## WOUNDED QUARKS AT THE LHC\*

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(Received February 16, 2017)

We review the results of the wounded-quark model with a stress on eccentricity observables in small systems. A new element is a presentation of symmetric cumulants for the elliptic and triangular flow correlations, obtained in the wounded-quark approach.

DOI:10.5506/APhysPolBSupp.10.513

This paper is largely based on [1] where more details and results can be found. Historically, the wounded-quark model [2–5] was introduced shortly after the phenomenological success of wounded-nucleon model [6]. The approach stems from the Glauber model [7] adapted to inelastic production [8]. Our results contribute to the on-going discussion on the nature of the initial stages of the ultra-relativistic collision and the relevant degrees of freedom taking part in the early production of entropy/energy in the fireball: Are these nucleons, quarks, partons, random fields, gluonic hot-spots? As the combinatorics of the production depends on the number of constituents, showing in the dependence of the particle production on centrality, it offers a possibility to assess the number of active constituents without a detailed reference to their physical nature.

In [9], it was noticed that the RHIC data can be explained within a wounded-quark model, where the linear scaling  $\frac{dN_{ch}}{d\eta} \sim Q_W$  is used, with  $Q_W$  denoting the number of wounded quarks. The model was further advocated by the PHENIX Collaboration [10, 11]. The wounded-quark scaling also works for the SPS [12]. More recent development in this direction can be found in [10, 11, 13–17].

<sup>\*</sup> Presented by W. Broniowski at the "Critical Point and Onset of Deconfinement" Conference, Wrocław, Poland, May 30–June 4, 2016.

Modeling at the subnucleonic level allows for examination of p + p collisions with the techniques typically used for larger systems. In Fig. 1, we show eccentricities resulting from the wounded quarks in fireballs formed in p+p collisions at the LHC. We note that both the ellipticity and triangularity are large, hence may lead to substantial harmonic flow, in accordance to the collectivity mechanism expected for collisions of small systems [18–21] with sufficiently high multiplicity. The two sets of curves, thinner and thicker, correspond to two values of the source smearing parameter [1]. We note that the eccentricities do not strongly depend on the number of produced particles (centrality), which shows that they originate from fluctuations and not the geometry of the collision.



Fig. 1. Ellipticity  $\varepsilon_2$  (solid lines) and triangularity  $\varepsilon_3$  (dotted lines) of the fireball in p+p collisions at  $\sqrt{s} = 7$  TeV. The thin and thick lines correspond to the Gaussian smearing parameter  $\sigma = 0.4$  fm and 0.2 fm, respectively. The eccentricities are plotted as functions of the mean charged multiplicity at  $|\eta| < 2.4$ .

As a new result of the simulations in the wounded-quark model as implemented in [1], we present the ellipticity-triangularity correlations. We use the measure in the form of a symmetric cumulant, as introduced by the ALICE Collaboration [22]

$$\frac{\mathrm{SC}(a,b)}{\langle a^2 \rangle \langle b^2 \rangle} = \frac{\langle a^2 b^2 \rangle - \langle a^2 \rangle \langle b^2 \rangle}{\langle a^2 \rangle \langle b^2 \rangle} \,. \tag{1}$$

As for a given reaction, the elliptic and triangular flow coefficients are roughly proportional to the corresponding eccentricities

$$v_n \simeq \kappa(n,c)\epsilon_n, \qquad n=2,3,$$
 (2)

with c indicating centrality, one obtains the approximate relation

$$\frac{\mathrm{SC}(v_2, v_3)}{\langle v_2^2 \rangle \langle v_3^2 \rangle} \simeq \frac{\mathrm{SC}(\epsilon_2, \epsilon_3)}{\langle \epsilon_2^2 \rangle \langle \epsilon_3^2 \rangle},\tag{3}$$

where the dependence on  $\kappa(n,c)$  has been canceled out.

The results for Pb+Pb collisions at the LHC are presented in Fig. 2. We note that the predictions from the wounded-nucleon and wounded-quark models are similar to each other and follow the trend of the data [22], indicated with points. The negative sign and the fall-off with centrality are properly reproduced.



