

# NON-STRANGE DIBARYONS STUDIED IN COHERENT DOUBLE NEUTRAL-MESON PHOTOPRODUCTION ON THE DEUTERON\*

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We have investigated the  $\gamma d \rightarrow \pi^0 \pi^0 d$  reaction to study dibaryon resonances. The total cross section as a function of the  $\gamma d$  center-of-mass energy shows resonance-like behavior peaked at around 2.47 and 2.63 GeV. The measured angular distribution of deuteron emission is rather flat, which cannot be reproduced by kinematics for quasi-free  $\pi^0 \pi^0$  production with deuteron coalescence. A clear peak is observed at 2.14 GeV in the  $\pi^0 d$  invariant-mass distributions. The present work shows a sequential process  $\gamma d \rightarrow R_{\text{IS}} \rightarrow \pi^0 R_{\text{IV}} \rightarrow \pi^0 \pi^0 d$  is dominant with two 2.47- and 2.63-GeV isoscalar dibaryons ( $R_{\text{IS}}$ ) and a 2.14-GeV isovector dibaryon ( $R_{\text{IV}}$ ).

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## 1. Introduction

The structure of hadrons is one of the most important subjects to be studied in the non-perturbative domain of quantum chromodynamics. A  $B = 2$  system (dibaryon) [1] is an interesting object to study its basic configuration from a molecule-like state consisting of two baryons to a spatially-compact hexaquark hadron state. Understanding dibaryons would not only give a clue to the solution of the current problem in hadron physics, but also provide an insight into the nuclear equation of state and the interior of a neutron star [2].

Recent observations of the  $d^*(2380)$  dibaryon [3, 4] have opened the door to study non-strange dibaryons. It is important to establish the excitation spectrum of dibaryons to understand their internal structure. We study the  $\gamma d \rightarrow \pi^0 \pi^0 d$  reaction to observe intermediate dibaryon states. Possible production mechanisms for the  $\gamma d \rightarrow \pi^0 \pi^0 d$  reaction are classified as follows:

1. Dibaryon: sequential  $\pi^0 \pi^0$  emission from the deuteron with intermediate isoscalar dibaryon  $R_{IS}$  and isovector dibaryon  $R_{IV}$  ( $\gamma d \rightarrow R_{IS} \rightarrow \pi^0 R_{IV} \rightarrow \pi^0 \pi^0 d$ ).
2. QF- $\pi\pi$ :  $\pi^0 \pi^0$  is photoproduced on the quasi-free (QF) participant nucleon  $N_p$ , followed by coalescence of  $N_p$  and spectator nucleon  $N_s$  into a deuteron.
3. QF- $\pi$ :  $\pi^0$  is photoproduced on  $N_p$ , followed by coalescence of  $N_p$  and  $N_s$  into  $R_{IV}$ , finally  $R_{IV}$  decays into  $\pi^0 d$ .
4. Direct- $\pi\pi$ :  $\pi^0 \pi^0$  is directly photoproduced from the deuteron.

In the QF- $\pi\pi$  mechanism, the second  $\pi^0$  should be emitted to compensate for the momentum given to  $N_p$  by the first emitted  $\pi^0$  to coalesce into a deuteron. In this case, the angular distribution of deuteron emission in the  $\gamma d$  center-of-mass (CM) frame shows strongly backward peaking. As for the QF- $\pi$  mechanism, the condition to coalescence of  $R_{IV}$  makes the distribution sideway peaking at the incident photon energy around 1 GeV. A completely different rather flat distribution is obtained in the dibaryon and direct- $\pi\pi$  mechanisms.

## 2. Experiment

The total and differential cross sections were measured for the  $\gamma d \rightarrow \pi^0 \pi^0 d$  reaction using an energy-tagged bremsstrahlung photon beam ranging from 0.75 to 1.15 GeV [5] at the Research Center for Electron Photon Science (ELPH), Tohoku University, Japan. The target used in the experiment was liquid deuterium with a thickness of 45.9 mm. All the final-state particles in the  $\gamma d \rightarrow \pi^0 \pi^0 d$  reaction were measured with the FOREST detector [6].

FOREST consists of three different electromagnetic calorimeters (EMCs), and a plastic-scintillator hodoscope (PSH) is placed in front of each EMC to identify charged particles.

### 3. Results

The analysis of the  $\gamma d \rightarrow \pi^0 \pi^0 d$  reaction was made in the same way as in Ref. [7]. We selected the events containing four neutral particles and a charged particle, and applied a kinematic fit with six constraints: four-momentum conservation, and every  $\gamma\gamma$ -invariant mass being the  $\pi^0$  mass. Events in which the  $\chi^2$  probability is higher than 0.4 were selected to reduce those from background processes. Figure 1 shows the total cross section  $\sigma$  as a function of the  $\gamma d$  CM energy  $W_{\gamma d}$ . The excitation function is not monotonically increasing but shows resonance-like behavior peaked at around 2.47 and 2.63 GeV.

