

FIRST ESTIMATION OF THE FISSION DYNAMICS OF THE SPECTATOR CREATED IN HEAVY-ION COLLISIONS*

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In peripheral high-energy heavy-ion collisions only parts of colliding nuclei interact leading to the production of a fireball. The remnants of such nuclei are called spectators. We estimated the excitation energy of nuclear remnants as a function of impact parameter. Their excitation energy is of the order of 100 MeV. The dynamical evolution of hot nuclei is described by solving a set of Langevin equations in four-dimensional collective coordinate space. The range of nuclear masses and excitation energies suits very well the ability of our model. Thus, for the first time, we investigate dynamically the fission and evaporation channels in deexcitation of the spectators produced in heavy-ions collision.

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1. Introduction

The physics of high-energy heavy-ion collisions has developed very effectively for many years. At CERN, *e.g.* $^{208}\text{Pb}+^{208}\text{Pb}$ collisions were considered at the SPS and LHC. Two lead nuclei flying with velocities comparable with light velocity produce the fireball in the place of collision, but depending on the centrality of the collision, the Pb parts which are not participating in the collision (called “spectators”) follow their initial path and deexcite later on. There are many experimental and theoretical [1, 2] investigations of the processes connected to the fireball — a piece of quark–gluon plasma.

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The density of the matter in the contact area is so high that the quark–gluon plasma is created and the high-energy particles such as pions, kaons, and others emerge in the hadronization process. This part of the reaction requires dedicated models and methods that are beyond the scope of the present discussion.

The interesting issue is the behavior of the spectators produced as the remnants of the collision [3]. There are not so many experimental observables published up to now, which can help with constraining the physics of spectator evolution. The fast collision of two spherical ions provides exotic, for nuclear physics, shapes of the remnants which are strongly dependent on the impact parameter. The deformation energy of the produced system turns into excitation energy. Thus, we treat the spectator as an excited nucleus, which can deexcite by shape changes or emission of light particles and/or γ -rays. This allows to apply the stochastic approach based on the solving transport equations in collective coordinate space [4–7].

2. Theoretical method

The stochastic approach assumes that the evolution of the nucleus in collective coordinate space is similar to the Brownian motion of particles which interact stochastically involving a large number of internal degrees of freedom, which constitutes the surrounding “heat bath”. The motion of the nucleus is slowed down by a friction, considered here as the hydrodynamical friction force. The transport coefficients are derived from the random force averaged over a time larger than the collision time scale between collective and internal variables. The Gaussian white noise is taken for the random part, producing fluctuations of the physical observables such as the mass/charge distribution of the fission fragments or evaporation residua. The coupled Langevin equations describing the fission reaction have the form of

$$\begin{aligned} \frac{dq_i}{dt} &= \mu_{ij} p_j, \\ \frac{dp_i}{dt} &= -\frac{1}{2} p_j p_k \frac{\partial \mu_{jk}}{\partial q_i} - \left(\frac{\partial F}{\partial q_i} \right)_T - \gamma_{ij} \mu_{jk} p_k + \theta_{ij} \xi_j(t), \end{aligned} \quad (1)$$

where \mathbf{q} are the collective coordinates, \mathbf{p} are the conjugate momenta, $F(\mathbf{q}, K) = V(\mathbf{q}, K) - a(\mathbf{q})T^2$ is the Helmholtz free energy, $V(\mathbf{q})$ is the potential energy, $m_{ij}(\mathbf{q})$ ($\|\mu_{ij}\| = \|m_{ij}\|^{-1}$) is the tensor of inertia, and $\gamma_{ij}(\mathbf{q})$ is the friction tensor. The $\xi_j(t)$ is a random variable satisfying the relations

$$\begin{aligned} \langle \xi_i \rangle &= 0, \\ \langle \xi_i(t_1) \xi_j(t_2) \rangle &= 2\delta_{ij} \delta(t_1 - t_2). \end{aligned} \quad (2)$$

Thus, a Markovian approximation is assumed to be valid. The strength of the random force θ_{ij} is given by the Einstein relation $\sum \theta_{ik}\theta_{kj} = T\gamma_{ij}$. The temperature of the “heat bath” T is determined by the Fermi-gas model formula $T = (E_{\text{int}}/a)^{1/2}$, where E_{int} is the internal excitation energy of the nucleus and $a(\mathbf{q})$ is the level-density parameter.