Fig. 2. The correlation between the elliptic and triangular eccentricities for Pb+Pb collisions at the LHC, expressed via the symmetric cumulant. The lines indicate the calculations in the wounded-nucleon and wounded-quark models, whereas the points correspond to the ALICE data [22].

The predictions for small systems are shown in Fig. 3. We also plot there the standard Pearson's correlation coefficient

$$\rho(a,b) = \frac{\langle ab \rangle - \langle a \rangle \langle b \rangle}{\sqrt{(\langle a^2 \rangle - \langle a \rangle^2)(\langle b^2 \rangle - \langle b \rangle^2)}} \,. \tag{4}$$

For p+Pb and d+Au reactions, the coefficient  $\rho(\epsilon_2, \epsilon_3)$  follows closely the symmetric cumulant measure, whereas for p + p collision it is substantially different. Such correlations as displayed in Fig. 3, which are large in the considered approach, may be checked in future data analyses of flow correlations in small systems.



Fig. 3. Predictions for the correlation measures between the elliptic and triangular eccentricities in small systems, obtained from the wounded-quark model for three sample reactions.

We wish to thank Ante Bilandzić for a useful discussion concerning symmetric cumulants. This research was supported by the National Science Centre, Poland (NCN) grants 2015/17/B/ST2/00101, 2012/06/A/ST2/00390, and 2015/18/M/ST2/00125.

## REFERENCES

- [1] P. Bożek, W. Broniowski, M. Rybczyński, *Phys. Rev. C* **94**, 014902 (2016).
- [2] A. Białas, W. Czyż, W. Furmański, Acta Phys. Pol. B 8, 585 (1977).
- [3] A. Białas, K. Fiałkowski, W. Słomiński, M. Zieliński, Acta Phys. Pol. B 8, 855 (1977).
- [4] A. Białas, W. Czyż, Acta Phys. Pol. B 10, 831 (1979).
- [5] V.V. Anisovich, Yu.M. Shabelski, V.M. Shekhter, *Nucl. Phys. B* 133, 477 (1978).
- [6] A. Białas, M. Błeszyński, W. Czyż, Nucl. Phys. B 111, 461 (1976).
- [7] R.J. Glauber, in: Lectures in Theoretical Physics, Vol. 1, W.E. Brittin, L.G. Dunham (eds.), Interscience, New York 1959, p. 315.
- [8] W. Czyż, L.C. Maximon, Ann. Phys. 52, 59 (1969).
- [9] S. Eremin, S. Voloshin, *Phys. Rev. C* 67, 064905 (2003).
- [10] S.S. Adler *et al.*, *Phys. Rev. C* **89**, 044905 (2014).
- [11] A. Adare *et al.*, *Phys. Rev. C* **93**, 024901 (2016).
- [12] P. Kumar Netrakanti, B. Mohanty, *Phys. Rev. C* 70, 027901 (2004).
- [13] C. Loizides, J. Nagle, P. Steinberg, *SoftwareX* 1-2, 13 (2015)
  [arXiv:1408.2549 [nucl-ex]].
- [14] R.A. Lacey et al., arXiv:1601.06001 [nucl-ex].
- [15] L. Zheng, Z. Yin, Eur. Phys. J. A 52, 45 (2016).
- [16] J.T. Mitchell, D.V. Perepelitsa, M.J. Tannenbaum, P.W. Stankus, *Phys. Rev. C* 93, 054910 (2016) [arXiv:1603.08836 [nucl-ex]].
- [17] C. Loizides, *Phys. Rev. C* 94, 024914 (2016) [arXiv:1603.07375 [nucl-ex]].
- [18] P. Bożek, *Phys. Rev. C* **85**, 014911 (2012).
- [19] P. Bożek, W. Broniowski, *Phys. Lett. B* **718**, 1557 (2013).
- [20] P. Bożek, W. Broniowski, *Phys. Rev. C* 88, 014903 (2013).
- [21] P. Bożek, W. Broniowski, G. Torrieri, *Phys. Rev. Lett.* **111**, 172303 (2013).
- [22] J. Adam et al., Phys. Rev. Lett. 117, 182301 (2016) [arXiv:1604.07663 [nucl-ex]].