$ shifts to-

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Fig. 2. Linear thermal expansion coefficient $\alpha vs T$ in different fields $B \leq 8 \text{ T}$. Inset: $B$-$T$ diagram with superconducting ($T_c$), and, presumably, quadrupolar ordering ($T_Q$) phase transitions. Lines are guides to the eye.
wards lower temperatures upon increasing the field to 5 T corresponds to
the Schottky anomaly found in specific heat. The feature can also be fitted
by a Schottky term with an additional term taking into account the T de-
pendence of the Grüneisen parameter. The low T up-turns are caused by
the hyperfine splitting of the Pr nucleus.

Upon increasing the magnetic field to \( B \geq 5 \) T, a second phase transition
occurs. The transition temperature \( T_Q \) and the jump height in \( \alpha(T) \) increase
upon increasing the field. The resulting \( B-T \) phase diagram is shown in the
Inset of Fig 2. The high field transition at \( T_Q \) is presumed to be a transition
into quadrupolar ordering. A similar field dependence of the phase-transition
line was found for PrPb\textsubscript{3} which is known to exhibit quadrupolar order \[6\].

To summarize, we confirmed the occurrence of heavy-fermion supercon-
ductivity in PrOs\textsubscript{4}Sb\textsubscript{12} by an analysis of the low-temperature thermal ex-
ansion and Grüneisen parameter. Two superconducting phase transitions
were found in thermal expansion and specific heat measurements. This
points to unconventional superconductivity in PrOs\textsubscript{4}Sb\textsubscript{12}. For \( B \geq 5 \) T, the
quadrupolar moments presumably order at \( T_Q \).

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