The energy conservation law is used in the form of

$$E^* = E_{\text{int}} + E_{\text{coll}} + V + E_{\text{evap}}(t) \quad (3)$$

at each time step during a random walk along the Langevin trajectory in the collective coordinate space. The total excitation energy of the nucleus is E^* and the collective kinetic energy is taken as: $E_{\text{coll}} = 0.5 \sum \mu_{ij} p_i p_j$. The energy carried away by the evaporated particles by the time t is marked as $E_{\text{evap}}(t)$.

The collective coordinate space is constructed with three deformation parameters: elongation of the nucleus (q_1), its neck (q_2) and octupole shape (q_3) based on the “funny hills” parametrization proposed initially by Brack [8] and the fourth degree of freedom is the projection of the angular momentum vector on the fission axis. This four collective coordinates allow to describe a rich ensemble of shapes and, additionally, it allows to orient the nucleus in the laboratory space, because K is the projection of the spin on the fission axis.

The 4D Langevin method describes the behavior of the hot nucleus in the reference frame connected with the center-of-mass of each nucleus, thus in the case of the fast moving systems in the laboratory frame, additional transformations are necessary as the velocities of nuclei are in relativistic regime.

The initial conditions to start evolution of the nucleus are usually the mass and charge of the beam and target nuclei, their excitation energy and maximal angular momentum, which can be produced in the reaction.

Our idea of investigating the spectator with 4D Langevin method is very recent. There are no direct measurements that could identify any fission/evaporation reactions occurring in the remnant of heavy-ions collisions, but new developments in the detection techniques [9] could make possible the investigations of new effects accompanying the main reaction.

The estimation of the shape of the spectator just after the ion collision is the gateway to introduce the fission dynamics. We propose three possible scenarios of the non-central collisions. Let us assume that the beam and target nuclei have the spherical forms in their rest frame. When they fly in opposite directions, they collide and graze each other. The central part produces the quark–gluon plasma and all kinds of high-energy particle phenomena. Remnants could have extremely exotic forms and sometimes it is assumed that they are a collection of free nucleons. Here, we discuss the

following scenarios of initial shape production: “sphere–plane” means the sphere cut by the plane at various impact parameters; “sphere–sphere” gives the shape of the spectator which is sphere but with missing spherical cap; “sphere–cylinder” is the globe shot by the bullet.

These three scenarios give shapes, very strange for nuclear physics. These exotic deformations are impossible to describe with any standard nuclear shape parameterizations. Thus, we make an assumption that the time to change nuclear deformation from initial (complicated) to spherical is small compared to the time scale of the rest of the dynamics. This rough assumption will be investigated in the future in detail. The difference in the surface energy for spherical and deformed nucleus gives the preliminary excitation energy of the spectator.

The potential energy of the nucleus is obtained with the Lublin–Strasbourg Drop formula [10, 11]

$$\begin{aligned} E_{\text{lsd}}(Z, N; q) &= E_{\text{vol}} + E_{\text{surf}} + E_{\text{curv}} + E_{\text{Coul}} & (4) \\ &= b_{\text{vol}} (1 - \kappa_{\text{vol}} I^2) A + b_{\text{surf}} (1 - \kappa_{\text{surf}} I^2) A^{2/3} B_{\text{surf}}(q) \\ &\quad + b_{\text{curv}} (1 - \kappa_{\text{curv}} I^2) A^{1/3} B_{\text{curv}}(q) + \frac{3}{5} e^2 \frac{Z^2}{r_0^{\text{ch}} A^{1/3}} B_{\text{Coul}}(q), & (5) \end{aligned}$$