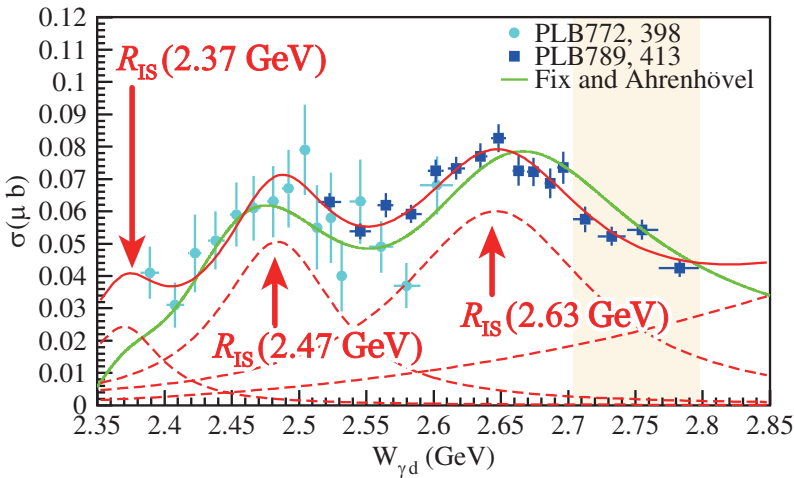


Fig. 1. (Color online) Total cross section  $\sigma$  as a function of  $W_{\gamma d}$ . The squares (blue) and circles (cyan) show  $\sigma$ s obtained in the previous [7] and in the present work [8], respectively. The horizontal error of each data point corresponds to the coverage of the incident photon energy, and the vertical error shows the statistical error of  $\sigma$ . The solid curve (grey/green) shows a theoretical calculation given in Ref. [9] based on the QF- $\pi\pi$  mechanism. The solid curve (dark grey/red) shows the fitted function expressed by a sum of three Breit-Wigner peaks and phase-space contributions. Each contribution is shown as a dashed curve (dark grey/red).

A naive interpretation of this behavior may be a QF excitation of the nucleon, followed by coalescence into the deuteron. However, a rather-flat angular distribution of deuteron emission completely differs from the QF- $\pi\pi$  and QF- $\pi$  mechanisms. In addition, the  $\pi^0 d$  invariant-mass distributions

shows a peak at 2.14 GeV. The present work suggests a sequential process  $\gamma d \rightarrow R_{\text{IS}} \rightarrow \pi^0 R_{\text{IV}} \rightarrow \pi^0 \pi^0 d$  is dominant with two 2.47- and 2.63-GeV  $R_{\text{IS}}$  and a 2.14-GeV  $R_{\text{IV}}$ .

#### 4. Summary

The total and differential cross sections have been measured for the  $\gamma d \rightarrow \pi^0 \pi^0 d$  reaction at  $E_\gamma = 0.75\text{--}1.15$  GeV. The measured angular distribution of deuteron emission is rather flat, suggesting that a sequential process  $\gamma d \rightarrow R_{\text{IS}} \rightarrow \pi^0 R_{\text{IV}} \rightarrow \pi^0 \pi^0 d$  is dominant. The total cross section as a function of  $W_{\gamma d}$  shows isoscalar dibaryons  $R_{\text{IS}}$  with masses of 2.47 and 2.63 GeV. The  $\pi^0 d$  invariant-mass distributions show an isovector dibaryon  $R_{\text{IV}}$  with a mass of 2.14 GeV. The details of the analysis and discussion can be found elsewhere [8].

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#### REFERENCES

- [1] H. Clement, *Prog. Part. Nucl. Phys.* **93**, 195 (2017).
- [2] A. Akmal, V.R. Pandharipande, D.G. Ravenhall, *Phys. Rev. C* **58**, 1804 (1998).
- [3] M. Bashkanov *et al.* [CELSIUS/WASA Collaboration], *Phys. Rev. Lett.* **102**, 052301 (2009).
- [4] P. Adlarson *et al.* [WASA-at-COSY Collaboration], *Phys. Rev. Lett.* **106**, 242302 (2011).
- [5] T. Ishikawa *et al.*, *Nucl. Instrum. Methods Phys. Res. A* **622**, 1 (2010); **811**, 124 (2016); Y. Matsumura *et al.*, *Nucl. Instrum. Methods Phys. Res. A* **902**, 103 (2018); Y. Obara *et al.*, *Nucl. Instrum. Methods Phys. Res. A* **922**, 108 (2019).
- [6] T. Ishikawa *et al.*, *Nucl. Instrum. Methods Phys. Res. A* **832**, 108 (2016).
- [7] T. Ishikawa *et al.*, *Phys. Lett. B* **772**, 398 (2017).
- [8] T. Ishikawa *et al.*, *Phys. Lett. B* **789**, 413 (2019).
- [9] A. Fix, H. Arenhövel, *Eur. Phys. J. A* **25**, 115 (2005).