where $I = (N - Z)/A$. The geometric factors: B_{surf} , B_{curv} and B_{Coul} are defined as

$$B_{\text{surf}} = \frac{S(\text{def})}{S(0)}, \quad B_{\text{curv}} = \frac{C(\text{def})}{C(0)} \quad \text{and} \quad B_{\text{Coul}} = \frac{E_{\text{Coul}}(\text{def})}{E_{\text{Coul}}(0)} \quad (6)$$

that are the ratios between the surface (S)/curvature (C)/Coulomb energy (E_{Coul}) of the deformed and spherical nucleus. The details are presented *e.g.* in [10]. Thus for each of the scenarios, the surface and the volume of the spectator were calculated as functions of the impact parameter. The volume of the spectator and its mass were obtained using the geometrical properties. Assuming $N/Z_{\text{Pb}} = N/Z_{\text{spectator}}$ also the charge was assessed. The deformation together with the mass and charge gives the surface energy estimation.

3. Results

The surface energy of the remnant just after the collision reduced by the energy of the spherical nucleus gives the deformation energy which is presented in Fig. 1. The deformation energy depends on the initial condition. The highest deformation energies (almost 500 MeV) are predicted for the case of “sphere–sphere” grazing as the shapes are more curved than in other cases. The smallest energies are obtained for the sphere cut by the

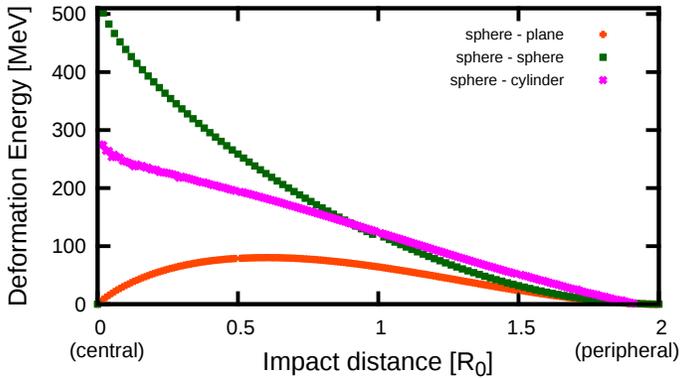


Fig. 1. (Color online) The deformation energy for three scenarios: “sphere–sphere”, “sphere–plane” and “sphere–cylinder”. The impact distance is normalized to the radius of the spherical lead nucleus.

plane as the forms are less curved and the surface of the deformed nucleus is closer to a spherical shape. The impact parameter is a good observable to discuss the central–peripheral collision, as it is presented in Fig. 1, but for fission dynamics calculation, the mass and charge of the spectators are necessary. Thus, in Fig. 2, the deformation energy is shown as a function of the mass of the remnant. The deformation energy has the main contribution to the excitation energy of the spectator. The black crosses mark the nuclei taken in further calculations. The 4D Langevin method is dedicated to describe the fission dynamics of the hot nuclei at excitation energies $E^* = 50\text{--}250$ MeV and the driving forces are derived macroscopically. The

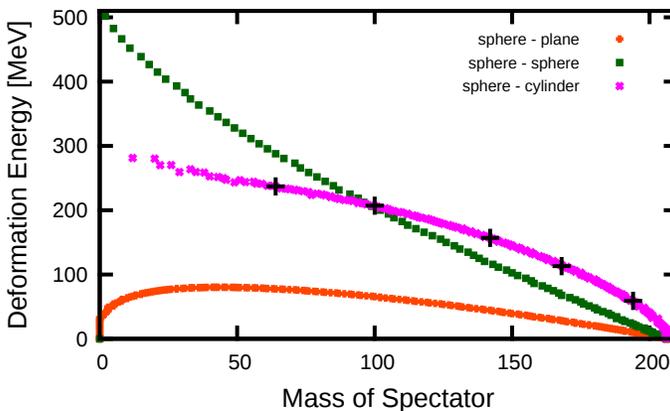


Fig. 2. (Color online) The deformation energy *versus* the spectator mass estimated from its volume. The crosses represent nuclei taken for dynamical evaluation.

excitation energy was taken as the deformation energy from the scenario: “sphere–cylinder” since the energy range fits exactly limitation of our model. From phenomenological point of view, this model is also the most realistic. Five nuclei are chosen arbitrarily for the fission dynamics studies: ^{64}Mn , ^{100}Y , ^{142}Ba , ^{168}Dy and ^{194}Os . The maximal angular momentum was taken as $2\hbar$ just to avoid possible technical problems. For each nucleus, around 300 000 trajectories were calculated and the results are presented in Fig. 3. The panel (a) compares the mean mass coming from an assumption of the symmetric division of the fissioning nucleus (dashed/red line) and centroid of the mass distribution whose width is shown in panel (b). The medium mass system demonstrates a slight deviation from the static assumption as during the path to fission various particles (protons, neutrons, deuteron, α particles) are emitted. Looking at panel (c) for the Total Kinetic Energy (TKE) of the fission fragments, it is visible that almost linear decrease of the excitation energy (Fig. 2) provides the exponential increase of the average TKE as the system is more and more heavy. The effects of inertia and/or the excitation energy is also visible in panel (f) for the increased fission time.

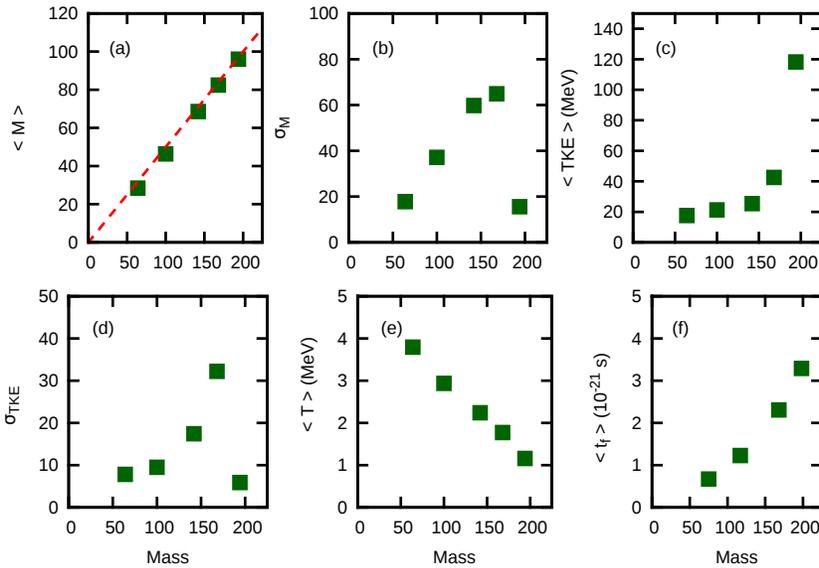


Fig. 3. (Color online) The average values of the mass (a), total kinetic energy (c), temperature (e) and fission time (f) resulting from initial systems with $M = 64$, 100, 142, 168 and 194. The widths of the mass (b) and TKE distribution (d) are also displayed. The dashed line in panel (a) shows the mean mass of the fission fragments assuming symmetric scission of the nucleus.

4. Summary

The preliminary results discussed here show a possibility of investigating the fission of spectators produced in ultrarelativistic heavy-ions collisions. The three scenarios of initial shape estimation for the remnants of the Pb–Pb reaction have been considered to estimate the excitation energies available for various impact parameters. Several observables have been extracted from the state-of-art Langevin calculations for fission dynamics. The average values and widths of mass, charge and total kinetic energy distribution of the fission fragments as well as mean values of the temperature and fission times have been shown for the first time.

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